## **ON-CHIP WIRES: SCALING AND EFFICIENCY**

A DISSERTATION SUBMITTED TO THE DEPARTMENT OF ELECTRICAL ENGINEERING AND THE COMMITTEE ON GRADUATE STUDIES OF STANFORD UNIVERSITY IN PARTIAL FULFILLMENT OF THE REQUIREMENTS FOR THE DEGREE OF DOCTOR OF PHILOSOPHY

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## Abstract

Recent years have seen an increase in the importance of on-chip wires, as they have slowed down and gates have sped up. This dissertation takes a close look at the story of wire scaling. It forecasts wire and gate characteristics from the Semiconductor Industry Association roadmap and combines them into performance metrics, showing how the ratio of wire delays to gate delays scales slowly for scaled-length wires and grows rapidly for fixed-length wires.

This duality of "fast local wires" contrasted with "slow global wires" affects how we approach VLSI designs. First, CAD place-and-route tools must improve to keep up with growing die complexity and more local blocks gathered on a chip. Second, modular architectures can effectively exploit the dual nature of wires, using wide global buses of high bandwidth to offset long wire latencies.

Using such wide and long global buses can burn a great deal of power, especially if built with traditional delay-optimal CMOS repeaters. Traditional repeaters can be sized and spaced to save about 30% in energy for only a 10% delay penalty. Because this 30% of energy savings is not a lot, techniques for running global wires at a reduced voltage can be very important. These include NMOS drivers, overdrive pre-emphasis, and voltage pre-equalization. Using these circuit techniques offers an order-of-magnitude in energy savings for no effective slowdown. Experimental results on a 180nm testchip validate this 10x savings in energy over 10mm long on-chip buses, running at 1 token per 10 gate delays. Further experimental data shows receiver input offsets under 90mV, with input offset compensation leading to residual input uncertainties of around 15mV.

## Acknowledgements

The title page of this doctoral thesis lists my name as the single author, although that is a deception. The ideas and concepts in the pages that follow arose from the work of many people, and to them I owe my thanks.

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## **Chapter 1**

## Introduction

The importance of on-chip wires has dramatically risen over the past decade. Prior to the early 1990s, chip designers could treat on-chip wires as purely capacitive loads of logic gates; these wires had no intrinsic delays of their own. As technologies scaled into the mid-1990s, growing wire resistance coupled with shrinking native gate speeds made wire delays increasingly important. For example, the principal and enduring speedpath on the PentiumPro/II/III architectures at Intel, designed in the early-to-mid 1990s, was a long write-back bus whose wire length spanned much of the chip.

About this time, a now-famous graph projected wire and gate delays into what was then the distant semiconductor future (see Figure 1.1). James Meindl noted wryly that this graph, based on an early study from the 1980's [1], should have been designated the logo of the Semiconductor Research Corporation, so often was it invoked at industry meetings [2]. It showed a clear divergence between gate delays and wire delays, and the use of copper rather than aluminum only delayed the eventual cross-over. This graph raised more questions than it answered, such as how long were the wires, what kind of gates could run so fast, and whether speed-up technologies like repeaters were used. Regardless, it conveyed its point clearly: like it or not, wires are getting slower as gates are getting faster.

The seemingly inexorable divergence between improving gate delays and degrading wire delays seems to paint a bleak future for chip designers. But does it? After all, the industry continues to produce more and more chips of enormous complexity, with designers

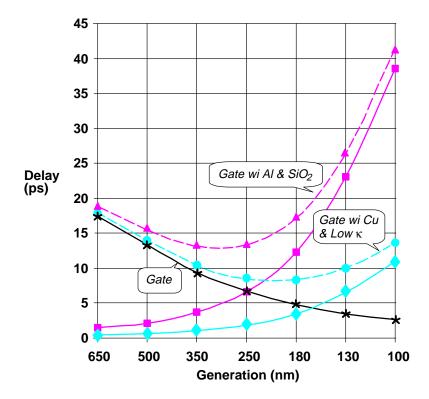


Figure 1.1: The view of the future, circa 1997

constantly creating new architectures and circuits for ever-increasing performance. Chip builders, it seems, must know things that Figure 1.1 does not express.

This dissertation takes a close look at the story of wire scaling: how do wires perform now, and how will they perform in the future? What are designers doing today and what will they do tomorrow? As we shall see, managing wire delay as technology scales is important, but doing so in an energy-efficient manner is just as critical.

### 1.1 Organization

To understand the effects of technology scaling on on-chip wires we must first decide what wire metrics are important. Chapter 2 describes models of wire and gate delays. These models allow us to project the wire and gate delays in future technologies. Notice that only their *relative* performance matters. If gate delays and wire delays changed identically, either dramatically or even not at all, then the overall system design would still be balanced, and neither gates nor wires would individually limit overall performance.

Chapter 3 takes the results from Chapter 2 and looks at the implications of such scaling trends. Although wire scaling affects the entire VLSI design space, I limit my attention to two areas in particular: CAD tools and their ability to automate design, and computer architectures and how they can leverage or exploit the scaled characteristics of wires. In discussing these latter architectural implications of scaling, I will consider long on-chip wires and how energy efficiency on these wires will become an important design constraint.

Figuring out how to drive long wires is not a new problem. A number of solutions using optimally-repeated CMOS repeaters are well-known. However, the problem of driving these wires efficiently is rarely considered. Being stingy with delay typically means spending extravagently in power, and repeaters with slightly sub-optimal delays can offer potentially large energy savings. Chapter 4 examines these tradeoffs.

Tweaking repeater sizing and placement can buy only limited energy savings. Chapter 5 departs from the full-swing CMOS repeater discussion to consider low-swing circuit architectures. It discusses design issues related to low-swing drivers, receivers, and the wires themselves. Chapter 6 follows and gives results from testchip experiments.

## **Chapter 2**

## Metrics, models, and scaling

In order to lay a consistent foundation for the following sections, this chapter considers how to model the delays of wires and gates. A contemporary  $0.18 - \mu m$  technology will provide a framework for our discussion of gate and wire delays. Gates will use a delay metric called an *FO4*, which is based on feature size and described below; wires will use simple cross-sectional-area-based models for resistance and capacitance that lead to correspondingly simple wire delay metrics. These measures of gate and wire performance can help project trends of wire delays relative to gate delays, including side effects such as noise. As technologies scale, some wires scale in length and others do not, and we can apply our metrics to both of these types of wires. As we shall see, the scaled-length wires do trend with gate delays, but the fixed-length wires do not.

### 2.1 A simple gate delay model

Transistors are very complicated devices that can be connected in a myriad of ways, so characterizing gate performance in a way that facilitates comparisons to wires initially appears difficult. However, for reasons of productivity, CAD tool support, and robust behavior, VLSI designers use transistors in only a very limited set of topologies; static and simple dynamic CMOS gates dominate digital designs. As a result, metrics that characterize the performance of these gates will suffice.

To measure gate speed, we use the delay through an inverter driving four identical

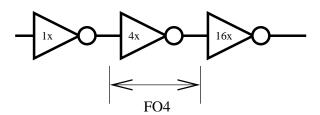


Figure 2.1: A fanout-of-four inverter delay

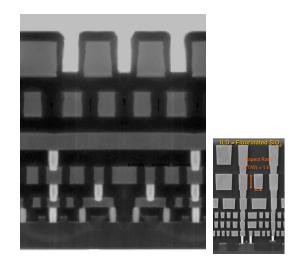
copies of itself, shown in Figure 2.11. This is called a "fanout-of-four inverter delay," or simply an *FO4*. In our benchmark 0.18- $\mu$ m technology, an *FO4* is 90ps.

Combinational delays, composed of many different static and dynamic CMOS gate delays, can be normalized to this *FO4*. The resulting relative delay holds constant over a wide range of process technologies, temperatures, and voltages. We can therefore treat the *FO4* as our proxy for CMOS gate delays: to understand how gates compare with wires, we need estimate the delay of only a single loaded inverter.

*FO4* delays are generally constant for a given feature size; that is, one company's 180nm technology will produce an *FO4* delay reasonably close to another company's 180nm technology. This simplifies estimating *FO4*s. However, specialized in-house technologies (such as those from Intel or IBM) typically offer 30-50% faster *FO4* delays due to highly optimized—and expensive—gate engineering. We will discuss this further in Section 2.4.

### 2.2 Wire characteristics

Two generations of Intel process technologies, shown as cross-sectional photographs in Figure 2.2, reflect advances in on-chip wiring. A 0.25- $\mu$ m technology from 1997, the vehicle for 700MHz Pentium III processors, used five layers of aluminum metal, with shiny tungsten plugs connecting them [3]. A 0.13- $\mu$ m technology from 2002, the vehicle for 3GHz Pentium 4 processors, uses six layers of copper metal with copper vias [4]. Between these two technologies, spanning three generations, the wires' cross-sectional areas and spacings fell dramatically. The importance of cross-sectional area and spacing lies in their effects on the wire electrical characteristics that we care about: resistance and capacitance.



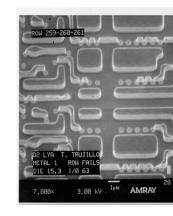


Figure 2.2: Drawn to relative scale: A  $0.25-\mu m$  aluminum interconnect technology (left) and a  $0.13-\mu m$  copper interconnect technology (right)

For the forseeable future, resistance and capacitance will determine wire delay and noise behavior. The following sections describe geometric models for resistance and capacitance that are based on cross-sectional area and spacing. The simplicity of these models does not preclude them from including many of the non-idealities of real wires. We will also discuss wire inductance and why we can ignore its role in performance modeling.

### 2.2.1 Resistance

All wires have a finite conductance, representing the ability of the wire to carry a charge flow. Aluminum wires have a resistivity of  $3.3 \text{m}\Omega \cdot \text{cm}$ , while newer thin-film copper wires, used in most contemporary processes, have a resistivity of  $2.2 \text{m}\Omega \cdot \text{cm}$ . Resistance (per unit length) may be approximated by the material resistivity divided by the conductor's cross-sectional area, but several wire non-idealities affect this model.

For copper wires, a thin barrier layer on three sides of the wire, required to prevent copper from diffusing into the surrounding oxide (see Figure 2.3) raises resistance by decreasing the wire's effective cross-sectional area. Similarly, due to surface planarization

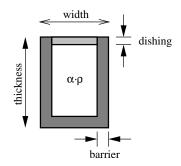


Figure 2.3: A simple resistance model

and proximity effects to neighboring lines, wide wires may be over-polished, or "dished," again reducing effective cross-sectional area. Finally, electrons inelastically scatter off lattice bonds at the edges of wires. As the wire dimensions grow smaller and smaller, this will reduce the mean free path of electrons, effectively increasing the material resistivity [5]. A modern 0.18- $\mu$ m technology has a copper barrier thickness of 17nm, negligible dishing for normally-sized wires, and negligible carrier scattering [6].

Assuming that the barrier layer is conformal, a reasonable assumption with fabrication techniques such as atomic layer deposition, we can write wire resistance as

$$R_{wire} = \alpha_{scatter} \cdot \frac{\rho}{(thickness - barrier - dishing)(width - 2 \cdot barrier)}$$
(2.1)

We can ignore skin effects for the vast majority of on-chip wires, because wires are under a few skin depths thick and wide. Digital gates rarely swing with transition times faster than an *FO4*, or 90ps in our 0.18- $\mu$ m technology, which corresponds to frequency components of 1.5GHz<sup>1</sup>. At that frequency, copper has a skin depth of 1.7 $\mu$ m, or nearly 20 $\lambda$ , exceeding most wire dimensions. A signal's edge rate, and hence frequency content, scales with technology, but so do wire dimensions. Because skin depth falls with the square root of frequency, wires will shrink faster than their skin depth, making this issue less important under scaling.

<sup>&</sup>lt;sup>1</sup>For a signal switching with a transition time of *t*, the "knee" frequency of that signal's harmonics sits at  $(2\pi t)^{-1}$ . This the frequency at which the spectral amplitude is half that of the standard 20db/decade rolloff. A common misconception is that high-frequency effects become less important if we slow down clock rates. A more accurate statement is that they become less important only with slow edge rates.

At upper layers, metal widths may well exceed two skin depths, especially for those signals carrying power or ground. Designers typically route power and ground wires right next to each other to maximize decoupling capacitance, and by the "proximity effect," currents in the two wires flow as close to each other as possible, making horizontal skin depth important. For those specialized wires, skin effects will prompt designers to break the wide wires up into fingers.

Plugs, or vias, between aluminum metal layers were made of tungsten, and tended to be fairly resistive; in a 0.25- $\mu$ m process a M1-M2 via resistance was about 5 $\Omega$  and vias from M5 down to the substrate added up to more than  $20\Omega$ . This may seem large considering a 1 $\mu$ m wide, 1mm long M5 line itself had a total resistance of only  $20\Omega$ , but by arraying many vias together, designers could easily reduce plug resistance; in most cases, self-heat and electromigration checks required arrayed vias for long wires anyway. Copper processes improve via resistance by depositing vias at the same time as wires. These copper vias are much less resistive and do not need to be as aggressively arrayed, although some recent experience has shown that copper vias have their own electromigration concerns: they serve as nucleation sites, gathering voids that flow down the copper wires much like tumbleweeds [7]. Copper vias thus tend to be arrayed like aluminum vias, only for improved reliability and not reduced resistance.

### 2.2.2 Capacitance

All wires have capacitance, modeling the charge that must be added to change the electric potential on the wire. Some analytical models approximate the capacitance of a wire over a plane; more accurate ones combine a bottom-plate term with a fringing term to account for field lines emerging from the edge and top of the wire. However, wires today are taller than they are wide, and will grow even taller to reduce resistance as technologies scale. At minimum pitch their side-to-side capacitances are a significant and growing portion of the total. Capacitance is thus better modeled by four parallel-plate capacitors for the top, bottom, right, and left sides, as shown in Figure 2.4 [8], plus a constant. This extra term lumps all the fringing field terms together and approximates their sum as a constant.

Here, the only non-ideality we need to consider are vertical and horizontal capacitors

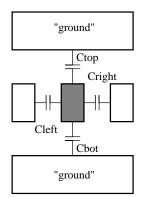


Figure 2.4: A simple capacitance model

that have different relative dielectrics. This may be due to air gaps in the intra-layer oxide (faintly visible in the left picture in Figure 2.2), or due to intentional differences in technologies that leverage low- $\kappa$  materials [9]. In this case, we can use

$$C_{wire} = \varepsilon_0 \left( 2M \varepsilon_{horiz} \frac{thickness}{spacing} + 2\varepsilon_{vert} \frac{width}{ILD_{thick}} \right) + fringe(\varepsilon_{horiz}, \varepsilon_{vert})$$
(2.2)

That we can lump the fringe terms into a constant value (though one which changes for different technologies as the dielectric constants  $\varepsilon_{horiz}$  and  $\varepsilon_{vert}$  change) is somewhat surprising. This is really just a result of curve-fitting, and comes about because the fringe terms vary only logarithmically with spacing.

The "far" plates for the top and bottom capacitors are typically modeled as being grounded: they represent a collection of orthogonally-routed conductors that, averaged over the length of the wire, maintain a constant voltage. This capacitance would be multiplied by an appropriate factor if the orthogonal wires switched simultaneously and monotonically, as with a precharge bus. Capacitors to the left and right, on the other hand, have data-dependent effective capacitances that can vary: if the left and right neighbors switch in the opposite direction as the wire, the effective sidewall capacitances double, and if they switch with the wire, the effective sidewall capacitances approach zero. We model this multiplication effect by varying the M parameter in Equation 2.2 between 0 and 2 (our simple model ignores the effects the M term has on the fringe capacitance). These left and right neighbors are also the worst offenders for noise injection. The fringe term depends only

weakly on geometry and for today's  $0.18 \ \mu m$  technologies with homogenous dielectrics is about 40 fF/ $\mu m$ . For the very top layers of metal with no upper layers, we can use three parallel plates with extra fringing terms on the two horizontal capacitors.

#### 2.2.3 Inductance

All wires also have inductance, representing an inertia against changing current through the wire. Unlike resistance or capacitance, inductance has no handy closed-form models. Freshman physics taught us to think about the inductance of a loop, and how a changing magnetic flux through that loop induces a voltage on it. This view of inductance does not easily model on-chip wires, however, because we do not always know what structures form the "loop": if we send current down an on-chip wire, the return currents may flow in adjacent wires, parallel power supply buses, or even the substrate. In fact, due to return currents flowing in the paths of least impedance, the actual current loops will change with the frequency content of the signal. At low frequencies, wide low-resistance power buses, even if far away, have low impedance ( $Z = R + j\omega L$ ), leading to fairly large loops and hence higher inductance. At high frequencies, far-away return paths have unattractively high impedances, and return currents will bypass them to return in local, capacitively-coupled wires, implying lower inductance but higher path resistance.

This "chicken-and-egg" problem is the basic challenge in calculating inductance: we cannot know the inductance until we know what the loop (or loops) are. But we cannot discern the correct loops until we know what the inductance is. To bypass this problem, today's tools define return paths to be at a fixed common reference,<sup>2</sup> and the resultant "partial inductances," when combined with wire resistances and capacitances, can yield accurate results inside circuit simulation [11][12].

Figure 2.5 shows an example of using partial inductance. Here, wire *ab* carries current from a driver to a load, and this current can return in wires *cd* or *ef* (two power supply lines). There would be many more potential return paths in a realistic layout. Figure 2.6 shows the partial inductances for each of these wire segments. The partial "loops" terminate

<sup>&</sup>lt;sup>2</sup>The common return path can be arbitrarily picked, so long as it is consistent. The most mathematically convenient return path is at infinity; however, this choice leads to physically non-intuitive numbers. As Larry Pillegi noted, "visualizing a loop at infinity is somewhat like drawing God." [10]

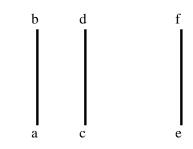


Figure 2.5: A signal wire *ab* and two potential returns *cd* and *ef* 

at infinity, so their flux areas extend out to the right past the page edge. These partial inductances come only from the geometric information about these wires and not from any assumptions of the current loops. Suppose current upwards in *ab* returns downwards

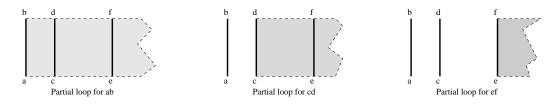


Figure 2.6: Partial loops for the three wires

in *cd*. Then *cd*'s flux area has the inverse polarity, and when overlaid atop the flux area for *ab*, cancels with *ab*'s flux areas to the right of *cd*. This leaves only the area between *ab* and *cd*, leading to a simple loop inductance calculation and hence an impedance for that loop. The same calculation can be performed for current returning in *ef*, similarly leading to an impedance for that loop. Then in a SPICE simulation, the return current will split, preferring the loop with less overall impedance, and allowing overall loop inductance calculations.

Most inductance calculation tools overestimate inductance because they assume all wire current uniformly flows to the end of the line, while in VLSI circuits, most current actually returns through distributed wire capacitances [13]. A larger problem with inductance calculation involves data explosion: inductance falls slowly with distance inside the return loop, so all wires within several pitches—*i.e.* several wires—must be included in the computation. For each extracted wire we must calculate the mutual inductance to multiple

neighbors, and the amount of data to store and compute quickly becomes unmanageable. Sparsification schemes try to reduce this data without destabilizing the resultant coupling matrices [14][15].

To determine whether or not wire inductance is important we need to consider two questions [16][17][18]:

- Does the driver-end of the wire swing slowly enough to avoid transmission-line effects? Quantitatively, does the driver impedance exceed the line impedance (*R<sub>gate</sub>* > 2 · Z<sub>0</sub>)?
- Do resistive losses in the wire outweigh any transision line effects? Quantitatively, does the wire's attenuation factor  $(0.5 \cdot l \cdot R_{wire})/Z_0$  exceed one?

Short signal wires meet the first condition, making inductance unimportant for them. This is because in a contemporary 0.18- $\mu$ m technology, FastHenry simulations of a well-gridded bus show  $L_{wire} = 0.3$ nH/mm with  $C_{wire} = 0.3$ pF/mm, making  $Z_0$  approximately  $30\Omega$ ,<sup>3</sup> much less than short-line driver resistances. Typical on-chip inductance values range from 0.2 - 0.5nH/mm [19]. In a 0.18- $\mu$ m technology, a drive resistance less than 60 $\Omega$  (twice  $Z_0$ ) must be at least 85 $\mu$ m in width, a driver size appropriate for a 2mm-long wire. Under scaling, both drive resistance and  $Z_0$  remain constant, maintaining this driver-impedance inequality.

The second condition is satisfied by wires that span a typical repeater distance or longer. As will be seen in following chapters, an optimally-repeated wire (assuming RC behavior) has a wire length of

$$l = 3 \cdot \sqrt{\frac{R_{gate}C_{gate}}{R_{wire}C_{wire}}}$$
(2.3)

Thus, the attenuation factor can be written

$$\frac{0.5 \cdot l \cdot R_{wire}}{Z_0} = 1.5 \sqrt{\frac{R_{wire}(R_{gate}C_{gate})}{L_{wire}}}$$
(2.4)

<sup>&</sup>lt;sup>3</sup>We approximate  $Z_0 = \sqrt{\frac{L_{wire}}{C_{wire}}}$ , while more accurately,  $Z_0 = \sqrt{\frac{R_{wire} + j\omega L_{wire}}{j\omega C_{wire}}}$ . With edge rates faster than 100ps, and hence signal frequency content exceeding 1.5GHz (10G-rad/s), this approximation holds quite well.

In our 0.18- $\mu$ m technology, with  $R_{wire} = 30\Omega/\text{mm}$  and  $R_{gate}C_{gate} = 10\text{ps}$ , the optimal repeater distance is about 3mm, and the attenuation term shows that line resistance overwhelms inductance unless the wire is shorter than two-thirds the repeat distance, or 2mm. Under scaling, the attenuation constant for a given wire will increase, making it increasingly resistive.

Hence, in a 0.18- $\mu$ m technology, inductance can be safely ignored for wires shorter than 2mm and for wires longer than 2mm. For wires that are exactly 2mm long, simulations show the delay of an *LRC* wire differs from the delay of an *RC* wire by under 3%, much less than the uncertainty in capacitance or resistance extraction. Wire inductance is therefore unimportant for the delay of typical signal wires.

With much wider wires having much lower resistance, such as power lines, or with systems very sensitive to the exact delay modeling, such as clock networks, inductance does play a role in design. For most signal wires, however, inductance effects on delay are largely irrelevant. Inductive noise, which depends on  $M\delta i/\delta t$ , is not as easily dismissed and will be discussed in more detail below.

### 2.3 Wire performance metrics

The discussion of wire characteristics above provides the groundwork for an examination of wire performance. This section will first consider robustness by exploring signal coupling noise issues, and then discuss delay and bandwidth metrics.

### 2.3.1 Signal coupling

Coupling noise is a serious problem for a chip designer, as both mutual capacitance and inductance terms for wires can be large. To understand the magnitude of coupling noise problems, we need to compare the induced noise to the noise margins of the receiving gate. Static and dynamic CMOS gates are voltage controlled—they switch their output voltage when the input voltage exceeds some threshold. Thus we are concerned about the voltage noise on the wire relative to the voltage margins of the receiving gates.

Capacitance noise coupling is a larger effect so we will look at it first. The large aspect

ratios of modern wires mean that for a wire surrounded by neighboring wires on either side, the cross-capacitance to these sideways neighbors dominates the total capacitance; sideways cap can exceed 70% of the total. When these sideways neighbors (the "attackers") switch, the current that flows through the coupling capacitors must then flow through the center wire (the "victim"), inducing noise on it. The familiar model of  $V_{noise} = V_{swing} \frac{C_{coupling}}{C_{total}}$  gives a pessimistic upper bound on the noise, because this is the noise voltage only if the victim line is left floating. Many recent papers have modeled this noise more carefully, and have shown that the noise voltage depends on both the coupling capacitance to total capacitance ratio as well as on the ratio of the strengths of the gates driving the two wires [21][22][23]. A convenient model simple enough for first-order hand calculations is:

$$V_{noise} = V_{swing} \cdot \frac{C_{coupling}}{C_{total}} \cdot \frac{1}{1 + \frac{\tau_{att}}{\tau_{vic}}}$$
(2.5)

where  $\tau_{att}$  and  $\tau_{vic}$  are the time constants of the attacker and victim drivers, respectively. If the attacker has a much smaller time constant than the victim (and is hence much stronger), the noise approaches the pessimistic worst-case. Typically, however, the transition times of different gates are matched, which gives an attacker-to-victim time constant ratio that is greater than one. If the two wires are identical, with identical drivers, the time constant ratio will be set by the difference between the effective resistance of a MOS transistor in the saturated region, driving the aggressor wire, and a transistor in the linear region, trying to hold the value of the victim wire stable. This ratio is usually between two and four,<sup>4</sup> which greatly reduces capacitive coupled noise for most nodes.

However, the limitation of this model is that it does not account for distributed line resistance. Adding this effect makes deriving analytical results difficult, leading researchers to use approximations like lumping the wire resistance with the driver resistance [22]. However, for the special case where the wires are identical, the most common case where coupling is a problem, there is a way to view the problem using superposition that gives a simple and intuitive view of coupling. This model starts by assuming that the driver resistances are the same, as shown in Figure 2.7.

<sup>&</sup>lt;sup>4</sup>The ratio hinges on the degree of velocity saturation of the attacking transistor. Since nMOS gates suffer more from velocity saturation, the ratio for nMOS gates is generally closer to 4.

#### 2.3. WIRE PERFORMANCE METRICS

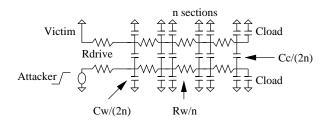


Figure 2.7: Bus coupling noise model

The key to the analysis is to break the driving input into a symmetric, or even mode, input (both sides are driven by a  $\frac{V}{2}$  ramp), and an antisymmetric, or odd mode, input (attacker driven by  $\frac{V}{2}$  ramp and the victim driven by  $-\frac{V}{2}$ ). In the even component, both attacker and victim see a half-amplitude input, and because the two lines now have identical responses, the coupling capacitors conduct no current and can be zeroed out. In this case, the response at the end of the victim is the same as that of a single wire in isolation, with total line capacitance  $C_{total} = nC_s + C_l$ .

For the odd component, the attacker sees a positive half-amplitude step, and the victim sees a negative half-amplitude step. In this case, the two lines have exactly opposite responses, so the coupling capacitors see twice the voltage difference and can be replaced by double-size capacitors referred to ground. Thus we can once again treat the victim wire as an isolated single wire, with total line capacitance  $C_{total} = n(C_s + 2C_c) + C_l$ .

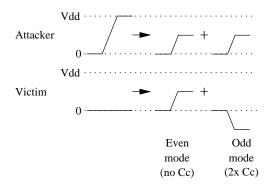


Figure 2.8: Decomposition of attacker and victim waveforms

The combination of the even and odd modes, as in Figure 2.8, will place a full step

on the attacker driver and hold the victim driver to ground, so we need add only the two decoupled responses to get the true victim waveform. In other words, the victim response can be written as the sum of two isolated wire responses, one with no coupling, and the other with double coupling. These two isolated responses can be derived from a number of models, ranging from simple single time-constant exponentials to more complicated moment-matched asymptotic waveforms [24]. The key idea is that symmetry properties allow us to break the highly-coupled circuit into two isolated circuits that are more easily handled.

Note that this model requires that the attacker and victim lines have completely identical resistances and capacitances; in particular, we need them to have the same driver resistances. Yet the driver of the victim wire, a transistor in its linear mode, typically has a lower (stronger) resistance than the saturated transistor driving the attacker wire.

We avoid this limitation by observing that a driving resistor that sees a step input can be transformed into a larger (weaker) resistor by using a slower exponential input. In other words, from the perspective of the downstream wire, a properly-chosen exponential input driven into a resistor is almost indistinguishable from a step input driven into a larger (weaker) resistor. Thus if we use an appropriate exponential input instead of a step input, and the smaller (stronger) victim resistance for both of the wire models, we will effectively increase the attacker driving resistance while maintaining the proper victim resistance.

The mathematical derivation using simple single-time constant models for the wire responses reduces to a peak noise given by:<sup>5</sup>

$$V_{peaknoise} = \frac{C_{coupling}}{C_{total}} \cdot \left(\frac{1+M}{k+M}\right)^{\frac{k+M}{k-1}}$$
(2.6)

$$M = \frac{nR_{wire}}{2R_{att}} \tag{2.7}$$

where k is the ratio of attacker to victim driving resistances (typically between two and four). For reasonable wire lengths, the driver resistance ratio does a good job of attenuating the noise pulse, making it a small issue for static CMOS circuits. However, capacitance coupling is a large problem for weakly-driven nodes, and CAD tools must be used to check

<sup>&</sup>lt;sup>5</sup>Note that this formula reduces to a slightly different result than Equation 2.5 when the wire resistance is 0 (*i.e.* when M = 0). In these cases, this equation gives a better result.

#### 2.3. WIRE PERFORMANCE METRICS

for coupling on such weakly-driven or dynamic nodes.

Noise from inductive coupling can also present problems for VLSI wires. The current flowing down the aggressor wire generates a magnetic field which causes a backwards return current to flow in the victim wire. Inductive coupling pushes the victim in the opposite direction from capacitive coupling: a rising attacker capacitively couples a victim up, but inductively couples the victim down. While capacitive coupling is mostly a "nearest neighbor" phenomenon, inductive coupling has a much larger range. Inductive noise becomes a problem only when a large number of wires switch at the same time in bus-like situations [25][26][27]. The worst-case noise vector would have multiple wires switching, with near neighbors switching in one direction, and far neighbors switching in the opposite direction. This causes the capacitive and inductive noises to add, and the accumulated noise can be enough to cause failures [26].

Designers cope with inductive coupling by adding power planes or densely gridded power supplies to reduce the number of wires that can couple into a victim. Power planes, or dense power grids, effectively reduce both self and mutual inductances for wires in the direction of the grid, because they provide very nice return paths within a few microns of the wire itself and thus limit the extent of the magnetic coupling [28]. Most companies have design rules for buses, such as requiring every fifth wire to be a power supply wire, which makes inductive noise much less than capacitive noise and under 5% of the power supply.

#### 2.3.2 Wire delay

The delay of an on-chip wire can be modeled by a simple RC formulation. Here, we treat a CMOS driver as a simple resistor  $R_{gate}$  with a parasitic load  $C_{diff}$ . The CMOS receiver at the other end of the wire presents a capacitive load  $C_{gate}$ .

$$Delay \propto R_{gate}(C_{diff} + C_{wire} + C_{gate}) + R_{wire}(\frac{1}{2}C_{wire} + C_{gate})$$
(2.8)

By approximating the CMOS driver with a simple resistor, this model ignores both non-linear drive resistance as well as the effect of slew rate on delay.

This model takes advantage of the small effects inductance has on delay: it includes

only R and C terms. Inductance can affect wire delay in four ways, but all are insignificant. First, signal propagation is limited by the speed of light down the wire, set by  $\sqrt{LC}$ . When the driver end of the wire switches, the receiver end cannot begin to transition until at least  $l\sqrt{LC}$ . This effect is insignificant: as discussed previously, wires long enough to make this propagation delay important also have wire resistance that swamps out wire inductance. Second, dramatically under-damped systems will ring, making single-time-constant models poor predictors of delay. But keeping driver fanouts reasonable (*i.e.* not smaller than unity) prevents wires from ringing and keeps the "sharpening" effect of inductance to a small percentage of total line delay. Third, inductive coupling, much like capacitive coupling, can push out delay by forcing a victim to absorb induced transients before swinging. With adherence to some inexpensive design heuristics, however, inductive coupling can be made trivial in comparison to capacitive coupling. Fourth, loop inductance can force return currents into tighter loops with higher resistivity than wider loops. This extra "return path resistance," often overlooked by designers, is typically modeled by simply increasing the  $R_{wire}$  term. This model is imperfect because the return path resistance term actually appears after the load capacitor, so the increase in  $R_{wire}$  is typically done with a correction factor that has been curve-fit to match accurate delay simulations of wiring templates. This scheme, though somewhat inaccurate, has the virtue of easily fitting into standard CAD timing flows.

The first term in Equation 2.8 is about 1*FO4*, as simple sizing heuristics aim for gate sizes to have a fanout of about four for optimal delay [73]. Long wires with large capacitive load might thus imply huge gates, but designers typically use higher fanouts for such long wires; because wire resistance shields downstream capacitance from the drivers, higher fanouts are more efficient. In these long wire cases, the 1*FO4* approximation is somewhat optimistic. We will also assume that our wires are fairly long, so that  $C_{wire} \gg C_{gate}$ . Our metric for delay is therefore simply  $1FO4 + \frac{1}{2}R_{wire}C_{wire}$ . These assumptions do not hold for wires driving large or many gate loads, such as repeated wires (which we will consider later) or control wires driving each bit of a wide datapath. Representative delay numbers for a 0.18- $\mu$ m technology are shown in Table 2.1; this table uses a total capacitance whose cross-capacitance term is Miller-multiplied by a factor of two, to simulate a data-dependent worst-case delay.

#### 2.3. WIRE PERFORMANCE METRICS

	Copper wire delays FO4/mm <sup>2</sup>
Local wire	0.56
Semi-global wire	0.22
Global wire	0.05

Table 2.1: Sample  $\frac{1}{2}R_{wire}C_{wire}$  delays, 0.18- $\mu$ m technology

As Section 2.5 describes in more detail, modern technologies optimize their metal layers for three different tasks. Local wires run on the lowest level of interconnect; the semiglobal wires, on mid-level layers of metal, typically run within functional units; and the global wires, on the top layers of metal, route power, ground, and global signals. The wire delay for all three classes of wires are given in the table. For a copper 0.18- $\mu$ m technology, long unbuffered wires with small loads are not too slow. A 10mm route takes 1+56=57 *FO4*s on local wires, but 1+22=23 *FO4*s on semi-global lines, and 1+5=6 *FO4*s on global wires.

We can also estimate the bandwidth of an unbuffered wire by asking how long we must wait between successive transitions on a wire. If we switch a wire once, we need to wait until residual currents from that transition have mostly died away, or else we will see intersymbol interference when we switch the wire again. We can do this by waiting for three propagation delays before sending the next signal.

Figure 2.9 shows idealized (linear driver, no noise and no process variations) waveforms from a model of an inverter driving a long wire<sup>6</sup> in a wave-pipelined system. Waiting three time constants ( $\tau$ 's) between tokens allows each waveform step to transition to 95% of the swing; going any faster closes the data "eye" unacceptably. The perfectness of the waveforms underscores the need to wait at least three  $\tau$ 's between tokens—any slight noise or imperfections in the system would close the data eyes further. In reality, real designs, even if wave-pipelined, never achieve these rates: the margining to account for process, voltage, and temperature variabilities would significantly slow the token rate. However, we can still use this three- $\tau$  repeat rate as a safe upper limit for wire bandwidth.

<sup>&</sup>lt;sup>6</sup>This example uses a linear driver resistance of  $36\Omega$ , total gate capacitance of 0.75pF, diffusion capacitance of 0.375pF, wire resistance of  $200\Omega$  and wire capacitance of 0.8pF. The wire is written as a  $\pi$  model.

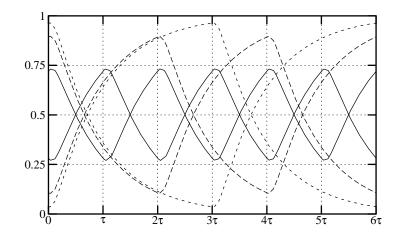


Figure 2.9: Idealized wave-pipeline with 1-, 2-, and 3-t repeat rates

In the equation below, we assume the propagation delay to be a gate delay (FO4) plus the distributed wire delay. Increasing a wire's pitch will monotonically increase that wire's bandwidth, because it decreases the wire RC product, leading to the misleading result that fatter wires are always better. Therefore, we will actually examine the bandwidth across a routing area. In this case, making wires excessively fat will reduce the number of wires available, and hence potentially reduce bandwidth over that area:

$$BW_{area} = \frac{1}{3(1FO4 + \frac{1}{2}R_{wire}C_{wire})} \cdot \frac{Blockwidth}{Wirepitch}$$
(2.9)

This formulation allows us to examine unrepeated bandwidth in both local and global contexts. For module-length wires, we run semi-global layer metals across a square that holds around 50,000 gates. For global wires, we run top-level metals across a 2cm die and thus consider the bandwidth across a die-sized square.

Figure 2.10 shows module and global unrepeated bandwidth. In Equation 2.9, the lefthand term rises with increasing wire pitch, but the right-hand "number-of-wires" term falls with increasing pitch. Whether or not designers should increase the wire pitch depends on the wire length: if the wire is short enough that its delay is dominated by gate delay, then the bandwidth improvement from increased pitch tends to be less than the bandwidth degradation from fewer wires. If the wire is long enough that its delay dominates gate delay,

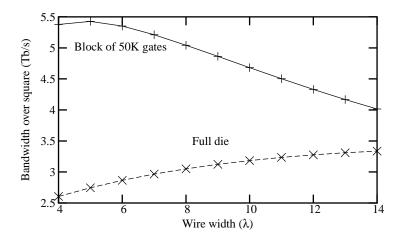


Figure 2.10: Unrepeated bandwidth, in an  $0.18 - \mu m$  technology

then bandwidth is improved by increasing pitch. In Figure 2.10, we see that increasing wire width does not improve local bandwidth, but it slightly improves global bandwidth.

The long delay and low bandwidth of the global wires clearly indicates a problem caused by the large resistance of these wires. Fortunately, there is a simple way to dramatically reduce the effect this resistance has on circuit performance—we can break these long wires into a number of shorter segments by adding gain stages between the segments. These stages are called repeaters.

#### 2.3.3 Repeaters

Because the delay of an uninterrupted wire grows quadratically with wire length, designers can add repeating elements periodically along the wire. When added in a way to optimize delay, repeaters make the total wire delay equal to the geometric mean of the total wire delay and the individual repeater stage delay. Hence, the length-squared term in wire delay falls out of the square root, making total delay linear with total wire length:

$$Delay_{rpt} \propto \sqrt{\frac{1}{2}R_{wire}C_{wire} \cdot FO4}$$
 (2.10)

Using repeaters is far more attractive for long wires, although they add some design complexity. First, the simplest repeaters are inverting elements, so an even number of repeaters is necessary to maintain logic levels<sup>7</sup>. Second, repeaters for global wires require many via cuts from the upper-layer wires all the way down to the substrate, potentially congesting routes on intervening layers. Third, designers are rarely afforded the luxury of placing repeaters in their optimal locations, because they require active area; designers usually floorplan repeaters in pre-planned clusters. Finally, even with delay-power optimizations, repeaters are still large devices, and repeating an entire bus takes an impressive amount of silicon area. Fortunately for these last two complications, delay and capacitance curves for repeater insertion have fairly shallow optimizations, so that adding or removing a single repeater stage, moving repeaters back and forth, or resizing repeaters have fairly small costs.

Repeated wires offer substantially increased performance. After sending one signal down a wire, we need wait only until that signal fully transitions on the first repeater segment before we send the next signal; the bandwidth of a repeated wire does not depend on the entire wire length. Also, increasing wire pitch makes the repeated segment length longer but does not change the segment delay, so wider wires simply reduce the number of available routing tracks and hence do not improve bandwidth.

In Chapter 4 we will examine repeater sizing, placement, and bandwidth more closely. For now we merely note that repeaters offer an alternative wire design structure that is far more attractive than uninterrupted wires for long wire lengths. For either unrepeated or repeated wires, simple geometric models for wire resistance and capacitance, when coupled with gate delay lead directly to useful wire delay metrics. Next, we consider how these metrics will scale with technology, beginning with gates.

## 2.4 Gate metrics under scaling

Historically, gates have scaled linearly with technology, and a useful model of FO4 delays has been  $500 \cdot L_{gate}$  ps under worst-case environmental conditions (typical devices, low  $V_{dd}$ , high temperature). In this expression,  $L_{gate}$  is in microns. Figure 2.11 shows FO4 delays for a number of different process technologies running at the worse-case environment corner.

<sup>&</sup>lt;sup>7</sup>Designers may opt to use buffered repeaters, which are two back-to-back inverters. Buffers avoid logical inversion complexity, but, as we will see in the next chapter, are slightly less efficient.

#### 2.4. GATE METRICS UNDER SCALING

This trend may continue for future generations of transistors, as devices seem scalable down to drawn dimensions of 0.018  $\mu$ m [29]. Whether or not such devices will continue to obey the above delay model is uncertain, due to issues in scaling gate oxide,  $V_{dd}$  and  $V_{th}$ . These concerns mean 500· $L_{gate}$  ps is a lower limit for future FO4 delays. Because we are considering wire delays relative to gate delays, faster gates provide the worst case for wire issues, and thus we will use this model as our gate delay projection.

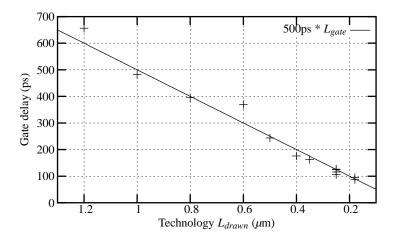


Figure 2.11: FO4 scaling at TTLH (90%  $V_{dd}$ , 125 degrees).

In a commodity 0.18- $\mu$ m technology, our model accurately predicts *FO4* to be about 90ps. In advanced in-house foundries, such as Intel's, *FO4*s are considerably faster: due to notches in the poly gates, the physical gate length  $L_{gate}$  ends up quite a bit smaller than the drawn (and advertised) feature size.<sup>8</sup> For example, Intel's 0.18- $\mu$ m gates have a physical gate length of slightly under 100nm [20], and Intel's 0.13- $\mu$ m gates have a physical gate length of 65nm [4]. Our model still holds as long as we use the correct physical gate length, so an *FO4* in the Intel 0.18- $\mu$ m process is about 50ps. This gap between commodity and advanced transistor technologies appears to be a fairly constant 2x multiplier, allowing us to continue to use our scaling formula.

We assume that other device parameters, such as gate and diffusion capacitance, will also continue to scale. Gate capacitance (per unit width) will remain around 1.5-2 fF/ $\mu$ m;

<sup>&</sup>lt;sup>8</sup>This is not the same as simply saying that "electrical gate length is shorter than drawn gate length," because it is a physical modification of the gate length.

although this would seem to demand too-thin gate oxides, high- $\kappa$  dielectrics may permit this aggressive scaling of the effective  $T_{ox}$  [30]. We project diffusion capacitance to stay at about half the gate capacitance for legged devices, although trench technologies and/or SOI can reduce this [31].

## 2.5 Wire characteristics under scaling

Before we look at how wire characteristics will scale, we will first examine the geometry assumptions in our baseline 0.18- $\mu$ m technology. This process has multiple layers of copper interconnect, with upper layers wider and taller than lower ones. The lowest metal layer, M1, has the finest pitch and hence the highest resistance, and it predominantly connects nets within gates or between relatively close gates. The middle layers, M2 through M4, have a wider pitch than M1 and connect both short- and long-haul routes, typically within functional units. The top layers, M5 and M6, have the widest pitch and hence the lowest resistance and they usually carry global routes, power and ground, and clock. Table 2.2 shows the pitches for these various layers in a contemporary 0.18- $\mu$ m technology. The pitches are described in technology-independent  $\lambda$ 's, where a  $\lambda$  is half of the drawn gate length. We will use similar wire pitches in our scaled technology projections: for our purposes, local wires have a pitch of  $5\lambda$ , semi-global wires a pitch of  $8\lambda$ , and global wires a pitch of  $16\lambda$ .

Metal	Pitch, µm	Pitch, $\lambda$
M6	1.76	20
M5	1.6	18
M4	1.08	12
M3	0.64	7
M2	0.64	7
M1	0.5	5.5

Table 2.2: Wire pitch dimensions for an Intel 0.18- $\mu$ m technology [36]

Predicting the future of wire technologies is tricky; whatever we say will almost certainly turn out to be wrong. Hence, we take a two-sided approach in this section. First, we consider wire performance given very optimistic, or aggressive, projections of technology scaling. This would include minor or insignificant resistance degradation from dishing or scattering, aggressive low- $\kappa$  dielectrics, and tall wire aspect ratios. Second, we also consider wire performance given very pessimistic, or conservative, projections. This would include significant scattering and dishing effects, very limited low- $\kappa$  dielectrics, and small wire aspect ratios. Pushing either projection to extremes allows us to confidently state that future technologies will fall inside the resulting broad range. This approach will be useful if both extremes still tell us a consistent story, and we shall see that they do.

		Technology, in $\mu$ m									
		0.18	0.13	0.1	0.07	0.05	0.035	0.025	0.018	0.013	
Common	Material	Cu	Cu	Cu	BulkCu	BulkCu	BulkCu	BulkCu	BulkCu	BulkCu	
Wire	ρ, mΩ · µm	0.022	0.022	0.022	0.018	0.018	0.018	0.018	0.018	0.018	
Properties	$\varepsilon_r$ for $C_c$	3.750	3.188	2.709	2.303	1.958	1.664	1.414	1.202	1.022	
	Pitch, µm	0.720	0.520	0.400	0.280	0.200	0.140	0.100	0.072	0.052	
	Aspect ratio	2.0	2.2	2.4	2.7	3.0	3.0	3.0	3.0	3.0	
Semi-	Thickness, µm	0.720	0.572	0.480	0.379	0.300	0.210	0.150	0.108	0.078	
Global	ILD, µm	0.750	0.540	0.480	0.405	0.315	0.210	0.150	0.108	0.078	
Wires	Barrier, µm	0.017	0.013	0.010	0	0	0	0	0	0	
	Dishing	0	0	0	0	0	0	0	0	0	
	$\alpha_{scatter}$	1	1	1	1	1	1	1	1	1	
	Pitch, $\mu m$	1.440	1.040	0.800	0.560	0.400	0.280	0.200	0.144	0.104	
	Aspect ratio	2.2	2.5	2.7	2.8	3.0	3.0	3.0	3.0	3.0	
	Thickness, µm	1.584	1.300	1.080	0.784	0.600	0.420	0.300	0.216	0.156	
Global	ILD, µm	1.5	1.08	0.96	0.81	0.63	0.420	0.300	0.216	0.156	
Wires	Barrier, µm	0.017	0.013	0.010	0	0	0	0	0	0	
	Dishing	0	0	0	0	0	0	0	0	0	
	$\alpha_{scatter}$	1	1	1	1	1	1	1	1	1	

Table 2.3: Aggressive wire predictions, technology terms only

Tables 2.3 and 2.4 summarize the technological parameters used in the conservative and aggressive technology scaling projections. These assume room temperature sheet resistances. Happily, recent SIA roadmap trends do fall within these bounds [37]. In both sets of scaling projections, we keep the semi-global pitch to be  $8\lambda$  and the global pitch to be  $16\lambda$ . One design possibility not pursued here is the use of "superwires": for performance and power delivery reasons, designers may choose to give the very top layers of metal a thickness and pitch that stays constant in microns. These global wires thus scale upwards in size relative to the rest of the metal layers, and will have superior current-carrying and delay characteristics, enabling global delays to scale with gate delays. These superwires

		Technology, in $\mu m$								
		0.18	0.13	0.1	0.07	0.05	0.035	0.025	0.018	0.013
Common	Material	Cu	Cu	Cu	Cu	Cu	Cu	Cu	Cu	Cu
Wire	ρ, mΩ · <i>μ</i> m	0.022	0.022	0.022	0.022	0.022	0.022	0.022	0.022	0.022
Properties	$\varepsilon_r$ for $C_c$	3.750	3.375	3.038	2.734	2.460	2.214	2.104	1.998	1.899
	Pitch, µm	0.720	0.520	0.400	0.280	0.200	0.140	0.100	0.072	0.052
	Aspect ratio	2.0	2.0	2.0	2.0	2.0	2.0	2.0	2.0	2.0
Semi-	Thickness, µm	0.720	0.520	0.400	0.280	0.200	0.140	0.100	0.072	0.052
Global	ILD, µm	0.750	0.540	0.480	0.405	0.315	0.210	0.150	0.108	0.078
Wires	Barrier, µm	0.017	0.012	0.008	0.006	0.004	0.003	0.003	0.002	0.002
	Dishing	0	0	0	0	0	0	0	0	0
	$\alpha_{scatter}$	1	1	1	1	1	1	1.05	1.05	1.05
	Pitch,µm	1.440	1.040	0.800	0.560	0.400	0.280	0.200	0.144	0.104
	Aspect ratio	2.2	2.2	2.2	2.2	2.2	2.2	2.2	2.2	2.2
	Thickness, µm	1.584	1.144	0.880	0.616	0.440	0.308	0.220	0.158	0.114
Global	ILD, µm	1.980	1.430	1.100	0.770	0.550	0.385	0.275	0.198	0.143
Wires	Barrier, µm	0.017	0.012	0.008	0.006	0.004	0.003	0.003	0.002	0.002
	Dishing	10%	10%	10%	10%	10%	10%	10%	10%	10%
	$\alpha_{scatter}$	1	1	1	1	1	1	1.05	1.05	1.05

Table 2.4: Conservative wire predictions, technology terms only

were first envisioned by Song and Glasser [38] for electromigration and voltage drop considerations. Because their consumption of wire resources (per gate) grows worse under scaling, our discussion does not include their usage.

## 2.5.1 Resistance under scaling

Under ideal scaling in all dimensions, wire resistance grows rapidly, as wire cross-sections fall by 2x each generation (only approximately, because of cladding and dishing correction terms). To prevent this resistance penalty, designers have been scaling wires in only the lateral dimension and increasing wire aspect ratios. Doing this makes the scaling penalty for resistance ( $\Omega$  per unit length) a single scale factor. Not considered here is the possibility of active cooling: refrigeration can lower copper resistance by almost an order of magnitude as temperatures drop from  $300^{\circ}K$  to  $77^{\circ}K$ , although today such cooling is prohibitively expensive. Resistance scaling is shown in Table 2.5 and in Figure 2.12.

### 2.5.2 Capacitance and inductance under scaling

Under ideal scaling, per-unit-length capacitance would fall due to incremental improvements in dielectric constants; it depends principally on the ratio of dimensions, so as long

		Technology, in $\mu$ m								
	0.18	0.13	0.1	0.07	0.05	0.035	0.025	0.018	0.013	
Aggr. semi-global wires	0.096	0.168	0.260	0.340	0.600	1.224	2.400	4.630	8.876	
Aggr. global wires	0.020	0.035	0.054	0.082	0.150	0.306	0.600	1.157	2.219	
Cons. semi-global wires	0.096	0.184	0.307	0.627	1.220	2.509	5.014	9.821	19.458	
Cons. global wires	0.023	0.044	0.073	0.150	0.292	0.598	1.183	2.298	4.474	

Table 2.5: Resistance, in  $\Omega/\mu m$ , with technology scaling

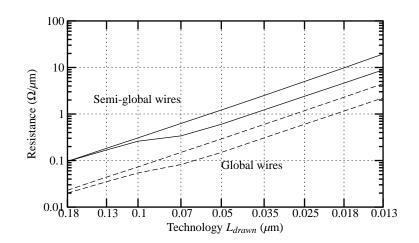


Figure 2.12: Resistance scaling, optimisitic and pessimistic trend curves

as all the dimensions are equally scaled, their ratios will not change. However, as mentioned above, wires are not ideally scaled, but typically shrunken only laterally. In this case, sidewall capacitance increases in a manner offset by advances in low- $\kappa$  dielectrics, meaning that capacitance changes only very slowly over technologies. See Table 2.6 and Figure 2.13. The numbers reflect worst-case capacitance: the side-to-side capacitances are "Miller multiplied" by a factor of two to model simultaneous switching of neighboring wires.

		Technology, in $\mu$ m								
	0.18	0.13	0.1	0.07	0.05	0.035	0.025	0.018	0.013	
Aggr. semi-global wires	0.414	0.397	0.374	0.359	0.345	0.315	0.288	0.266	0.247	
Aggr. global wires	0.440	0.430	0.403	0.367	0.345	0.315	0.288	0.266	0.247	
Cons. semi-global wires	0.414	0.387	0.359	0.333	0.311	0.295	0.287	0.280	0.272	
Cons. global wires	0.432	0.403	0.377	0.353	0.332	0.313	0.304	0.296	0.288	

Table 2.6: Capacitance, in  $fF/\mu m$ , with technology scaling

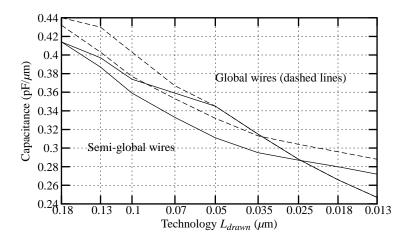


Figure 2.13: Capacitance scaling, optimisitic and pessimistic trend curves

Like capacitance, inductance per length should be roughly constant with scaling [39]. In fact, the rising aspect ratios of the wires will cause the value to slightly decrease. More important than the wire aspect ratio is how the power and ground networks scale, because current returns limit the inductive coupling of the wires. While the design of the supply is chip-dependent, the trend is for denser power distribution networks to lower the supply

impedance for each technology shrink [40], and about a 2x reduction in supply impedance is needed in each generation to maintain the same relative amount of supply noise [41]. These trends will prevent wire inductance from increasing under technology scaling, so we will continue to ignore inductance in these studies.

### 2.5.3 Noise as a limiter to scaling

A problem with the type of scaling discussed above, in which wires aspect ratios continually increase, is that with these stalagmite-like structures, noise coupling will quickly affect design. Today, wire aspect ratios are around 2, a value facilitated by the development of copper interconnect. When copper wires replaced aluminum wires on chips, their reduced resistivity allowed designers to cut down on wire aspect ratios in a one-shot reduction. This allowed the first generation of copper wires, when compared to their predecessors, to consume less power and have reduced noise coupling for the same overall resistivity. But as technologies scale, wire resistance will continue to increase at the same rate, leading again to higher and higher aspect ratios.

At today's aspect ratios of 2, the ratio of sidewall capacitance to total capacitance is between 65%-75%, depending on the dielectric constants between wires and between wire layers. In the SIA projections of wire technology, this aspect ratio has been capped at 2.2-2.5, to limit noise sensitivity. This limit is somewhat arbitrary: at today's ratio of 2, the noise problem is bad enough that designers already apply various techniques to overcome it. Upping the aspect ratio limit to 3 (for example) or beyond would not fundamentally alter the noise problem. We will next discuss a number of circuit techniques that help designers limit noise sensitivity of their long wires.

## 2.5.4 Noise minimization techniques

First, simply widening the wire-to-wire spacing for noise-critical nets decreases the coupling ratio, and can be effective when applied selectively. Moving the repeaters in a wide bus around such that each bit's repeaters are staggered from its neighbors forces capacitive noise to cancel itself down any potential victim wire. Because half of the injected noise must propagate down the *RC* wire to negate the other half, this cancellation is not perfect, but still effective. Charge compensation techniques, in which an explicit coupling capacitor injects reverse noise to combat parasitic coupling [26], requires some device area and extra power but can minimize noise as well as reduce data-dependent delay variation. These techniques are illustrated in Figures 2.14 and 2.15.

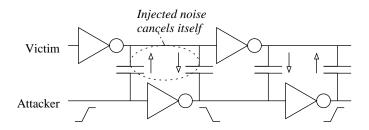


Figure 2.14: Staggering repeaters minimizes injected noise

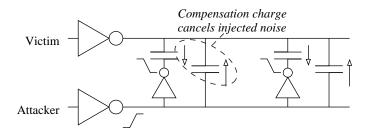


Figure 2.15: Charge compensation devices inject "negative" noise

The most effective method of reducing capacitive coupling noise runs each bit on two wires differentially, so the output comes from the voltage difference between the wires. Alternating wires are also twisted periodically. In this scheme, injected noise affects both wires equally, so that while the common-mode voltage might vary, the differential voltage is largely unchanged. In addition, these systems have minimal inductive coupling to the rest of the system, because each wire acts as the other's return path, creating the smallest possible current loops. While this requires some design effort to create differential receivers and area overhead for the drivers, the wire costs themselves decrease over time, as wire layers increase and the number of wire tracks per gate grows. With differential and twisted bits, wires can easily reject noise even if the coupling ratio approaches 90 to 100%. Figure 2.16 illustrates this approach.

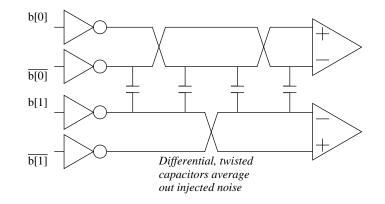


Figure 2.16: Twisted differential wires effectively eliminate noise

With a combination of these noise minimization techniques, designers can build robust wire systems in the presence of capacitive noise even larger than that suggested by today's aspect ratio limits. This analysis shows that fabrication and manufacturing restrictions should limit wire aspect ratio, not circuit noise considerations. In our "aggressive" predictions of technology scaling, we set a maximum aspect ratio of 3 to account for the limited efficacy of purely anisotropic etching.

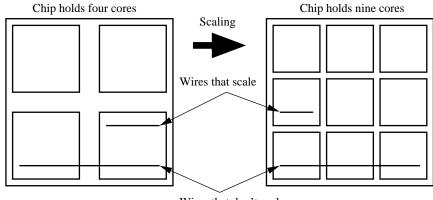
With or without use of these noise minimization techniques, we can consider the scaling of wire noise (both capacitive and inductive) with technology. Noise coupling should be mostly unchanged under scaling as long as the wires scale in length, and once the aspect ratio is "pegged" to its maximum value. In these cases, the overall capacitive coupling noise depends on the scaling of the ratio of the wire resistance to the driver resistance. If the wire lengths scale, the wire resistance scales down slowly, if at all. Driver resistance remains constant with scaling, thus leading to a coupling noise relative to the power supply (and hence to gate noise margins) which is either constant or slowly scaling down.

Inductive noise, depending as it does on a superposition of  $M \cdot di/dt$  terms, stays constant relative to the power supply for scaled-length wires. This is because the mutual inductance, M, stays constant per unit length, so that the total mutual inductance scales down with shorter wires. The total capacitance, and thus total current, also scales down, so the di/dt term stays constant. Thus the product of the two scales downward, along with the power supply  $V_{dd}$ .

If the wire lengths remain constant, the increase in wire resistance will cause the capacitive coupling noise to increase slightly and the inductive coupling noise to grow relative to the power supply. This increase in noise for long wires is another reason to use repeaters.

## 2.6 Wire performance under scaling

To properly discuss wire delays under technology scaling, we will first make a distinction between two kinds of wires, as shown in Figure 2.17. The first kind of wire connects gates locally within blocks. We can think of these as wires that span a block of fixed complexity; for example, a wire that runs across a block of 1000 gates. As technology scales, these gates get smaller, and this wire therefore gets shorter. This is the salient characteristic of these "local wires"—they get shorter under scaling.



Wires that don't scale

Figure 2.17: Two kinds of wire on a chip: local and global

The second kind of wire connects together sections of a die, and spans a block of fixed physical distance. For example, we can consider a wire that spans a 5mm length (approximately a quarter of a maximal die) regardless of technology. Under scaling, die size does not typically decrease—we simply add more functionality to the chip, so these "global wires" have the characteristic of remaining constant in length.

Note that this distinction between wire types does not assume specific wire lengths. There can be short global wires that span 1/1000 of a die, for example, and there can be long local wires that span 20,000 gates. The difference is simply in how they behave under technology scaling.

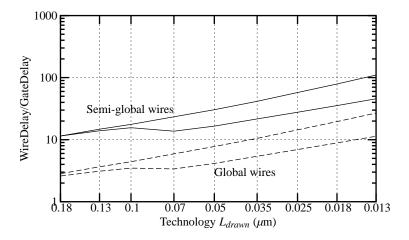


Figure 2.18: Local unrepeated wires, spanning 800 standard cells: 10x penalty over nine generations. The lines indicate optimistic and pessimistic scaling trends.

Here we show unrepeated wire delays for both local (Figure 2.18) and global (Figure 2.19) wires. In these graphs, local wires are defined as wires that span 800 standard cells. Global wires span 5mm. Along the x-axis are process technologies, plotted on log-scale to make them linear in time, and on the y-axis are wire delays, normalized to *FO4* gate delays. In both graphs, as we did for the resistance and capacitance scaling above, we show two pairs of trends, for mid-layer and upper-layer metals. Also as before, each trend has an upper and lower curve, representing the range of conservative (top curve) and aggressive (bottom curve) predictions of technology. Over this range of technology assumptions the conclusions are still consistent.

Because wire aspect ratios cap at either 2.2 (conservative) or 3 (aggressive), resistance grows quickly under scaling, leading to wire delays that increasingly lag gate delays. Normalized to gate delays, local wire delays degrade over nine generations by about 10x. Global wire delays are even worse, degrading over the same time span by more than 2000x.

Although these trends of wire performance paint a dismal picture of scaled designs, designers can use repeaters to improve wires. When applied to global wires, repeaters require a proportionally small area and design overhead and make these fixed-length wire

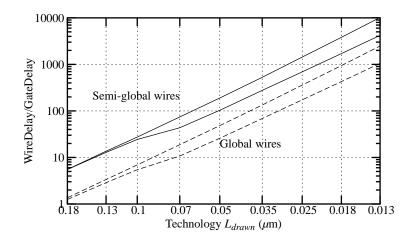


Figure 2.19: Global unrepeated wires, spanning 5mm: 2000x penalty over nine generations. The lines indicate optimistic and pessimistic scaling trends.

delays, relative to gates, grow only by a factor of 40x over nine generations (see Figure 2.20). Over nine generations, the number of gates covered by a fixed-length wire grows by  $0.7^{-9}$ , or 25x, so a repeated global delay penalty of 40x implies that communication costs versus logical span grows fairly slowly.

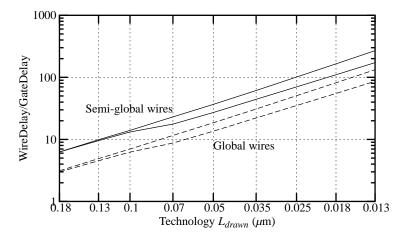


Figure 2.20: Global repeated wires, spanning 5mm: 40x penalty over nine generations

For local wires, a 10x penalty over nine generations is not crippling, because it means that the length of a wire *whose delay would otherwise be constant with gates* shrinks by  $\sqrt{10^{\frac{1}{9}}}$ , or about 14% per generation. This is not a lot. However, using repeaters can

eliminate this penalty for local wires as well. Of course, repeating all local wires requires an unwieldy number of repeaters, but designers can leverage the fact that the vast majority of local wires are also short: the distribution of wires in an 85,000-gate functional block extrapolated onto a 0.18- $\mu$ m process shows that less than 10% of the wires have delays that exceed half of a single gate delay [42]. Thus designers can ignore almost all of the local wires, and of the rest, only a few will be long enough to need repeaters; indeed, existing designs may already use repeaters for such wires. While under scaling the optimal repeater distance shrinks about 10% per generation, repeater delay is sufficiently insensitive to placement and sizing that this effect is minor.

With repeaters, local wires degrade relative to gates by a factor of only 2 or 3 (see Figure 2.21). Thus, the repeated propagation delay (in ps/mm) is basically constant under technology scaling, a fact that often surprises designers. This highlights the notion that wires themselves are not degrading in performance as much as chip complexity and performance are outpacing what wires can offer.

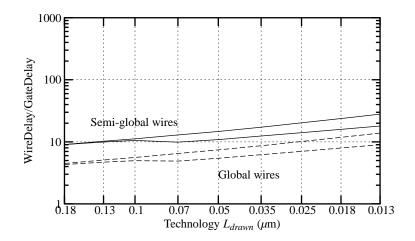


Figure 2.21: Local repeated wires, spanning 800 standard cells: 2-3x penalty over nine generations

## 2.7 Summary

This chapter examined how wire performance will scale as technologies advance. We first discussed a gate delay model, because wire delays are important only if they change relative to gate delays. We used the delay of a fanout-of-four inverter and modeled its speedup over time as linear with technology. Next, we discussed the two wire characteristics that are important to delay, resistance and capacitance, and why inductance is not important. Using very simple yet sufficiently sophisticated geometric models for resistance and capacitance, we constructed delay metrics and saw how they scaled with technology. On one hand, scaled-length, or "local," wires keep up with gate delays when repeated; on the other hand, fixed-length, or "global," wires cannot. Local wires get worse relative to gates by under 2-3X over nine generations of technology scaling, while global wires degrade by 40-50X. This bifurcated view of wires leads to a number of broad implications across VLSI design, and we shall consider a few of these effects in the next chapter.

# **Chapter 3**

# **Implications of Scaling**

The previous chapter painted a future with ever-increasing gate performance, local wires that can keep pace with these gates, and global wires that cannot. These scaling trends have implications for all aspects of VLSI engineering. In this chapter we focus on how scaling will affect design in two areas in particular. First, we will consider design "at the ground level," and examine scaling implications for gate-level designers. The vast majority of design today happens at the gate level, with place-and-route tools automatically creating physical instances of gates and wires. We will show that these computer-aided design (CAD) tools will need to improve as wires scale, despite oft-cited arguments to the contrary.

Second, we will consider design "at the 50,000-foot level," and examine scaling implications for computer architecture. Modern processors follow the dominant model of single-chip, monolithic architectures—a computing paradigm ultimately unsustainable under wire scaling. We will argue for architectures that are more "wire-aware," and advance modular systems as one such candidate.

One interesting offshoot of these modular architectures is that global communication becomes mediated by a very dense, regular, and structured mesh of global wires. The bad news with these global wires is that they can potentially consume hundreds of watts of power. The good news is that their highly ordered structure lends itself nicely to aggressive circuit design techniques designed to save wire power. We will briefly introduce the lowpower implications of wire scaling towards the end of the chapter to motivate the efficient circuits discussed in later chapters.

## **3.1** Design at the ground level: CAD tools

VLSI designers build chips today using one of two principal methodologies that differ in their extent of manual intervention. On one hand, microprocessors are designed by large design teams who hand-craft custom circuits for optimal performance and manually draw layout for high density. These designs that result have an extravagent productivity cost but demonstrate good performance; although they take more than three years to develop, leading-edge CPUs can run at close to 15 *FO4* delays per cycle, or 2.5 to 3.0GHz in a highly-optimized 0.13- $\mu$ m process.

On the other hand, application-specific integrated chips (ASICs) that target smaller markets are built by small design teams, who rely heavily on computer-aided design (CAD) tools to do much of the physical work for them. These highly automated CAD flows, used for designs such as graphics or network processors, generate structured netlists, place layouts from standard cell libraries, and route them together. The resulting designs can be completed very quickly—the Nvidia GeForce 4 graphics processor, with over 60 million transistors, went from initial concept to tapeout in just nine months—but do not clock as fast as the underlying technology permits; that same GeForce design runs at only about 40 *FO4*s per cycle [47].

A third design methodology, using gate arrays to compile chips, was historically important but is rapidly disappearing. While ASIC starts are predicted to be around 41,000 by 2005, gate-array starts are predicted to number only about 600 in the same year [48].

The ASIC design methodology, dominant outside of the few large CPU houses, is only as successful as its place-and-route CAD tools. Notably, interconnect scaling may force fundamental changes in how these tools are used: as wire delay grows, timing convergence will become difficult or impossible to attain with current fanout-based wire load models. In this section we explore the effects of wire scaling on CAD tools.

## **3.1.1 Claim: CAD tools need not improve**

In a now-famous 1998 tutorial at ICCAD, Sylvester and Keutzer analyzed interconnect scaling and its potential effects on CAD tools [49]. By examining the scaling of average-length wires, they concluded that delays for these wires will slowly decrease relative to

gate delays, and that today's CAD tools are therefore sufficient for future module-level designs. In the context of the wire scaling discussion from the previous chapter, they were examining short scaled ("local") wires.

We have two minor quibbles with their study. First, they conclude that average-length wires will not present a problem to CAD tools. But average-length wires have *never* presented a problem; *long* wires are the ones that constrain performance. Average wire length grows only weakly with gate count; a Donath model [43] with various Rent exponents (see Figure 3.1) shows that at an exponent of 0.7, average wire length increases less than 1.7x when gate count grows by an order of magnitude. At the same time, the semi-perimeter (half of the perimeter) of such a block, which models the longest interconnects, will grow by  $\sqrt{10} \approx 3.2x$ . Thus, modeling wire performance using optimistically short wires will return optimistic results. Second, their results depended highly on a number of scaling assumptions; under slight variations of the scaling of wire capacitance relative to gate and diffusion capacitances, their relative delays would actually *increase* slightly.

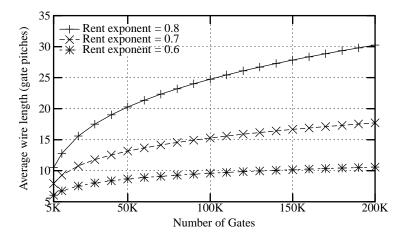


Figure 3.1: Average wire length versus gate count

These complaints are trivial. As seen in the previous chapter, even long scaled wires track gate performance fairly well, and so the idea underlying their key result still holds: that the delay of local scaled wires, relative to gates, changes little, if at all, and that wire delays for a module that scales in size do not get much worse (or much better) as technology progresses.

However, while the underlying idea is completely accurate, using it to conclude that CAD tools need not improve is specious. To understand why, we will discuss the fundamental weakness of synthesis CAD tools: discrepancies between pre-layout and post-layout wire loads.

#### **3.1.2 Underlying problem in synthesis**

The CAD methodology prevalent at the turn of the century<sup>1</sup> did not guarantee timing convergence to a design because of a discrepancy between pre-layout wire load estimates and the actual post-layout wire loads. Initial steps in synthesis decide logical structures using fanout-based wire load models, typically supplied by the standard-cell library vendor and based on statistical analyses of past designs using that library. Such a model associates with each gate fanout a single sample wire load from its overall distribution; for example, a gate that drives three other gates is assumed to see a representative capacitive load from all three-gate-driving cells in that library's history. The problem is that actual (post-layout) wire loading has a Poisson distribution with a narrow peak around the mean wire load and a long tail to the right. Figure 3.2 shows the discrepancy between post-layout and statistical wire loads for a small design, where the nets are sorted by fanout and then by post-layout wire load [50].

As the figure indicates, synthesis flows may reasonably estimate the wire load of short or average-length wires, but dramatically underestimate the load of the longest nets. For these long nets, the CAD flow may pick the wrong logical structure for the net and not account for the wire's intrinsic delay, leading to max-delay problems, or speed paths. These critical delays would in turn require incremental optimizations, such as gate sizing, buffer insertion, and critical path resynthesis. These resulting changes need to be merged back into the design without perturbing the unchanged logic, or else the incremental optimizations would themselves be inaccurate and lead to a new set of critical paths.

Compounding the problem is that many CAD flows do not account for, or poorly estimate, wire resistance. Wire resistance makes wire delays quadratic in length, so these unexpected extra delays easily create a new set of critical paths.

<sup>&</sup>lt;sup>1</sup>Several recent efforts, both in startups as well as in established CAD companies, are now working to improve this flow.

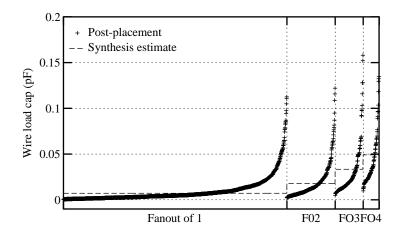


Figure 3.2: Estimated versus actual wire loads

Advances in synthesis placement algorithms to manage wire delays, such as by netlength constrained placement [51], attempt to attack this CAD-tool wire delay problem with varying degrees of success. Wires that nonetheless still fail to converge are called "wire exceptions" because they require designers to manually intervene in the synthesis process. Designer effort required per exception depends on specific CAD tools and methodologies but will not scale down with technologies, so to properly consider the effects of wire scaling on CAD tools, we must examine how the number of these exceptions will scale.

#### **3.1.3** Wire exceptions

Because wire problems arise from wire loads exceeding a fanout-based load prediction, we will use long wires as a proxy for wire exceptions, and examine how long wires scale. If we define a "long" wire simply as one whose intrinsic delay is some fraction of an *FO4* delay, we can define a *critical length* to be this wire's length. In the following expression,  $R_{wire}$  and  $C_{wire}$  are both per-unit-length quantities, giving  $L_{crit}$  units of length.

$$L_{crit} = \sqrt{\frac{\alpha FO4}{R_{wire}C_{wire}}} \tag{3.1}$$

In this discussion we will choose the fraction  $\alpha$  to be half of an *FO4*, although the actual value is not central to the analysis. In a 0.18- $\mu$ m technology, using mid-layer metals

with minimum pitch, the  $L_{crit}$  is 1mm—a distance certainly longer than an average wire's length, but still present in modestly-sized modules. For example, even with an optimistic area utilization of 100% and a cell pitch of  $7\mu$ m, a 50K-gate block still has a semi-perimeter of more than 3mm.

As  $L_{crit}$  depends not only on wire layer but also on wire geometries such as width and spacing, various wire engineering techniques can be used to mitigate both wire intrinsic delay as well as wire loading. In this manner, the actual  $L_{crit}$  may be increased and long wire exceptions thus handled. However, while today's CAD tools attempt these wire optimizations by tapering wires, widening wires, or promoting wires to higher layers, they do so only imperfectly, leaving much wire cleanup work to designers still.

Under technology scaling, we can plot how  $L_{crit}$  trends relative to gate pitches, assuming a baseline of  $7\mu$ m gate pitches in a 0.18- $\mu$ m technology. Following the strategy of the previous chapter, we plot both aggressive and conservative projections of technology. The results are shown below in Table 3.1 and in Figure 3.3. We can see that  $L_{crit}$  decreases very slowly, changing only by 2-3x over 9 technology generations. The slight bump in the agressive trend comes from the (optimistic) use of low-resistive bulk copper interconnects.

	Technology, in $\mu$ m									
	0.18 0.13 0.1 0.07 0.05 0.035 0.025 0.018 0.013							0.013		
Aggressive scaling	152	138	130	139	126	111	98	86	76	
Conservative scaling	152	134	123	106	93	80	68	58	49	

Table 3.1: *L<sub>crit</sub>* in gate pitches, with technology scaling

To bolster this conclusion, we can apply a wire-length distribution model from Davis *et al.* based on Rent parameters and gate count [44]. To pick the parameters of this model (a factor k, an exponent p, and a fanout f) we applied this model to four synthesized blocks: three units from the M-Machine multiprocessor [45] and the Magic controller chip from the Flash multiprocessor [46]. From Figures 3.4 and 3.5 (we show only one M-Machine plot for brevity) we see that the model fits wire length distributions well for a wide variety of gate counts.

Using these same model parameters, we can show the percentage of wires that exceed

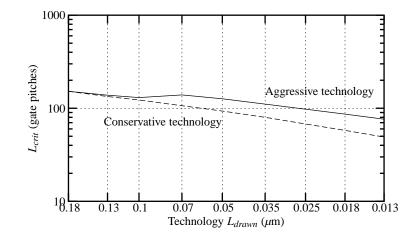


Figure 3.3: *L<sub>crit</sub>* scaling

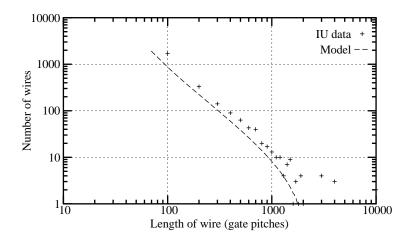


Figure 3.4: M-Machine IU unit versus model (k=3, p=0.7, f=2.3)

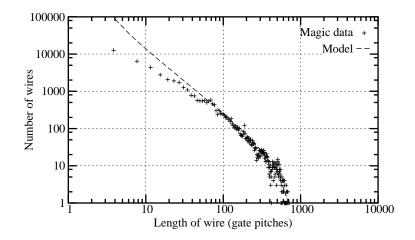


Figure 3.5: Flash MAGIC node versus model (k=3, p=0.7, f=2.5)

 $L_{crit}$  for a number of gate design sizes and technologies, as shown below in Figure 3.6. In this figure we again use both conservative and aggressive technology scaling assumptions, and we also vary the design size by the number of gates.

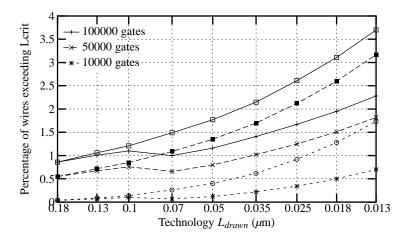


Figure 3.6: Percentage of wires that exceed  $L_{crit}$  under scaling. The two sets of curves show optimistic and pessimistic technology scaling trends.

Figure 3.6 shows that the number of wires that exceed  $L_{crit}$  in each technology will grow somewhat slowly. In the most pessimistic case, the percentage of wires exceeding  $L_{crit}$  grows only by about 4x over nine generations. We can therefore argue that the number of CAD tool exceptions will also grow slowly and keep module-level design manageable. Yet we cannot declare victory over wires and CAD tools. The reason is an exponential growth in die complexity.

Scaling effects design in two ways: first, as we have just discussed, it increases the percentage of wires that exceed  $L_{crit}$ , albeit only slightly. Second, and more importantly, it exponentially increases the number of wires in a design. Over five generations of scaling, for example, the percentage of wires that exceed  $L_{crit}$  only doubles, but the total number of wires grows by  $2^5 = 32$ . Figure 3.7 below illustrates what might happen in one generation of scaling. Assume that nine modules have eighteen wire exceptions. In the next generation we can squeeze twice as many modules on the same die. Those eighteen modules, at the same wire exception *rate*, will generate nearly forty wire exceptions<sup>2</sup>.

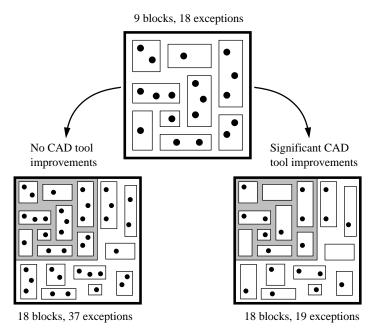


Figure 3.7: Wire exceptions under scaling

Fixing this problem requires either larger design teams, longer design times, or improved CAD tools that result in a *lower wire exception rate*. In Figure 3.7 we show an ideal result: CAD tool improvements that cut the wire exception rate in half, resulting in constant designer effort. To achieve such a result, however, requires CAD tools that can handle

<sup>&</sup>lt;sup>2</sup>These numbers, of course, are purely for illustration. Actual wire exceptions for large blocks can number in the hundreds or thousands.

increasing  $L_{crit}$ . Figure 3.8 below shows how  $L_{crit}$  must dramatically increase in order to maintain constant designer effort, and how CAD tools must therefore comprehend the resulting increasingly longer wires. This figure shows how making a CAD tool handle just 1% more wires in a design can easily double the maximum wire length with which the tool must deal. A number of efforts today attempt to address this long-wire problem, including merged synthesis and layout and ECO flows like buffer insertion and gate relocation [51].

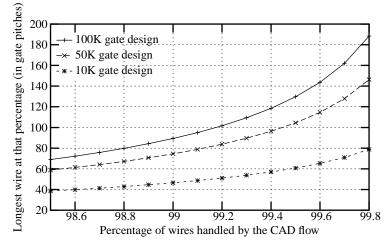


Figure 3.8: CAD tools handling more wires must also handle longer wires

Hence, claims to the contrary notwithstanding, CAD tools do have a wire problem. It arises from die complexity growth: absent tool improvements, the total number of wire exceptions will double each generation. We can double neither design team size nor design times, so future synthesis flows must handle an increasingly larger percentage of long wires without designer intervention. This is the key challenge confronting CAD tools.

## **3.2 Design at 50,000 feet: Architectures**

The story of wires told in the previous chapter, with a bifurcation between scaled wires that track gate performance and fixed-length wires that hopelessly fall behind gates, bodes poorly for today's computer architectures. Today's processor architectures follow a model that optimized computation at the cost of communication, but as global wires (the communication) become slower and slower, and transistors (the computation) become cheaper and cheaper, this architectural model will become ultimately unsustainable.

### **3.2.1** A historical perspective

Modern processor architectures evolved when transistors were precious and on-chip communication was cheap. As a result, a theme of functional centralization and reuse prevailed, with an emphasis on monolithic global resources made available to multiple consumers [54][55]. For a long time, wire access to such resources was low-latency, making this model attractive to architects and programmers: it presented the most uniform computational framework for the best utilization of functional units. This focus on function rather than communication pervades all aspects of design. Consider hardware description languages like Verilog, for instance: in Verilog code, each line represents a function, and communication between different functions, or lines of code, are implied only by variable names.

In older technology generations, few functional units actually fit on a die, so maximal use of these few functional units was paramount, and fitting all the needed functions on the chip was the critical design goal. The number of gates that fit on a  $2.0-\mu$ m technology chip was less than 50K, the size of today's synthesized blocks. Cross-chip wires were not a problem, because their logical span was so short, and wire resistance was not an issue.

As technologies improved, wire resistance remained small, but the increasing capacitance of long wires became significant. Floorplanning of high-performance designs became an important design step, so that proper device sizing could keep communication costs low. Continued technology scaling led to the situation where global wire delays were non-trivial but still much less than a clock cycle. In this design period, the programming model remained one of globally shared resources, but with micro-architectures increasingly partitioned. For example, the instruction fetch unit, while logically part of the datapath, physically migrated to the cache to minimize branch latencies. The address adder in many machines was duplicated: one in the datapath where it logically belonged, and a smaller version near the data cache to generate the cache index, again to reduce latency. Wires delays were still modest (much less than a cycle), and these micro-architecture changes were mostly invisible to the user. Designers developed many tools and methodologies to handle the increasing importance of wires during this period (the beginning of sub-micron design), including analytical wire *RC* models and more accurate wire simulation tools such as asymptotic waveform evaluation [56]; floorplanning techniques such as delay-driven segregation of local and global routing; and local circuit generation techniques such as layout-driven synthesis. A modern 0.18- $\mu$ m design utilizes some or all of the above techniques. While local routing within reasonably-sized blocks has negligible wire delays, global routes between such blocks are closer to half a cycle. The cost of communication is becoming more explicit. Chips are partitioned early in the design process, and the delays of global lines are rolled into timing models.

As the complexity of digital systems has continued to increase, architects have responded to higher communication costs by further partitioning internal micro-architectures and adding internal latency (pipe stages) in locations that will least damage machine performance. While some researchers imply that the delay of global wires sets cycle times [57][58], in high-performance machines this is not true; communication on these global wires is simply pipelined. This additional internal latency allows the machines to absorb the penalty of on-chip wires while still taking advantage of their high bandwidth. These added latencies are now visible to the user, but have small effects on the programming model. For example, in the Alpha 21264 processor, the integer unit is partitioned into two clusters, and the latency for communicating between these clusters takes an additional cycle [59]; and in the latest Itanium-class processor, a pair of pipeline stages were added to account for wire delays to get to and from the 6MB level-three cache [52].

These micro-architectural techniques can hide minor global wire delays, but not if these delays continue to expand as described in Chapter 2. Wires delays to, from, and over monolithic architectural units that grow to multiple clock cycles will degrade system performance. As a result, increasing system parallelism and complexity with global models will be increasingly difficult: if we try to exploit technology scaling by increasing the size of centralized resources, we will pay the damaging performance costs of long wires.

#### **3.2.2 On-die signal range**

The cost in clock cycles of long wire delays depends on how clock frequency will scale. To predict this we can examine historical trends of numbers of *FO4*s per cycle. Table 3.2 shows approximate *FO4*s per cycle for Intel x86 microprocessors. Current machines cycle between fifteen and twenty *FO4*s per clock. Extrapolating future clock cycle times from this table can lead to unrealistic predictions: a semi-log fit would lead us to expect clock cycles of 3-4 *FO4*s per clock within a few generations. Performance-oriented architectural studies [33][34] showed that machines spinning at 6-8 *FO4*s per clock are best, but even these fast-cycling machines seem difficult to build.

Micro- architecture	Year	Technology	<i>FO4/</i> cycle
386	1985	1 <i>µ</i> m	100
486	1989	1 <i>µ</i> m	60
Pentium	1993	0.6µm	40
Pentium Pro, PII, PIII	1995	0.4µm	25
Pentium 4	2000	0.18µm	16

Table 3.2: Approximate FO4 per cycle for x86 Intel architectures

The design difficulties for very short-tick machines lie in both generating as well as in using fast clocks. First, creating a clean clock that rises and falls in just a handful of *FO4s* and that does not look like a sine wave (thus suffering from high power supply sensitivity), is hard—sharp rise and fall times require very short delays from the clock drivers to the clock loads. Since the clock loads on a chip—all the latches and flops—are currently enormous and increasing, this implies a very complicated and potentially high-power clock grid design.

Second, clock elements, whether latches, pulsed latches [53], or flipflops, all have through-delays of around 1 *FO4* or more, and these delays typically appear twice per cycle. If cycle times decreased to 6 *FO4*s, or worse yet, to 4 *FO4*s, these latch delays will consume a significant fraction of each phase, leaving precious little time in which to do real work. Including the effects of clock skew and other clock uncertainties exacerbates this problem, although some techniques help to minimize these effects [32].

Although we expect very short-tick machines to be extremely and perhaps prohibitively expensive, the actual future cycle time trends are unclear, and exactly how fast to run a machine is still an open question. Recently, architectural studies of both energy and performance concluded that the best number of FO4s per cycle was around 18 [35]. In the following discussions, we will follow this lead and limit the number of FO4s per clock cycle to be at least  $16^3$ . If machines spin even faster, then the trends in Figure 3.9 will look even worse, with even more limited on-die signal ranges.

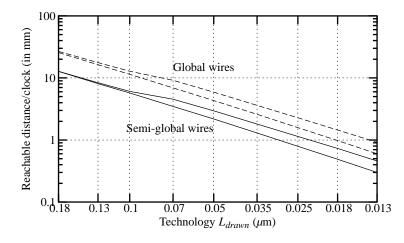


Figure 3.9: Reachable distance per 16 FO4 clock, in mm

Repeated wire delays determine the reachable on-die distance per clock, as shown in Figure 3.9. From an architectural perspective, the importance of this distance lies in implicit versus explicit latencies: spans that lie within this "signal range" need not be broken into pipestages or otherwise synchronized across cycles, while spans that cannot be crossed within a clock will have architecturally explicit latencies. While in Figure 3.9, we show this reachable span in microns, in Figure 3.10 we show it in  $\lambda$ 's. Notice that although the reachable span is decreasing in absolute distances, the logical span in  $\lambda$  is essentially constant over many technology generations, supporting the earlier conclusion that designs that shrink will have nicely-scaled performance.

Figure 3.11 maps these signal ranges to die sizes. The 2001-02 SIA roadmap projects

 $<sup>^{3}</sup>$ Note that the newest Intel micro-architecture, based on the 64-bit EPIC machines, cycle at nearly 30 *FO4*/cycle.

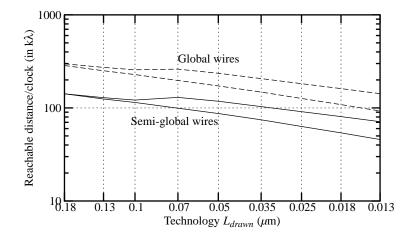


Figure 3.10: Reachable distance per 16 FO4 clock, in  $k\lambda$ 

maximal die sizes to be a constant 17.6mm on a side under scaling [37]. Seven technology generations appear in the figure; further projections would be hard to see. At each generation the shaded portion represents the reachable distance from the chip center, drawn as a diamond because on-chip routing is strictly Manhattan-style. At the 0.18- $\mu$ m technology generation, nearly the entire chip fits inside the reachable distance, but this quickly changes. To emphasize that this results almost entirely from growing chip complexity, we also draw how a 0.18- $\mu$ m chip would scale. The reachable distance falls only slightly faster than this scaled block of constant logical complexity.

As Figure 3.11 extends to future generations, we see that supporting global wire delays will become increasingly architecturally onerous and eventually infeasible. Solutions to this problem may rely upon alternate technologies, such as optical interconnects [19], to reduce global wire delays, although these communication methods have their own complications. On the other-hand, we might also consider architectures that leverage the delay characteristics of local and global wires in order to maintain high performance; we will discuss such "wire-aware" architectures next.

#### 3.2.3 Wire-aware architectures: modularity

An architecture that is "wire-aware" would maintain high performance under technology scaling by defining system speed with fast local wires and accepting the high latency of

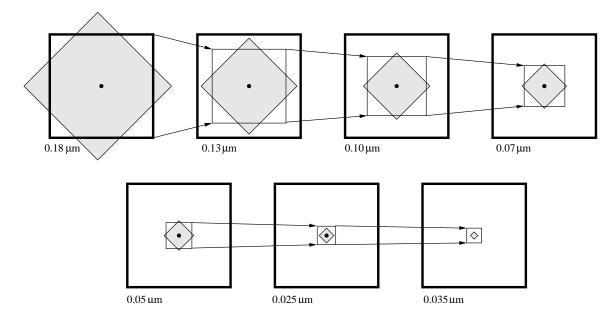


Figure 3.11: Reachable distance on a chip under scaling

global communication. In such a system, global cross-die communications do not limit frequency, and critical regions do not grow in complexity with scaling, so their wires remain local and scale down in length.

One way to do this is to exploit modularity: imagine a large die filled with a number of independent computing cores that use explicit task-level parallelism. Under technology scaling, we maintain the complexity of the cores, so that they shrink in dimensions. We can continue to improve overall system performance by simply adding more of these cores to the die with each generation. In such a system, the local cores' performance scales with technology because their scaled-length wires do not limit their clock frequencies. Core-to-core communication takes place over a "highway" of general and flexible wires, in which point-to-point virtual network connections are created and used for each packet. This communication between cores runs on similarly scaled-length ("local") wires as well, but global communication across the die requires an ever-increasing number of hops and thus an ever-increasing latency. However, these global communication costs are architecturally explicit and can be kept out of the critical paths of the chip.

A number of researchers have proposed such modular, wire-aware architectures. For

example, the RAW project at MIT divides a large chip into an array of identical tiles, distributing all computation and communication over a system of mesh networks. This structure allows scability in the face of wire delays, which would otherwise limit the operation of centralized resources [60]. Similarly, the recently-proposed Grid processor architecture project from the University of Texas puts forth an execution model of "chained" logic units, assembled in a systolic-like array, in order to expose wire-delay constraints at an architectural level [61].

Yet another example of a modular architecture is the Stanford Smart Memories project [62]. Targeted at a 0.10- $\mu$ m technology, it is organized as a collection of 64 compute tiles, organized into 16 quads of four tiles each (see Figure 3.12). Each tile, about 2.2mm on a side, has a reasonable amount of computational complexity, with two integer ALUs and one floating-point ALU; 128KB of reconfigurable on-die SRAM; and fast local interconnect. Each quad of four tiles also has a port to the flexible global routing network, organized as a mesh. The global routing network, the local tile SRAMs, and the local tile computation blocks are also dynamically configurable.

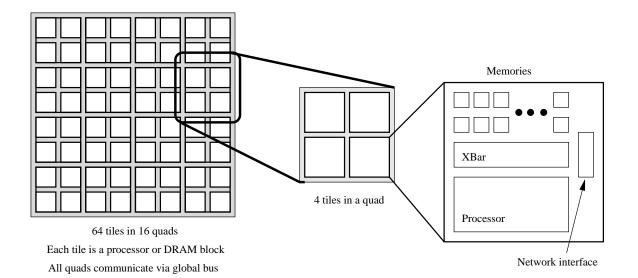


Figure 3.12: Smart Memories project

#### Benefits of a modular architecture

Modularity in architectures leads to a number of advantages. First and foremost, dividing a chip into local computation with global communication enforces a sense of wire-awareness. By elevating all cross-die communication to a mesh of global wires, in which request and reply packets get routed out and in with high latency, we make long wires architecturally visible in the programming model. Of course, on-chip wires are also cheap, and growing cheaper as advances in technology offer more wire layers, so on-chip wires still offer enormous bandwidth.

Second, modular architectures suggest a way to escape the spiraling costs of design and verification teams. A recent third-generation Itanium microprocessor with nearly half a billion transistors (25 million in the core and 385 million on the on-chip cache) consumed almost a thousand person-years in design and verification [52]. As we look towards chips with billions—or tens of billions—of transistors on them, unless productivity rates grow exponentially, we can forsee that design and verification costs will make large microprocessor design prohibitively expensive for all but a handful of companies. Using modularity to design a small block once and then replicating it over a die allows us to apply a "divide-and-conquer" tactic to design and verification.

Third, modularity invites flexibility. Utilizing all of the offered modules of an architecture becomes easier if the modules can be programmed to behave in different ways. In the case of the Smart Memories modular architecture, a great deal of flexibility arises from a concerted effort to make the machine dynamically reconfigurable: each CPU can be configured to act as a two-way superscalar machine, a VLIW machine, or an explicitly-controlled microcoded machine. In addition, the local SRAM blocks on each tile can be reconfigured to resemble regular random-access memories, caches, or stream buffers. Finally, the local and global interconnects can be reconfigured on the fly to dynamically route data within and between modules.

This reconfigurability and flexibility has an important cost benefit. Fabricating any chip requires the patterning of the lithographic masks used to "stencil" the chip substrate, and the cost of making these masks is growing rapidly: in two more technology generations,

a full mask set will cost upwards of ten million dollars [63]. This trend threatens "fineline" prototyping and raises doubts about the affordability of ASICs. However, flexible and dynamically reconfigurable architectures offer the possibility of using a single chip (and mask set) for solving a number of different problems.

Finally, modularity in architectures offers us the benefits of global wiring homogeneity. Because global core-to-core wiring consists of a regular and reconfigurable mesh network, these wires are highly structured and thus have very well-characterized layouts and environments. By constrast, monolithic architectures would have largely random global wiring. The advantages of this global wire regularity is that it offers a clean environment for circuit families that might otherwise be considered too risky or too design-intensive for random wiring, but that offer large performance, noise, or power gains. In the last section of this chapter we will consider one such circuit family, low-voltage-swing signaling, and why it might be useful.

## **3.3** Efficient global wiring networks

The scalable design cost of a very high-bandwidth, flexible global routing network is not area, because as wire layers increase with technology, the number of wire pitches per gate continues to grow. Rather, the cost of a global routing network lies in power consumption. The Smart Memories global network consists of a number of 128-bit buses, each spanning 2.5mm in length. These buses can support a peak aggregate bandwidth of over a terabyte per second, but in doing so would also consume nearly 80W of switching energy in the wires and their repeaters. Rarely do networks achieve their peak performance, but even at a more likely 25% or 30% utilization, the global network will consume from 20-25W. Under technology scaling, with chip voltages falling slower than feature sizes, this global network power will not scale down with technology, meaning that current delivery and thermal control will become increasingly difficult.<sup>4</sup>

A natural idea to reduce wire power is to reduce the voltage swing on the wires: this

<sup>&</sup>lt;sup>4</sup>Reduction of the global network power may not be the lowest-hanging fruit (the cores combine, on paper, to consume far more power than the global wires), but useful to explore nonetheless.

leads to a linear power reduction directly, and if the voltage supply driving the wires is lowered as well, to a quadratic reduction in consumed power. However, such low-voltage signaling is fraught with danger: induced noise and timing uncertainty make the performancerobustness tradeoffs for low-swing circuits difficult to navigate. In a later chapter we will discuss the principal issues involved with low-swing on-die wiring systems, but first we should ask if simpler schemes might be equally effective for much lower design costs, and in the next chapter we consider some of these simpler approaches.

## 3.4 Summary

In this chapter we considered some implications of wire scaling as they applied to two opposite corners of the VLSI design space. First, at the basic level of circuit and layout synthesis, we looked at how wire scaling will affect the use and effectiveness of CAD tools. Despite some claims to the contrary, wire scaling will require CAD tools to improve dramatically in their ability to route and converge on designs with long wires. This is not because modules are getting more and more complicated, but rather because dies are gathering more and more modules. Second, at the more global level of system architecture, we considered how an awareness of the disparity between scaled and fixed-length wires could be made explicit to the underlying machine by exploiting modularity. We briefly discussed three proposed machines: the MIT RAW project, the UT-Austin GPA machine, and the Stanford Smart Memories reconfigurable architecture. All are highly modular, and thus all provide the wire-awareness, the design productivity boost, and the potential cost benefit from the machine flexibility that arises with such modularity. Finally, we discussed how the global network of such modular machines is highly structured and characterized, making it a natural bed for otherwise risky circuits like low-swing interconnects. In the successive chapters we will consider both simple and complex methods of saving interconnect energy.

# **Chapter 4**

# **Efficient Repeaters**

At the end of the previous chapter we discussed the energy cost in switching a global network of flexible wires. While it is not the dominant consumer of energy in large chips, it is still significant. This chapter will examine some design tradeoffs that arise when we try to limit this wire energy. We will first briefly discuss coding methods to reduce energy usage, which can be implemented separately from any circuit optimizations. The bulk of the chapter will deal with circuit techniques: revisiting the well-known CMOS repeater sizing problem and solution and discussing the effects of including energy in the formulation. We will then extend this discussion to other types of repeaters. Because the energy gain from these types of optimizations are small at best, the next chapter will consider more the aggressive power-savings technique of low-voltage swing circuitry.

# 4.1 Coding for energy savings

By exploiting the property of CMOS circuits that energy is only dissipated when voltages change, designers can encode data transitions to minimize such events. On-chip wires only burn energy when bits flip, so in this encoding scheme, each *n*-bit data word is compared with the immediately previous *n*-bit data word. If more than half of the *n* bits need to flip from the previous word, then the inverse of the data word is transmitted instead. An extra single-bit control wire tells the receiver whether or not to flip the received data word [66].

To estimate the energy savings in this system, we can consider the number of wires

that flip between data words. Assuming a perfectly random data distribution, an *n*-bit bus carrying a particular data word will have  $2^n$  possible next data words with equal probability. One of those next-word possibilities will have no bit changes (*i.e.* the same word, repeated), and one will have all bit changes. For the rest, the number of next-word possibilities that have *k* bit changes equals the number of combinations of *n* items chosen *k* ways. On average, the total number of flipped bits is simply half the bus width:

Average bits flipped = 
$$\frac{1}{2^n} \sum_{i=0}^n \binom{n}{i} i = \frac{n}{2}$$
 (4.1)

If we now selectively invert the data word, some of these terms will have a lower "cost." For example, the next-word possibility that would normally flip all the bits will now flip none of the data bits, and just the control wire. We can express this as

Avg. bits flipped, inverted = 
$$\frac{1}{2^n} \left[ \sum_{i=0}^{\lfloor \frac{n}{2} \rfloor} {n \choose i} i + \sum_{i=\lfloor \frac{n}{2} \rfloor+1}^n {n \choose i} (n-i) \right] + \frac{1}{2}$$
(4.2)

where the last term of 1/2 represents the control wire flipping half of the time. We expect that the benefits to very small buses to be small, because the overhead of the extra control wire will be important. On the other hand, the benefits to wide buses should also be small, as the probability of hitting a "good" combination of words will be remote: most of the next word possibilities will be clustered towards n/2 bits flipping, where bus-invert helps, but only minimally. The energy savings, in terms of number of bits flipping, is shown in Table 4.1.

These results imply that bus-invert, done on buses in groups of eight bits, can be a reasonable power-savings scheme. Its costs include extra latency at the start of the transmission, for the driver to perform a bitwise XOR between the current and previous data words and then count the "1"'s that result, as well as extra latency at the end of the transmission, for the receiver to decide whether to invert the data word or not. For bandwidth-sensitive but latency-insensitive applications, these costs are not onerous, because we can hide the encoding and decoding costs in an extra clock period. The encoding and decoding logic also consume energy, to clock the registers that pipeline the data words, to do the XOR-ing

### 4.1. CODING FOR ENERGY SAVINGS

Bus width	Avg. # of bits flipped	Avg. # of bits flipped w/bus-invert	Savings
2	1	1.00	0.00%
4	2	1.75	12.50%
8	4	3.41	14.84%
16	8	6.93	13.39%
32	16	14.26	10.87%
64	32	29.32	8.37%
128	64	60.00	6.26%
256	128	122.12	4.59%
512	256	247.48	3.33%

Table 4.1: Savings from using bus-invert

and majority-counting (see Figure 4.1) on the transmission end, and to do the inverting and muxing on the receiver end. With sufficiently long interconnect, the wire capacitive load will dominate the switched capacitance in the logic, maintaining the energy savings. The majority counting amplifier can be implemented as a simple amplifier, with small sizes to minimize power consumption; the normally deleterious effects of transistor offsets only hurt by reducing energy savings.

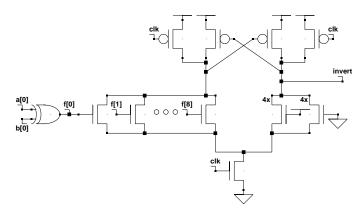


Figure 4.1: Bus-invert transmission logic

For dynamically signalled schemes, where wires toggle between a precharge and an evaluation voltage, buses can be encoded to minimize the number of bits that evaluate. For

example, codings that limit the number of "1"s can save power on a domino bus. Prior work focused on reducing chip I/O power looked at very similar schemes [67][68][69], because chip I/O's use resistive pullup loads and dissipate power on pull-down values. Those same techniques are applicable here as well.

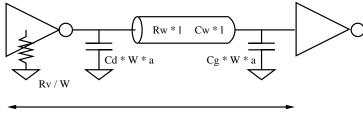
Side-to-side wire coupling dominates interconnect capacitance, so designers might be tempted to try encoding schemes that aim to minimize switched capacitance [70]. For example, we might map an *n*-bit bus onto *m* wires (m > n) such that wires with rising edges are not placed next to wires with falling edges (and *vice versa*). A simple example would use 2n wires to transmit *n* bits, weaving a power supply line between each actual bit. However, this and similar encoding schemes only reduce worst-case power (as well as worst-case noise and worst-case delay), and leave the average power dissipation unchanged; they also increase the minimum power dissipation. In this example, we would rather simply remove the interleaved power supply lines to significantly reduce both maximum and average power for the same area cost.

We might also try mapping *n* wires onto *m* wires (again, m > n) to save transitions between successive words. A trivial example would expand an 8-bit bus into 256 wires. On average, an 8-bit bus will have four bits transitioning between consecutive words, but the expanded 256 bits will be "1-hot," and so at the most have only two transitions between consecutive words. However, again by simply expanding our eight wires' spacing to consume more area (and not even 256 wires' worth) we can achieve much more energy savings.

Encoding for low-power can be seen as an optimization above, and orthogonal to, the physical layer of the actual interconnect circuitry. Therefore, we can assume that the actual bitstream to be transmitted down a wire has already been optimally encoded and concern ourselves with the task of physically signaling the bits down repeated wires. As we shall see in the next chapter, by reducing the signal swing, we can make the energy consumed in the wires a small percentage of the total energy required to transmit a token. Because coding schemes increase driver and receiver power in order to reduce wire power, they become less appealing with low-swing systems.

## 4.2 CMOS repeaters

Simple delay models of repeated wires are well-known [64]. They model a long wire broken up by CMOS inverters as a simple collection of linearized resistors and capacitors (see Figure 4.2). In this model, the only parameters of importance are the transistor width w, the wire length l, the normalized sum of NMOS+PMOS gate widths a, and some technology parameters representing transistor drive and parasitics<sup>1</sup>:  $R_v$ ,  $C_d$ ,  $C_g$ ,  $R_w$ , and  $C_w$ . Given a particular ratio of PMOS to NMOS device sizes (determining a), designers can typically only vary w and l to optimize performance.



Repeater unit

Figure 4.2: Simple model using inverters

This representation has many limitations: it cannot model the effects of risetimes upon delay, non-linear transistor currents, or bias-dependent capacitances. More detailed models could capture these effects, but are also too complex to provide intuition [71][72]. For our technology trending studies and first-order design decisions, this simple model's accuracy will suffice.

### 4.2.1 Optimizing for delay

The total delay down the repeated wire can be written as a product of *N*, the number of stages, and  $D_{stage}$ , the stage delay. Here, we are measuring delay at the 50% point, so that each term is multiplied by  $|\log 0.5| = 0.7$ .

$$N \cdot D_{stage} = N \cdot 0.7 [R_{\nu}a(C_g + C_d) + \frac{R_{\nu}}{w}C_w l + R_w \frac{C_w}{2}l^2 + R_w laC_g w]$$
(4.3)

<sup>&</sup>lt;sup>1</sup>Although wire parameters  $R_w$  and  $C_w$  depend on wire width and spacing, the device parameters are well approximated over different technologies by  $R_v=25000 \cdot L\Omega$ ,  $C_g=2fF/\mu m$ , and  $C_d=1fF/\mu m$ .

For center-skewed inverters, PMOS devices are twice the size of NMOS devices, making a = 3; and for legged devices,  $C_d = 0.5C_g$ . Following the derivation in Bakoglu, we can optimize for delay with respect to both the number of stages N as well as the driver width w, and end up with the following solution [64]<sup>2</sup>. In these expressions, the wire length and driver widths are in units of microns.

$$l_{opt} = 3\sqrt{\frac{R_{\nu}C_g}{R_wC_w}}$$
(4.4)

$$w_{opt} = \frac{1}{\sqrt{3}} \sqrt{\frac{R_v C_w}{R_w C_g}}$$
(4.5)

$$Delay/unit-length = 1.47\sqrt{FO4 \cdot R_w C_w}$$
(4.6)

In real circuits, designers rarely use the exact optimal segment length between repeaters, because the total end-to-end wire length is usually not an integer number of optimal-length segments. Similarly, they also often do not use the exact optimal device size due to cell libraries that have limited sizing granularity. What is the delay penalty for such non-optimal designs? If we write the actual segment length and device width as  $l = \hat{l} \cdot l_{opt}$  and  $w = \hat{w} \cdot w_{opt}$ , then the wire delay becomes

Delay = 
$$[0.34(\hat{l} + \frac{1}{\hat{l}}) + 0.4(\hat{w} + \frac{1}{\hat{w}})]\sqrt{FO4 \cdot R_w C_w}$$
 (4.7)

Writing this expression in terms of capacitive fanout is illuminating. Designers know that for logic gates driving large loads, a gain (product of capacitive fanout and logical effort [73]) of four is delay-optimal, but designers do not apply this heuristic to gates driving long wires. The analogy—sizing gates to break a large load into ramp-up stages versus sizing repeaters to break a long wire into serialized stages—is not perfect, because wire resistance shields downstream capacitance, and the wire itself has delay. However, we can still write the fanout as a ratio of capacitive load from wires and gates to the input gate capacitance:

$$Fanout = 1 + \sqrt{3} \frac{\hat{l}}{\hat{w}}$$
(4.8)

<sup>&</sup>lt;sup>2</sup>This result differs from that in Bakoglu because of the inclusion of diffusion capacitance, and it differs from that in Nose *et al.* because it separates the PMOS-to-NMOS sizing differences from  $C_g$  [64][65].

We note that the fanout is independent of the actual technology parameters, and depends strictly on the ratio of the segment length and driver width when normalized to their delay-optimal values (of course, the optimal length and width are highly technology-dependent). At optimal delay, when  $\hat{l} = \hat{w} = 1$ , the fanout is approximately 2.7. This formula points out that sizing by selecting a ratio of  $\hat{l}$  to  $\hat{w}$  (*i.e.* the fanout) does not necessarily return useful sizes and lengths. The fanout can remain near optimal for grossly under- or over-sized wires and devices.

Equation 4.7 is of the form  $x + \frac{1}{x}$ , so the penalty for grossly undersizing *w* and *l*—either underdriving the wires or overcrowding the wire with too many repeaters—is greater than grossly oversizing them. However, slightly mis-sized *w* and *l* lead to only slight performance penalties: underdriving the wires by 20% and over-spacing the repeaters by 40% leads to only a 4% delay penalty. In Figure 4.3 we see contours corresponding to 2% delay penalty increments as the stage length  $\hat{l}$  and driver width  $\hat{w}$  vary.

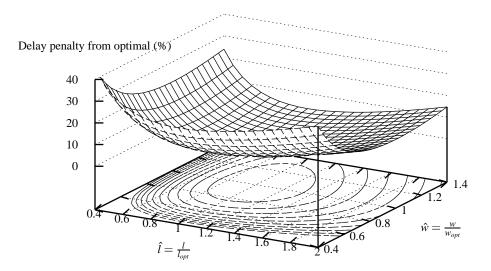


Figure 4.3: Delay sensitivity (2% contours)

We expect delay-optimal solutions to be inefficient. Delay optimization naturally sacrifices large power and energy costs in order to gain marginal delay benefits. At the optimal solution ( $\hat{l} = \hat{w} = 1$ ), especially, the shallowness of the contours implicates only minor performance gains for fairly large changes in  $\hat{l}$  and  $\hat{w}$ . From a design perspective, this delay optimality is wasteful; giving up small increments of performance to save energy leads to far better designs. In order to quantify gains in efficiency, we next construct models for repeater energy and optimize for energy-delay product.

### 4.2.2 Optimizing for energy-delay product

While a thorough energy analysis would not only consider switched capacitive current but also device leakage and crowbar (rush-through) current, this discussion will focus primarily on switched capacitances and ignore leakage currents, which are today an insignificant portion of total current. Crowbar currents will be approximated by a fixed multiplicative factor.

Simply adding up the various switched capacitances gives us the total switching energy below, where c represents a crowbar multiplier and L is the total wire length.

$$E_{tot} = C_w L V^2 \left( 1 + \frac{4.5}{3\sqrt{3}} c \frac{\hat{w}}{\hat{l}} \right)$$
(4.9)

Combining this result with the previous delay expression gives us the contours shown below in Figure 4.4 (using a crowbar of 1). The straight lines are 5% energy contours from the minimal-delay solution; there is no practical real "minimal-energy" solution outside of not doing the design. This picture helps to quantify the inefficiencies of a delay-centric design. For example, even within the first delay contour, when  $\hat{l}$  and  $\hat{w}$  are close enough to 1 to be within 2% of the delay optimum, the energy costs can swing by more than 15%. If, as suggested above, we set  $\hat{l} = 1.4$  and  $\hat{w} = 0.8$  and accept a 4% performance penalty, we can save 20% of the consumed energy.

The contours shown below in Figure 4.5 illustrate the optimal  $\hat{l}$  and  $\hat{w}$  for product of energy and delay. This graph arises from optimizing the products of Equation 4.7 and Equation 4.9, and in it, the benefits of longer wire segments and smaller drivers are clear. At the energy delay minimum, when  $\hat{l} = 1.59$  and  $\hat{w} = 0.62$ , the system has a constant fanout of 5.44, a delay of 11.5% worse than at the delay-optimal, and an energy of 28.4% better than that at the delay-optimal.

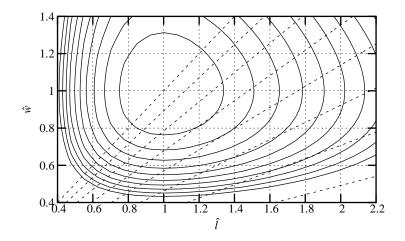


Figure 4.4: 2% delay contours (solid lines) and 5% energy contours (dashed lines). Energy increases towards the upper left direction.

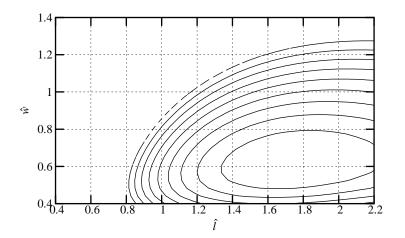


Figure 4.5: 2% contours for energy delay optimization

### 4.2.3 Other constraints and metrics

Simply optimizing for best energy delay, as we did above, will not yield a satisfactory repeater design, because it ignores other important design criteria. In this section we shall consider both reliability and bandwidth.

To maintain circuit reliability, designers need to limit signal risetimes at the end of wires to limit hot-electron degradation. As a transistor conducts current from its source to its drain, the highly energized carriers smack into the channel-gate interface, occasionally—though rarely—embedding themselves [74]. Over time, these trapped charges can gradually accumulate and eventually shift the threshold voltage of the transistor. Designers measure this degradation in a ten-year reduction in drive current by assuming that the transistor conducts only a certain percentage of the time. Ensuring this means that input risetimes (and falltimes) need to be kept short, typically under 40% of a clock cycle.

The previous expression for delay, Equation 4.7, uses a dominant time-constant model, and applying the same model to generate a 10%-to-90% risetime leads to:

RiseTime = 
$$FO4[1 + \hat{l}^2 + 1.2(\frac{\hat{l}}{\hat{w}} + \hat{l}\hat{w})]$$
 (4.10)

With a cycle time of 20 *FO4*s and a risetime limit of 40% of the cycle time, the expression above must be smaller than 8. Figure 4.6 shows the previous contours with the risetime constraint superposed. From this graph, a solution at  $\hat{l} = 1.6$  and  $\hat{w} = 0.6$  appears best.

For bandwidth, as we discussed earlier in Section 2.3.2, we should wait at least three time constants between tokens. For our inverter repeaters, a three- $\tau$  token separation can be expressed in two equivalent ways. First, our stage delay is defined at the 50% transition and thus represents 0.7 of a time constant, so waiting for three such time constants takes 3/0.7, or 4.3 stage delays. Mathematically, we can write:

Bandwidth = 
$$\frac{1}{T_{repeat}} = \frac{1}{4.3 \cdot FO4[\frac{1}{3}(1+\hat{l}^2) + \frac{\sqrt{3}}{4.5}(\hat{l}\hat{w} + \frac{\hat{l}}{\hat{w}})]}$$
 (4.11)

Alternatively, and perhaps more intuitively, we can specify a three- $\tau$  separation in terms of risetime at the output. If we wait for the output to rise to 95% of the transition (3  $\tau$ 's), this takes just slightly longer than the expression in Equation 4.10, which was based on a

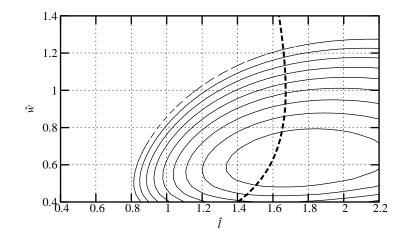


Figure 4.6: 2% contours for energy delay optimization with risetime constraint

10% to 90% transition, or 2.2  $\tau$ 's. Scaling the previous risetime expression by 3/2.2 gives an expression identical to that in Equation 4.11.

Notably, these expressions for repeat time do not depend on the wire parameters  $R_w$  and  $C_w$ . A repeated wire with skinny, high-resistance wires will have the same stage delay as one with fat, low-resistance wires. Of course, the low-resistance wires will have stages at farther intervals, so that more territory is covered per stage, leading to a faster wire with lower delay. However, for systems that are latency insensitive but which require high bandwidth, designers ought not engineer wires for better *RC* performance, but should instead maximally pack them for best density.

Finally, we note that optimizing for an energy bandwidth product would be uninteresting. Peak bandwidth comes from extremely short stage delays, so that the best bandwidth would come from stacking tiny inverter after tiny inverter, end upon end, with very little wire between stages. The energy cost of such a silly system would be large but offset by the gains in bandwidth. Adding a cost function for area (or equivalently, via utilization in "diving" down to the substrate from the wire layer) would make such a model more realistic.

To conclude this discussion of inverter repeaters, we list our figures of merit for the purely delay-optimal solution, compared to the energy delay optimum (under a risetime

	Value at	Value at	% worse than
	minimal	minimal	delay-optimal
	delay	energy∙delay	$(\hat{l} = \hat{w} = 1)$
Î	1	1.59	—
ŵ	1	0.62	_
Fanout	2.7	5.44	_
delay $(\sqrt{FO4R_wC_w})$	1.472	1.64	11.5%
Stage Delay (FO4)	1.44	2.55	77%
Energy ( $C_wL$ , $V_{dd} = c = 1$ )	1.87	1.34	-28.4%
Risetime (FO4)	4.5	8	77.4%
Repeat time (FO4)	6.14	10.9	77%
Energy delay product	2.75	2.192	-20.2%

constraint). We can easily save about 30% of the energy required per transition, by adjusting the device placement and driver size, but not much more.

Table 4.2: Minimal energy delay, under risetime constraints

# 4.3 General repeater models

In the previous section we used a inverter repeater model whose behavior was captured by a small collection of design parameters (driver size, segment length, and a PMOS-*vs*.-NMOS sizing term) in conjunction with some technology parameters (transistor resistance and parasitic capacitance, wire resistance and capacitance, and a crowbar current factor). Using those parameters we wrote expressions for delay, energy, stage-delay, risetime at the receiver, and fanout. We now extend this methodology by generalizing our repeater model.

Repeater circuit topologies may differ dramatically from simple inverters, but the functional differences that we care about can be distilled into a few parameters. First, a repeater may have multiple stages. This would give it non-zero internal delay and modify its input and output capacitances due to fan-up or fan-down. Second, a repeater might skew its PMOS-*vs.*-NMOS drive, affecting not only its input trip point but also its input capacitance and output drive. Third, a repeater may have more or less output drive than its gate width suggests, either from internal fan-up, from running the driver with different power supplies, or from putting the driver in a linear and not saturated region. Finally, a repeater, multi-stage or not, may burn more or less energy due to internal nodes switching.

To capture these possible differences, we modify our previous model by:

- Adding two terms to model nodes internal to the repeater. We use *d* (in *FO4*s) for the repeater's intrinsic delay and *e* (in units of gate capacitance) for the energy consumed in switching the repeater's internal nodes. Both are zero for the basic inverter case discussed above.
- Making the input capacitance term  $aC_gw$  instead of the inverter's  $3C_gw$ . This allows us to handle gates with skewed PMOS-*vs*.-NMOS sizes.
- Making the output capacitance term  $bC_gw$  instead of the inverter's  $1.5C_gw$ . This lets us handle repeaters with internal fan-up and/or output drive skews.
- Scaling the drive resistance by r to handle repeaters with increased current drive, either from fan-up or alternate circuit topologies. For the inverter, r = 1.
- Using a different time constant factor k instead of 0.7. This allows us to measure delays to thresholds other than 50% of the supply.

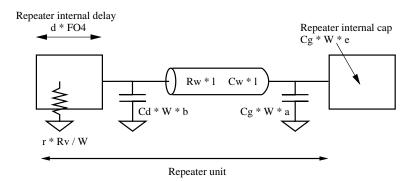


Figure 4.7: General repeater model

## **4.3.1** Delay and energy for the general model

Rewriting the delay equation results in

$$ND_{stage} = N[dFO4 + krR_{\nu}C_{g}(a+b) + kr\frac{R_{\nu}}{w}C_{w}l + \frac{k}{2}R_{w}C_{w}l^{2} + kR_{w}laC_{g}w]$$
(4.12)

As before, we can reduce this down into expressions for optimal segment length and driver width. The form of the resulting terms closely resemble those derived earlier in Equations 4.4-4.6. The optimal length  $l_{opt}$  still comes from balancing the repeater delay to the wire delay, and hence is again proportional to the square root of the gate-to-wire delay ratio. Similarly, the optimal driver width  $w_{opt}$  still comes from balancing the gate-wire and the wire-gate delays, and so is again proportional to the square root of the cross-delay terms. Only the proportionality constants have changed.

$$l_{opt} = \sqrt{\frac{18.9d + 2kr(a+b)}{k}} \sqrt{\frac{R_{\nu}C_{g}}{R_{w}C_{w}}}$$
(4.13)

$$w_{opt} = \sqrt{\frac{r}{a}} \sqrt{\frac{R_v C_w}{R_w C_g}}$$
(4.14)

When l and w might not be delay-optimal we can write the delay, fanout, and risetime as follows. Again, the structure of these equations closely parallels those previously derived, only with different constant terms.

Delay = 
$$\begin{bmatrix} \sqrt{d + \frac{kr(a+b)}{9.45}} \sqrt{\frac{k}{2}} \end{bmatrix} (\frac{1}{\hat{l}} + \hat{l}) \\ + \left[ k \sqrt{\frac{ar}{9.45}} \right] (\frac{1}{\hat{w}} + \hat{w}) \end{bmatrix} \sqrt{FO4 \cdot R_w C_w}$$
(4.15)

Fanout = 
$$1 + \sqrt{\frac{2(a+b)}{a} \frac{\hat{l}}{\hat{w}}}$$
 (4.16)

RiseTime = 
$$\left[\frac{2.2d}{k}\hat{l}^2 + \frac{r(a+b)}{4.3}(1+\hat{l}^2) + \frac{1}{4.3}\sqrt{\frac{ar(18.9d+2kr(a+b))}{k}}(\hat{l}_{\hat{w}} + \hat{l}\hat{w})\right]FO4$$
 (4.17)

As before we can also write an expression for bandwidth (or repeat-time). In the previous section we derived a lower-bound for repeat-time based on stage delay. But here, with an internal delay in the repeater, the maximum repeat time is set by the greater of the wire "delay" (a three- $\tau$  risetime delay) and the internal repeater delay. However, because the wire delay always exceeds the internal repeater delay, we can simply write repeat-time as a multiple of the risetime above.

Notice that the signal risetime is proportional to the repeater's internal delay, as written above in Equation 4.17. This makes sense because the inter-repeater segment length tracks the repeater's internal delay: the slower the repeater, the longer the wire between repeaters, and hence the slower the risetime.

# 4.4 Examples of repeater optimizations

The advantage of the general model formulation is that we can easily compare alternative repeater topologies and examine their sizing and placement from a delay and energy delay perspective. We will see that buffers and tristate repeaters require a small penalty in return for ease of placement (buffers) and some reconfigurability (tristates).

## 4.4.1 Example: Buffers

Buffers, or back-to-back inverters, may be used as repeaters to simplify the design of long interconnects, because they avoid logic inversions along the wire. Inserting an odd number of buffer repeaters causes no logic problems. However, they are slightly less efficient.

We assume the internal fanout inside the buffer to be 2.5; this provides for a good energy delay minimum (see Appendix). This makes our parameters:

- a = 3: PMOS devices are twice the size of NMOS devices.
- b = 3.75: At  $C_d = 0.5C_g$ , the buffer has a fanout of 2.5.
- r = 0.4: The output current drive is 2.5 times stronger.
- d = 0.66: At  $C_d = 0.5C_g$ , a FO2.5 is 3/4.5 times a FO4.

- e = 9: Internal load is  $C_d w$  of first stage plus  $2.5C_g w$  of the second.
- k = 0.7: We still use the 50% transition point for delay.

We can now plug these terms into our model and turn the crank. The delay-optimal stage length and driver width are

$$l_{opt} = 4.8 \sqrt{\frac{R_{\nu}C_g}{R_w C_w}}$$
(4.18)

$$w_{opt} = 0.36 \sqrt{\frac{R_v C_w}{R_w C_g}} \tag{4.19}$$

$$Delay = [0.551(\hat{l} + \frac{1}{\hat{l}}) + 0.249(\hat{w} + \frac{1}{\hat{w}})]\sqrt{FO4 \cdot R_w C_w}$$
(4.20)

$$D_{stage} = FO4[0.867(1+\hat{l}^2) + 0.393(\hat{l}\hat{w} + \frac{l}{\hat{w}})]$$
(4.21)

Compared to the inverter case, the length sensitivity is greater but the width sensitivity smaller. Because the buffer has internal delay, the optimal inter-stage wire length is longer, and hence varying the actual length has a greater effect due to its  $l^2$  term.

At the delay-optimal design point  $\hat{l} = \hat{w} = 1$ , we get a delay of  $1.6\sqrt{FO4 \cdot R_w C_w}$  and a stage delay of 2.518*FO4*s. This delay is 8.7% worse than at the inverter's delay-optimal.

The risetime at the receiver input, in FO4s, becomes

$$RT = 2.095\hat{l}^2 + 0.628(1+\hat{l}^2) + 1.232(\frac{\hat{l}}{\hat{w}} + \hat{l}\hat{w})$$
(4.22)

or 5.8 *FO4*s at  $\hat{l} = \hat{w} = 1$ , and the repeat time follows directly. The risetime constraint, superposed onto the delay contours in Figure 4.8 below, is much tighter than with inverters.

The energy consumed by the buffer is

$$E_{tot} = C_w L V_{dd}^2 [1 + 1.189c \frac{\hat{w}}{\hat{l}}]$$
(4.23)

or 2.189  $C_w LV^2$  at  $\hat{l} = \hat{w} = c = 1$ .

Figure 4.8 shows first the delay contours with a rise-time constraint. Figure 4.9 does the same, only with energy delay contours. Because the minimal delay-energy solution

results in a risetime that breaks the constraint, we instead use a solution that is 3.4% worse than optimum (but that is still 26% better than using  $\hat{l} = \hat{w} = 1$ ). At this optimal solution, buffers have segment lengths 16% longer than inverter segments, and drivers 55% smaller than inverter drivers. Table 4.3 compares the repeater characteristics with that from the original inverter system.

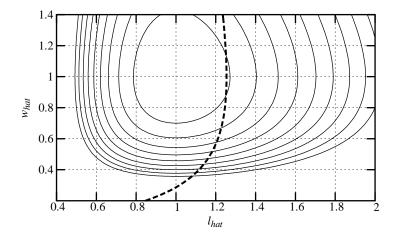


Figure 4.8: Delay sensitivity (2% contours) with rise-time constraint (dashed line)

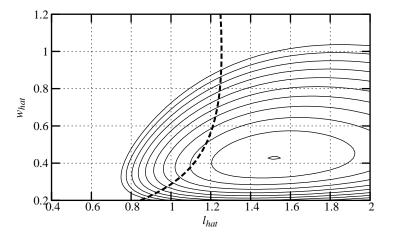


Figure 4.9: Energy delay sensitivity (2% contours) with rise-time constraint (dashed line)

From this analysis, buffered repeaters seem inferior to inverter repeaters: because of their internal delay, they have long stage separations, and thus are slower, have long repeat

		% worse than	% worse than
	Value	delay-optimal	inverters
		$(\hat{l} = \hat{w} = 1)$	$(E \cdot D \min)$
Î	1.141	_	(16% longer)
ŵ	0.437	—	(55% smaller)
Fanout	6.5	_	_
Delay $(\sqrt{FO4 \cdot R_w C_w})$	1.791	11.9%	9.1%
Stage Delay (FO4)	3.214	27.6%	26%
Energy ( $C_wL$ , $V_{dd} = c = 1$ )	1.456	-33%	9%
Risetime (FO4)	8	14%	0%
Repeat time (FO4)	10.9	14%	0%
Energy delay product	2.607	-25.6%	18.9%

Table 4.3: Buffer minimal energy delay, under risetime constraints

times, and burn more energy. However, buffered repeaters are still popular on microprocessor designs (or other large designs with irregular wiring patterns), because they preserve logic levels. This makes buffered repeaters ideal for back-end timing-driven repeater insertion, while inverter repeaters require the designer to beware of an odd number of repeating stages.

Another reason often cited for the continued use of buffered repeaters is lower crowbar current. With inverter repeaters, each stage is large and sees a relatively slow input slew, leading to significant crowbar current (the *c* term is nontrivial). With buffered repeaters, the small leading stage sees a slow input slew while the large trailing stage sees a fast input slew. This means that the trailing stage's *c* term is smaller than the leading stage's, a detail ignored by the previous analysis. This effect is small, however: if we assume a crowbar factor of 1.25 for the slow-slewing input and 1.05 for the fast-slewing input, then accounting for this effect at  $\hat{l} = 1.045$  and  $\hat{w} = 0.435$  leads to a difference of under 5% in consumed energy.

## 4.4.2 Example 2: Tristates

Tristate repeaters offer simple wire programmability. A simple NEWS set of repeaters allows for a signal to travel in each of four directions, as shown in the figure below. This scheme is insufficient as a network since it omits buffering or any control structures. However, for now we will assume that control logic happens outside of the critical path, and that the system is perfectly scheduled and so requires no buffering; the following results would therefore be optimistic.

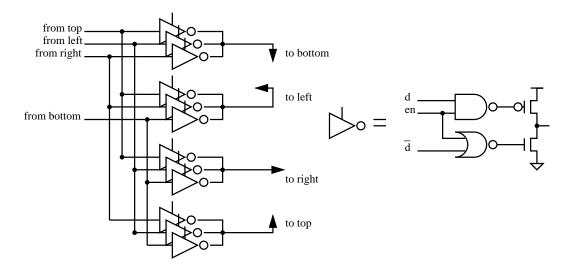


Figure 4.10: Simple model using tristate repeaters

These tristates have an output PMOS driver twice the size of the output NMOS driver. A 2-input NAND gate driving the PMOS has a fanout of 3, and a 2-input NOR gate driving the NMOS has a fanout of 2.5. Hence, the total input gate size (from one NAND and one NOR), relative to the output NMOS driver size, is  $\frac{2}{3} + \frac{1}{25}$ .

Fitting this system into our framework requires us to set

- a = (2/3 + 1/2.5) \* 3 = 8/7.5: We have  $C_{in}$  for each of three tristates.
- b = 3/2 \* 3 = 4.5:  $C_d$  is half of  $C_g$ , and there are three diffusions.
- r = 1. The output drive is no stronger or weaker.
- d = 1: Internal delay is approximately 1 *FO4*.
- e = 8/(7.5 \* 2) + 3 = 53/15: Internal nodes of only one tristate.
- k = 0.7: Still use the 50% transition point for delay.

These parameters result in the following expressions for stage length, driver width, and delays:

$$l_{opt} = 6.5 \sqrt{\frac{R_v C_g}{R_w C_w}}$$
(4.24)

$$w_{opt} = 0.56 \sqrt{\frac{R_v C_w}{R_w C_g}} \tag{4.25}$$

$$Delay = [0.741(\hat{l} + \frac{1}{\hat{l}}) + 0.407(\hat{w} + \frac{1}{\hat{w}})]\sqrt{FO4 \cdot R_w C_w}$$
(4.26)

$$D_{stage} = FO4[1.570(1+\hat{l}^2) + 0.863(\hat{l}\hat{w} + \frac{l}{\hat{w}})]$$
(4.27)

The delay-optimal length and  $\hat{l}$  sensitivity term is twice that of the inverter case, while the width and  $\hat{w}$  term about the same. The delay-optimal delay is 2.297, or 56% longer than the inverter's delay.

Because the tristate repeaters have relatively low current drive for the load they drive, the risetime constraint presents problems. We can write the risetime as

$$RT = 3.143\hat{l}^2 + 1.791(1+\hat{l}^2) + 2.709(\frac{\hat{l}}{\hat{w}} + \hat{l}\hat{w})$$
(4.28)

and see that at the delay-optimal point, the risetime at the receiver is 12.1 *FO4s*, which violates our 8 *FO4* constraint. Instead of using  $\hat{l} = \hat{w} = 1$ , our baseline needs to be at  $\hat{l} = 0.7$  and  $\hat{w} = 1$ , which makes the risetime 8 *FO4s*. Doing this invokes a delay penalty of 4% and an energy penalty of 21%.

Energy consumed by the tristate repeaters is 1.964  $C_w LV^2$  at  $\hat{l} = \hat{w} = c = 1$ .

$$E_{tot} = C_w L V_{dd}^2 [1 + 0.964c \frac{\hat{w}}{\hat{l}}]$$
(4.29)

As before we minimize the delay-energy product, and show the results in the following table and graphs. These contours are now at 4% steps. The minimal delay energy is far from the risetime curve, so that we need to pick a design point that leaves considerable performance on the table: at  $\hat{l} = 0.633$  and  $\hat{w} = 0.51$  the delay-energy product is 4% worse than at the delay-optimal and 35.6% worse than at the delay-energy minimum points. Also,

	Value	% worse than delay-optimal	% worse than inverters
	value	$(\hat{l} = \hat{w} = 1)$	$(E \cdot D \min)$
Î	0.633	—	(14% shorter)
ŵ	0.510	_	(20% smaller)
Fanout	3.726	_	—
Delay $(\sqrt{FO4R_wC_w})$	2.648	15.3%	61%
Stage Delay (FO4)	3.548	-27.1%	39%
Energy ( $C_wL$ , $V_{dd} = c = 1$ )	1.777	-9.5%	33%
Risetime (FO4)	8	14%	0%
Repeat time (FO4)	10.9	14%	0%
Energy delay product	4.705	4.26%	114.6%

at this design point, the segment length is 14% shorter than the inverter's design point, and the driver width 20% smaller.

Table 4.4: Minimal energy delay, under risetime constraints

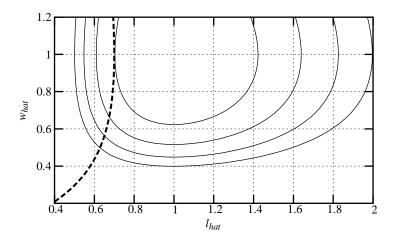


Figure 4.11: Delay sensitivity (4% contours) with rise-time constraint (dashed line)

This analysis reflects the costs of some primitive reconfigurability in wires. If we wanted to replace a regular grid of inverter repeaters with tristates for more flexible routing, we would need to shorten the segments and make the overall device sizes smaller. The resulting system would have moderately worse performance than inverters but have much

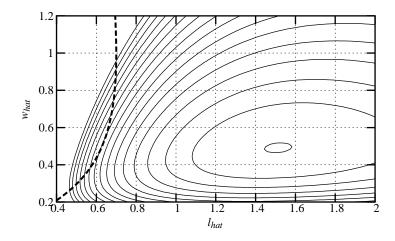


Figure 4.12: Delay-Energy 4% contours with rise-time constraint (dashed line)

more flexibility. The bandwidth (repeat time) is unchanged, because the risetime is the tightest constraint in the system.

# 4.5 Summary

This chapter examined the delay and energy costs of some basic repeater structures. We first considered a model for simple inverter repeaters, and then generalized this model for other circuits like buffer repeaters or tristate repeaters. We found that buffers are within 10% in both delay and power to inverter repeaters. This cost may be small when the repeated wires lack structure and homogeneity, since buffers are logically non-inverting and are hence much easier to insert. Tristates are 60% slower than inverters and cost 30% more energy, but do allow simple wire reconfigurability.

Our models also allow us to explore reoptimizations of these repeater schemes. The traditional solution for simple inverter repeaters provides best delay, but loses energy efficiency in the process. By re-optimizing for alternate metrics, such as energy delay product, we can save nearly 30% of energy for only an 11% penalty in delay. This resizing corresponds to a 20% improvement in energy delay product.

This 30% energy savings is rather small, and greater power savings will require more aggressive design techniques. In the next chapter we will consider the design of low-swing

## 4.5. SUMMARY

circuits, which hold the promise of power-savings much larger than those available from simple resizing.

# **Chapter 5**

# **Low-Swing Repeaters**

In the previous chapter we considered local optimizations on repeated wires by varying segment length and repeater driver width in order to make them more energy-efficient. This is akin to trying to save fuel by driving a Cadillac with the air conditioner turned off: the underlying vehicle's inefficiencies still dominate. In addition, reliance upon repeater schemes to overcome wire resisitance leads to layout and via congestion. The Smart Memories modular architecture has 128-bit, 4.5mm long links between module cores and, as shown in the previous chapter, would need four repeater stations along the links. Each repeater station would need 128 holes drilled down and then back up through all intermediate wiring layers, as well as a channel for 128 large inverters. While the area cost of the inverters, relative to the total die, lessens with technology scaling, the via and wire congestion costs remain high, significantly complicating physical design.

This chapter discusses how using low-voltage signaling effectively reduces energy consumption, and, if used to drive each Smart Memories link in a single stage, avoids the physical design complications of multiple repeater stages.

# 5.1 Benefits and costs of low-swing signaling

Reducing the voltage swing of our global interconnects linearly reduces the required energy, because capacitors need swing only a fraction of the original voltage. In addition, if we can provide the charge from an additional reduced power supply, then we can gain quadratic energy savings.

Simply lowering the supply for all of the elements in a global wiring system would certainly reduce the required energy, but at a performance cost, as drive resistance grows as supply voltages reduce. We can approximate the current of a transistor as proportional to  $(V_{gs} - V_t)^{\alpha}$ . The actual value of  $\alpha$  depends on the degree of velocity saturation, and 1.2 works well for today's transistors. We call the threshold voltage normalized to the power supply t, and a value of  $t = \frac{1}{4}$  is typical. In this case, the saturation drive resistance for lower power supplies, when normalized to full-swing drive resistance, becomes

$$R = \frac{x(1-t)^{\alpha}}{(x-t)^{\alpha}}, x = \frac{V}{V_{dd}}$$
(5.1)

and as Figure 5.1 shows, the performance penalty is not terribly bad for velocity saturated devices. Reducing the voltage by a factor of two, saving 4x the energy, doubles drive resistance; for gates driving long wires this less than doubles the delay because the wire resistance is unchanged. For devices that are not velocity saturated, the story worsens: reducing the voltage by only 30% doubles drive resistance. Using low- $V_t$  devices helps significantly, but leakage current trends will increasingly reduce the availability of low- $V_t$  transistors.

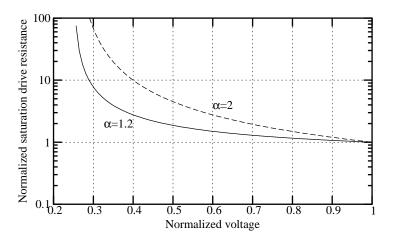


Figure 5.1: Drive resistance versus power supply, for  $V_t = \frac{1}{4}V_{dd}$ 

While this wholesale voltage reduction would certainly improve the energy-delay product, lowering the voltage swing and supply for *only* the long wires of a global interconnect system would do significantly more. It would maintain large energy savings: the long wires between repeater stations represent much of the capacitance in the system. But keeping full voltage swing for the logic stages leading up to the wire drivers retains high gate overdrive voltage, thus maintaining low drive resistance and high speed.

Using such a low-swing system, and designing it to drive the 4.5mm global interconnects of the Smart Memories chip in a single stage, would solve both power and layout congestion problems. It would also introduce a host of new issues, including how to strobe the receivers to tell them when to look at the data, and how to minimize the effects of noise that corrupts data.

However, we can turn the characteristics of the Smart Memories global interconnection network to our advantage and exploit them to simplify our problem. The homogeneity of the global fabric allows a fixed latency for all point-to-point connections, making a clocked receiver feasible and inexpensive. Also, because each link needs only to fit within a clock phase, the timing of each link can be relaxed somewhat and not squeezed to obtain minimal delay. In this case, a small timing overhead for the clocking of the driver and receiver is not only possible but in fact harmless to overall system performance. Finally, the regularity and structure of the global network means that wire environments are stable and wellcharacterized, making them well-suited for sensitive circuits like low-swing wires.

This chapter will outline a low-swing system suitable for the Smart Memories global network and focus on its component parts in depth. A sample design of this low-swing system, set for a contemporary 0.18- $\mu$ m technology, elucidates some of the engineering tradeoffs. This design emulates the target Smart Memories fabric: 10mm wires in a 0.18- $\mu$ m technology have the same resistance, although higher capacitance, than 4.5mm wires in a 0.10- $\mu$ m technology. The system is designed to run at an aggressive repeat rate of 10 *FO4s*, matching the target system speed of the Smart Memories machine.

# 5.2 Low-swing repeater systems

A low-swing repeated-wire system, simplified in Figure 5.2, contains three parts. First, a low-swing transmitter using overdrive and distributed equalization for higher performance transmits a signal down a long wire. Second, the wires that carry this signal are differential and twisted to reduce noise effects. Third, a clocked regenerative amplifier receives the signal and converts it to full-swing levels suitable for logical operations such as queueing or arbitration.

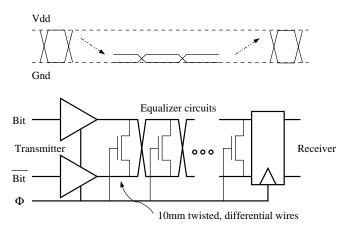


Figure 5.2: Generic low-swing repeated wire

## 5.2.1 Low-swing transmitter circuits

The transmitter circuits used in this low-swing system use a number of techniques to communicate quickly over long distances, including pre-emphasis, pre-equalization, and linear drive resistance. First we discuss how to limit voltage swing on the global wires.

### **Restricting swing**

In our proposed system we rely on a separate reduced power supply for increased energy efficiency [75]. This is a minor cheat: it removes complexity from the transmitter circuit, because voltage generation becomes free, and adds complexity to the power grid design. Although extra grids may require more layout work, they are certainly feasible. Power

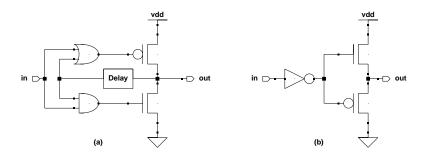


Figure 5.3: Cutoff drivers. (a) Logical cutoff. (b) Self-cutoff.

grids are robustly designed to maintain their voltage under peak current loads, minimizing  $R_{grid} \cdot i_{dc}$ , as well as under peak current swings, minimizing  $L_{grid} \cdot di/dt$ . If we replace global buses with low-swing versions requiring their own power supplies, the demands placed on the original full power supply will be reduced, and that grid can be made appropriately less robust, using narrower wires and sparser gridding. In that newly-freed space we can route the alternate power supply grid, which needs fewer resources than those we just made available, because they are low-swing and thus more energy-efficient. For example, suppose that without low-swing wires, a chip runs ten power lines every  $30\mu$ m on a particular metal layer. If the global buses to be replaced by low-swing buses account for 10% of the total capacitive load, and they now need their own supply, we can get by with nine full-voltage power lines and one reduced-voltage power line every  $30\mu$ m. This line of reasoning slips when discussing wide original power lines: a  $10\mu$ m wide wire cannot be divided into a  $9\mu$ m and a  $1\mu$ m wire without first consuming more area for wire-to-wire spacing. However, power lines today are already interdigitated to eliminate skin effect inefficiencies, making this situation uncommon.

A number of designers have proposed other means of lowering voltage swing, including constructing on-chip voltage regulators [76]; cutting off the driver early [77][78][79], [80][81]; and exploiting charge sharing [82][83][84][85][86][87]. We briefly consider these last two schemes here.

Designers can limit voltage swings by using cutoff circuits. In these schemes, the pullup or pulldown paths are either explicitly disabled by a logic chain or they turn themselves off by reducing  $V_{gs}$ . In either case, the energy savings is linear with swing.

### 5.2. LOW-SWING REPEATER SYSTEMS

Logical cutoff drivers use local feedback from the output of the driver to trip a logic path back to the driver's enable, as in Figure 5.3a. Once the output has gone "far enough," the driver is disabled and the output line left to float at its intended voltage [77][78]. To decide whether or not the output has indeed gone far enough, the driver can either use a replica of the receiver or simply "dead-reckon" the transmission with a local delay.

In the first case, when the driver uses a local copy of the receiver to measure the wire swing, a problem may arise from line resistance. Because the wire itself has intrinsic delay, using the "near-end" voltage at the driver as a proxy for the "far-end" voltage at the receiver can lead to timing errors and cutting off the driver too early. The second case, using a deadreckoned delay, is simpler, especially for homogenous and regular modular global wires that all have the same nominal delays.

In either case, cutoff circuits suffer from noise sensitivity. The driver stops driving once cut off, giving the signals little immunity to injected currents from other wires. In the noise framework of the previous chapters, the victim time constant becomes infinite, maximizing coupled noise. Engineering this noise demands increased wire-to-wire spacing, adding some area overhead.

Logical cutoff circuits are used in specialized circuits, like in memory arrays with small memory cells weakly pulling down bitlines. A timing signal cuts off the bitline drive once they get pulled down part-way, saving both delay and power. However, memory array bitlines are much shorter than repeated wire segments, and by the time their drive cuts off, the bitline sense-amps have already fired. Thus the bitline exposure to noise, both physical and temporal, is minimal. More importantly, the timing of their cutoff signal comes from a replicated memory path, which matches wire delays to wire delays and gate delays to gate delays, resulting in far closer delay matching than in the logical cutoff circuits.

The second type of cutoff driver stops itself when  $V_{gs}$  falls below  $V_t$ . Figure 5.3b shows an example [79]; a variant, not shown in the figure, restricts the pullup by adding an NMOS transistor in series with the pullup PMOS [81]. Poor performance cripples these drivers. An NMOS device pulling down maintains a good drive resistance all the way to the bottom, and this drive resistance actually improves as the device becomes linear. An NMOS device pulling up, by contrast, has a drive resistance that approaches infinity as the device gradually turns itself off, resulting in extremely long tails on the voltage waveforms and indeterminate rise- and fall-times. In addition, the high driver resistance at small  $V_{ds}$  means that small noise disturbances will not be restored. For our repeater model, the *r* term would rise dramatically, increasing delay and making the delay and risetime equations difficult to fit.

Another way to limit driver swings uses capacitive charge-sharing, in which the load capacitance connects to another reservoir of charge, balancing the voltages. This other reservoir of charge may be shared among many circuits globally, predefined to be between complementary wires in differential signaling, or shared amongst the bits in a bus.

Global charge-sharing drivers restrict voltage swing by shorting their outputs to a large shared capacitor [84][85]. On a falling edge charge gets shunted from the output node to the global capacitor, and on a rising edge the output node gets pulled to  $V_{dd}$  while the global capacitor sheds excess charge onto a dump capacitor. Since for small swings, the dump capacitor needs to be much larger than the circuit load capacitors, which are themselves long wires with large capacitances, these on-chip voltage regulators have a large area overhead.

Another scheme of global charge-sharing, more suited to a bus, shorts all the "0" bits together with a large pre-discharged shared capacitor [83]. The more bits that switch, the lower the swing. The drawback to this technique lies in the multiple common-mode levels required of the receivers, but it does offer  $\frac{1}{n}$  energy savings, where *n* is the number of bits in the bus.

Sharing charge between differential bits on a bus saves energy over schemes that would otherwise switch every cycle [87][75][82]. Extending this scheme to share charge amongst *all* the bits on a bus, so that each bit has a discrete range and voltage swing, offers  $\frac{1}{n^2}$  energy savings [86]. As with the previous technique, however, it also requires a host of different common-mode voltages for the receivers.

Both bus-shorting schemes trade design complexity for power supply distribution. Simply using a  $\frac{1}{n}$  power supply offers  $\frac{1}{n^2}$  energy savings for less complexity, but it also requires an alternate power supply to be shipped around the die.

#### **Overdrive pre-emphasis**

Inherent offsets and mismatches at the receiver determine the minimum signal that must be delivered to it. Using a drive voltage significantly higher than this minimum level, however,

### 5.2. LOW-SWING REPEATER SYSTEMS

can significantly increase performance [88]. If a particular receiver demands an input signal swing of 100mV, using a 200mV drive voltage implies that the signal is "done" at the 50% point, and hence a single-time constant model would predict a delay of  $0.7\tau_{wire}$ . Using a 400mV drive voltage to generate the same 100mV split, however, needs a 25% transition and this takes only  $0.4\tau_{wire}$ —a significant speedup.

This overdrive technique, using driver-side pre-emphasis to overcome wire loss, is illustrated below in Figure 5.4. In this example, an  $8.1-\mu m$  driver in a 1.8V,  $1.8-\mu m$  technology swings a wire 10mm long with a variable driver voltage. The speedup is bounded; multiple higher-order poles in a wire's transfer function limit the speedup to about a factor of 2.5x because they cause the "leading corner" of the transition to ramp slowly, as shown in the figure by the curves bunching up as the drive voltage increases. A sweet spot sits around 4x the target voltage (twice the typical drive voltage), giving about a 2x speedup.

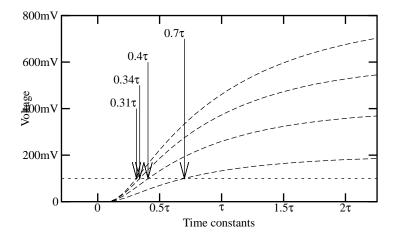


Figure 5.4: Overdrive transmitter benefits

By increasing the drive voltage we also increase the energy cost, because we have raised the supply from which current is drawn. If we overdrive in both high and low directions by doubling the required drive voltage, the energy cost will quadruple. Because the power consumed by the wires is small, this decrease in efficiency need not be fatal. However, we observe that when we use overdrive, we need not let the entire wire swing the entire overdrive voltage. Instead, by cutting off the driver *as soon as* the receiver sees the necessary voltage split, we can prevent most of the wire from switching, hence regaining some of that extra energy. Afterwards, we short the differential wires together in preparation for the next bit.

Figure 5.5 shows the results of a simulation of this scheme. On the left is drawn the transmitter ("near-end") voltage as well as the receiver ("far-end") voltage, and on the right is the voltage profile along the length of the wire, taken as a "snapshot" at time=6.4FO4s. The profile graph suggests that the energy consumed by the wire is bounded by  $C_{wire}V_{drive}0.5(V_{target} + V_{drive})$ . This rebalancing technique is known as pre-equalization, and is discussed next.

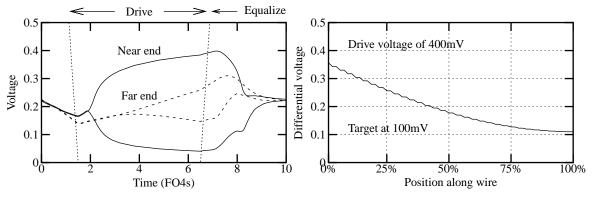


Figure 5.5: (a) Simulation of wire split at receiver end and (b) Wire differential voltage profile at t=6.4FO4

### **Pre-equalization**

As shown in the previous figures, the wire system uses differential signaling, mostly to convert potentially fatal coupling noise into fairly harmless common-mode level shifts. Because we have two wires per bit, we can pre-equalize the wires by shorting them and recycling charge before each new transmitted datum. Starting the two wires at a mid-way voltage allows for lower latency: developing a differential voltage of 100mV takes less time if the two wires both started at 200mV and immediately split apart, rather than if they started at 0 and 400mV and had to cross first.

Pre-equalization can also be viewed as simply another method of pre-emphasis. Overdrive modifies the voltage swing *final* voltage, while pre-equalization modifies the voltage swing *start* voltage: with it, the differential voltage across the two wires starts at zero, and

### 5.2. LOW-SWING REPEATER SYSTEMS

without it, at a negative voltage. In our design, pre-equalized latency down a 10mm wire dropped by another factor of two; thus both techniques of pre-emphasis contributed a total speedup of nearly 4x.

Pre-equalization requires twice the work on each pair of wires: for each token transmitted, the wires must first split apart and then recombine. However, twice the work need not require twice the energy, because pre-equalization recycles charge. During the evaluation phase of the wires, charge must flow from the positive supply to pull up one wire through half of the voltage swing, while the other wire discharges to ground. During the equalization phase, the high voltage wire shares, or recycles, its charge to the low wire, so that no charges must be pulled from the positive supply in order to balance the wires. Thus, ignoring for now clock power, pre-equalized wires consume the same energy as non-prequalized wires.

Each bit transmission now requires two half-swing wire transmissions, and is twice as "busy": we need to ping-pong between an "active" phase and a "balance" phase. Fortunately, we avoid a bandwidth penalty because the wires can cycle twice as fast as before. Figure 5.6 shows a simulation comparing pre-equalized versus non-equalized drivers. This comparison is not perfectly fair, as the non-equalized drivers were not re-optimized for sizing, but it does show that equalization lowers latency enough to do both drive and balance in the same time it would take to drive non-equalized wires.

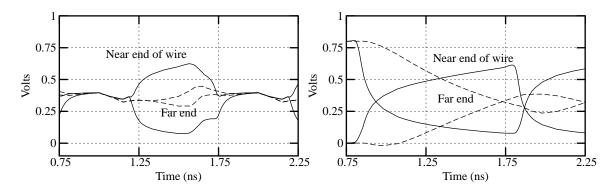


Figure 5.6: Equalized versus non-equalized drivers

One notable drawback of pre-equalization is that it makes energy consumption constant and not data-dependent. For data patterns that simply repeat bits, a non-equalized system would not consume any energy, because the wires would not transition. The pre-equalized system, on the other hand, still precharges and then evaluates each cycle, consuming a smaller but consistent amount of energy. The benefits of pre-equalization still outweigh the drawbacks; pre-equalization simplifies the overshoot driver pre-emphasis. Without pre-equalization, a consistent series of data bits ("1111111...") will force the data lines farther and farther apart, until they actually reach the maximum swing allowed by the reduced power supplies. Then, the next bit flip ("0") will have a much larger delay than normal, because it has to overcome the "history" effect of repeated data bits. One way to avoid this problem is to stop driving the wire on repeated data, so that the wires "sit" at the proper voltage. Another way is to use two possible drive levels, one level when driving a bit that differs from the previous bit, and another level when driving a repeated bit. These solutions either require careful noise minimization or multiple power supply levels. Pre-equalization provides an alternate solution.

Another challenge with pre-equalization lies in the difficulty in clamping two wires together. Wires act similar to limp strands of cooked pasta lying on a table: if we grab the ends of two such noodles and tried to pull them together, they would meet very slowly. We will discuss the circuits required for pre-equalization in the wire engineering section below.

### Linear-mode drive transistors

The actual drivers, shown in Figure 5.7, use NMOS pullups and pulldowns to exploit their much lower linear drive resistance at small  $V_{ds}$ . Their sizes and the overdrive supply voltage were chosen to enable a 100mV separation on 10mm-long wires within 5 *FO4*s. This required 8.1- $\mu$ m drivers in an 0.18- $\mu$ m technology and a drive voltage of 0.5V.

Using NMOS pullups and NMOS pulldowns with a dedicated power supply is a simplification of the idea of "stacked" devices, in which two inverters sit one atop the other, splitting the voltage supply evenly between them [89][90]. These NMOS-only drivers are similar to those used by Burd [91], and have a resistance that can be illustrated in Figure 5.8. This graph shows driver resistance of NMOS and PMOS devices in a MOSIS  $0.18-\mu$ m technology as common-mode voltage and swing extend vary.

In this graph, the x-axis is the center of the swing—typically 50% in full-swing CMOS and the different curves represent voltage swings from 10% to 100% of the supply. In other

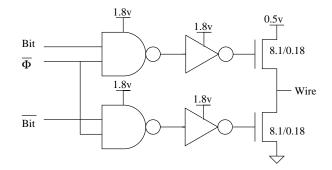


Figure 5.7: Transmitter, one of a pair, sized for a 0.18- $\mu$ m technology

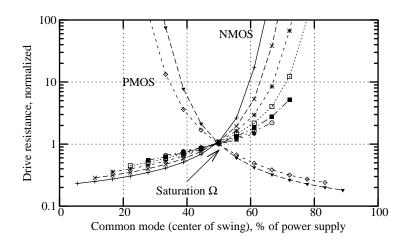


Figure 5.8: Driver resistance in a MOSIS 0.18- $\mu$ m technology

words, the voltage ranges around  $V_{CM} \pm V_{swing}$ . The y-axis is drive resistance, normalized to that of a full-swing CMOS pulldown; that is, to the resistance gotten from a swing of  $\frac{V_{dd}}{2}$  and a center voltage of  $\frac{V_{dd}}{2}$ . In all curves, the gate is held at  $V_{dd}$ .

The two lonely curves arcing the other direction are representative curves for PMOS devices, for swings of  $0.2 \cdot V_{dd}$  and  $0.3 \cdot V_{dd}$ .

This graph shows two implications for driving low-swing signals. First, in part due to the difference between saturation and linear resistances, NMOS devices have their lowest driving resistance at very low common mode and swing voltages, and *vice versa* for PMOS devices. At an extremely low swing voltage of  $0.1 \cdot V_{dd}$  centered at  $0.1 \cdot V_{dd}$ , the driving resistance is one-fourth of the CMOS saturation resistance. In the repeater formulation of the previous chapter, this leads to a smaller drive factor *r*.

Second, as the center voltage increases past  $\frac{V_{dd}}{2}$ , causing the device's  $V_{gs}$  to fall, the NMOS drive resistance escalates rapidly; the converse holds true for PMOS devices. Hence, NMOS pulldowns for small swings centered above  $\frac{V_{dd}}{2}$  are a poor choice, as are PMOS pullups for swings centered below  $\frac{V_{dd}}{2}$ . Much better would be swings centered below  $\frac{V_{dd}}{2}$  that use both NMOS pulldowns and NMOS pullups, or swings centered above  $\frac{V_{dd}}{2}$  that use both PMOS pullups and PMOS pulldowns. Either of these systems would attain a lower r drive factor, although the NMOS version would perform slightly better. Centering small swings at  $0.5 \cdot V_{dd}$  would give about the same drive resistance as full-swing CMOS logic.

## 5.2.2 Wire engineering

In order to engineer the wires to carry low-swing signals we used three techniques: noisesuppressing techniques of differential signaling and twisted routing, extra width to overcome resistive loss, and distributed balance drivers to enable pre-equalization.

### Differential and twisted wires

The regularity and homogeneity of the global wires leads to well-characterized capacitive noise profiles. It does not, however, necessarily ensure a safe-by-construction environment, suggesting the use of noise minimization techniques. The wires are both differential and twisted as a result. Running differential wires provides two benefits. First, noise sources

#### 5.2. LOW-SWING REPEATER SYSTEMS

equally exposed to the two wires, such as a wide precharged bus running perpendicularly over the wires, injects not differential but rather common-mode noise, which most receivers can reject. Second, a wire whose complement sits next to it has a very close—and thus low impedance—return path. The current loop for that wire becomes minimal, especially as much of the current flow returns along the wire length as displacement current across coupling capacitors, thereby minimizing inductive coupling.

Differential wires do cost twice as many wires for routing, but under technology scaling, the number of total wire layers continues to increase slowly. As a result, the number of wire pitches per gate increases, making the incremental cost of differential signaling smaller with technology scaling.

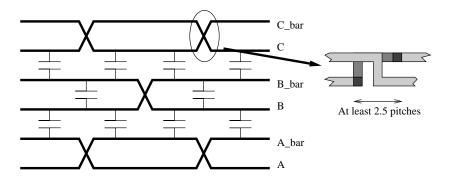


Figure 5.9: A typical twist solution

Wire twisting at regular intervals down the bus eliminates differential noise injection from neighboring wires and also helps to average out Miller capacitance to other bits, thus minimizing data-dependent delays. Figure 5.9 shows a typical twisting pattern that uses few metal resources from a neighboring metal layer. In this case, wire "B" couples equally to both "A" and "A\_bar," as well as equally from "C" and "C\_bar." It suffers from high total capacitance, because "B" and "B\_bar" are tightly coupled and always moving in opposite direction. Other twisting schemes can reduce this worst-case Miller capacitance, but typically consume more wires in a lower or upper layer for farther twists.

Using twisted wires demands some care. Each twist requires vias placed in-line with the wire. Without using arrayed vias to reduce via resistance, the twists may add non-trivial additional loss to the line. In this figure, the twists happen at regular and evenly-spaced intervals. However sometimes a better twisting arrangement can be used. Figure 5.10 shows this situation, where capacitor C1 and capacitor C2 ideally push and pull equivalent current onto and off of the victim.

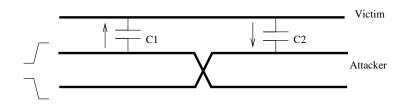


Figure 5.10: Optimized twisting would put the twist at 70% down the wire

The problem is that the noise contributions from C1 and C2 are not equal. Current from C1 will split: most of it will drain harmlessly down the victim's drive resistor at the left of the line, and the rest will flow down to the far end, causing potential problems at the victim's receiver. Conversely, negative current from C2 will also split: some will go the long route to the left, and most will go to the right, towards the victim's receiver. Thus, for a twist point at mid-way down the wire, the currents from C2 will overpower the currents from C1, simply because C2 is closer to the victim's receiver than C1 is. Moving the twist point to the right may help, by reducing the proximity-aided influence of C2 on the victim's receiver. Too far, however, and the diminishing value of C2 will make C1 too important.

A first-order model of this assumes a constant attacker edge rate down the wire, so that injected currents are proportional to C1 and C2. The current due to C1 must flow down the victim wire to the receiver before it matters. This delay matches to first order the attacking edge flowing down the attacker wire, so that C1's current should reach the victim receiver somewhat aligned with C2's current. Hence we can simply consider the peak current levels due to C1 and C2.

Injected currents see a resistor divider whose values are simply the left and right portions of the victim wire. In this case, balancing the currents that flow to the right of the victim wire, if the twist is at x (normalized between 0 and 1), requires:

$$C1 \cdot \frac{x}{2} = C2 \cdot (x + \frac{1-x}{2})$$
 (5.2)

$$\frac{x^2}{2} = (1-x)(x+\frac{1-x}{2})$$
(5.3)

$$x = \frac{\sqrt{2}}{2} \tag{5.4}$$

$$x \approx 0.7 \tag{5.5}$$

Figures 5.11 and 5.12 show simulations of a 5mm long wire twisted as above, and the twist location varied along the 5mm length. In this case, the switching attacker has a  $100\Omega$  drive resistance, and the quiet victim a  $50\Omega$  drive resistance. In the simulation, the attacker's wavefront has a decreasing slope and hence *C*2 injected current is correspondingly less, so the optimal twist point moves slightly to the left. However, the 70% twist point is a good first estimate.

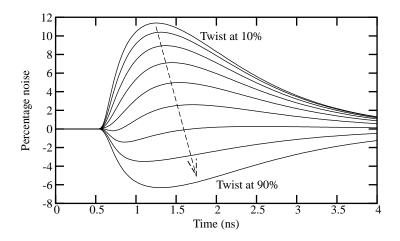


Figure 5.11: Simulated noise, 5mm wires,  $100\Omega$  and  $50\Omega$  drive resistances

#### Increased width to combat resistance

The wires were drawn at twice the minimum pitch to avoid resistive loss. This extra width, combined with the techniques of transmitter overdrive and pre-equalization, allowed the system to avoid intermediate repeater stages along the wire and cycle in 10 *FO4*s. Eliminating the intermediate repeaters was a huge win from a layout complexity perspective. On-chip wires are cheap and plentiful, so unlike conventional off-chip networks, on-chip

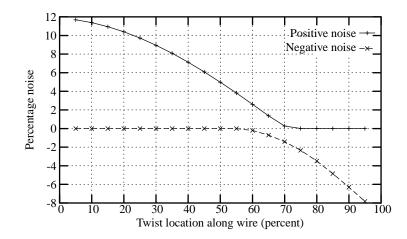


Figure 5.12: Peak of simulated noise, with optimal around 70%

networks are generally not wire limited. Thus avoiding intermediate repeaters by using wider wires is a good tradeoff for these systems.

#### **Circuits for pre-equalization**

Achieving the charge recycling pre-equalization mentioned above required additional care. The equalization cannot use the same pre-emphasis technique as the data driver to achieve low-latency, because overshooting the balancing of the two wires wastes energy.

An aggressive pre-equalization technique would launch a wave of wire balancing from the driver end to the receiver end shortly after sending the data. By the time the receiver sees the data token, the driver has already shorted the near-end of the differential wires back together, and as the wire itself has significant delay, it can pipeline both data and balanced signals. Due to loss in the wire, balance gates would need to be distributed along the wire, but the control could be propagating down the wire and consistently trailing the data signals.

Although this scheme would maximize token repeat rate, it would also complicate the design. Margins for such a system across wire and gate loading variations would necessarily be large and we predicted significant overdesign. As described in a later section, the system was eventually designed to run at the chip clock rate, making the pre-equalization task naturally at a clock phase edge. Hence, we simply used the pre-existing global clock

#### 5.2. LOW-SWING REPEATER SYSTEMS

signal to control multiple equalization circuits distributed along the wire. These distributed balance devices presented an additional load to the clock.

In our design for a 10mm wire we used five evenly-spaced 3.6  $\mu$ m-wide balance devices, at 0, 25%, 50%, 75%, and 100% down the line. This many devices allowed smaller sizes, but even so, the balance devices had a total gate width double that of the driver. Fewer balance devices would have been correspondingly larger: SPICE simulations confirmed that three equalizers would have been 30- $\mu$ m wide each; two equalizers would have been 140- $\mu$ m wide each; and a single equalizer could not have worked, even at sizes exceeding 250- $\mu$ m. These relationships arise from a cubed relationship between balance driver sizes and the distance between drivers: as the distance between balance drivers grows by *x*, the difference between our 4*FO4*s time limit and the wire *RC* shrinks by  $x^2$ . In addition, the balance driver must drive a portion of the wire *C* in this shrinking time, and that wire capacitance has also grown by *x*. These two terms combine for an  $x^3$  dependency between the balance device spacing and the balance device sizes.

In a real design we would also provide small—1.2  $\mu$ m in the testchip—passgates to a shunt supply at the balance voltage, to fix DC drifts from imperfect balancing.

### 5.2.3 Low-swing receivers

A simple regenerative sense amplifier, drawn in Figure 5.13, followed by an optimized R/S latch [92] performs full logic amplification of low-swing inputs. This section will discuss three design considerations for this particular circuit: common-mode and swing level of the input voltage, offset voltage and mismatch, and clocking.

#### Amplifier delay and common-mode voltage

In Section 5.2.1 above, we discussed the use of NMOS devices at low common-mode voltages, and their drive resistance advantage over PMOS drivers at high common-mode voltages. This should motivate us to use PMOS input gates in the regenerative amplifier to properly receive low input voltages. As it turns out, the amplifier would also prefer low common-mode input voltages.

Considering only the sense-amplifier itself, and not the R/S latch, which has a fixed

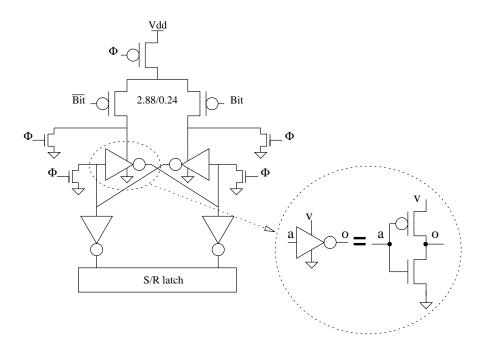


Figure 5.13: Receiver amplifier, in 0.18- $\mu$ m technology

delay, we can model the amplifier latency in two parts. First, the clock trips and the two input gates pull current up each side of the amplifier. These devices are saturated, and so their currents are constant. These currents charge up the inverters from their pre-discharged states. After some time, the amplifier's cross-coupled inverters kick into a regenerative feedback loop and the amplifier switches; during this stage, the input devices play only a very minor role in the amplifier's operation.

A highly simplified model of this amplifier can help illustrate its delay. This model assumes that the devices are velocity saturated; while PMOS gates in today's technologies are not completely velocity saturated, they are close, and this assumption dramatically simplifies the model. In this case, we can write the input voltages as  $V_1 = V_{cm} + \Delta V$  and  $V_2 = V_{cm} - \Delta V$ , and the corresponding currents as:

$$I_1 = K_p(V_{cm} + \Delta V - V_t) = I_{cm} + \Delta I$$
(5.6)

$$I_2 = K_p(V_{cm} - \Delta V - V_t) = I_{cm} - \Delta I$$
(5.7)

$$I_{cm} = K_p (V_{cm} - V_t)$$
(5.8)

#### 5.2. LOW-SWING REPEATER SYSTEMS

$$\Delta I = K_p \Delta V \tag{5.9}$$

The time required for the first stage, as the constant currents pull up either side of the cross-coupled amplifiers, can be written as

$$t_1 = \frac{C_{tot}V_{swing}}{I_{cm}} = \frac{C_{tot}V_t}{I_{cm}} \approx \frac{C_{tot}\frac{1}{4}V_{dd}}{I_{cm}}$$
(5.10)

where the current is the common-mode current; the capacitance is the total load capacitance; and the voltage swing is a threshold drop  $V_t$ . This last term arises because once the output nodes exceed  $V_t$ , the inverters begin to regenerate, driving one of the input transistors out of saturation (the input devices stay saturated for a while, governing the current through the cross-coupled inverters, so this is only an approximation). Here we assume that the threshold voltage is roughly one-fourth the power supply.

The time required for the second regeneration stage depends logarithmically on the voltage gain, where the regeneration time constant is a lightly-loaded inverter delay (0.35FO4s). The voltage gain depends on how far the nodes have split when the regeneration kicks in, and this can be written as

$$t_2 = 0.35 \log \frac{V_{dd}}{V_{start}} \tag{5.11}$$

$$V_{start} = \frac{1}{C_{tot}} \int_0^{t_1} 2\Delta I dt = \frac{2K_p \Delta V}{C_{tot}} t_1$$
(5.12)

We can consider two cases. In the first case, the input voltages are centered about  $0.5V_{dd}$ , so that the common-mode voltage is simply  $0.5V_{dd}$ . Assuming that a threshold voltage is one-fourth the power supply leads to an expression

$$D_{fixedCM} = \frac{2C_{tot}}{K_p} + 0.35 \log \frac{1}{swing}$$
(5.13)

In a full-swing case, when the second term is zero, the first term (an *RC* product) looks like the delay of a three-input NOR, since the circuit has three series PMOS gates. This "NOR" drives a load whose size makes it look like it has a fanout of 2.25. This results in a delay of

about 1.3FO4s.

$$D_{fixedCM} = 1.3 + 0.35 \log \frac{1}{swing}$$
(5.14)

In the second case, the common-mode voltage falls along with the swing, increasing the constant current. In this case, the higher common-mode current pulls up the inverter nodes faster—shorter  $t_1$ -but also gives the differential current  $\Delta I$  less time to split apart the regenerative nodes—longer  $t_2$ . Here, taking  $V_{cm} = V_{dd} - \Delta V$ , so that the common-mode falls with the swing, and again assuming that  $V_t = \frac{V_{dd}}{4}$  leads to an expression

$$D_{variableCM} = \frac{1.3}{3 - 2x} + 0.35 \log \frac{3 - 2x}{swing}$$
(5.15)

Figure 5.14 compares this simple model to simulation, and shows that the amplifier indeeds operates quickest for low common-mode inputs and small swings. Because the amplifier regeneration speed easily outstrips a wire time constant, a strategy for low latency would be to transmit *very* low swings and let the amp, not the wire drivers, do all the work. However, as seen in the next section, swings that are too small run the risk of failure.

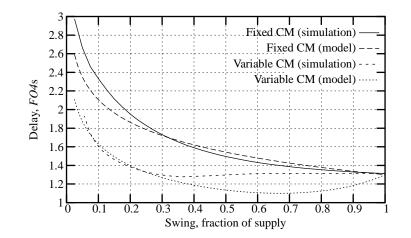


Figure 5.14: Sense amplifier delays

#### **Mismatch and offsets**

The problem with swings that are too small is that they might not overcome inherent, systematic inaccuracies in the receiver [93][94]. If the receiver has a bias to evaluate one way over the other, then any inputs whose difference does not overcome that bias will cause the amplifier to give the wrong answer.

This intrinsic offset in the amplifier is dominated by geometric and  $V_t$  mismatches in the two PMOS input devices, as well as by geometric and  $V_t$  mismatches in the inverter's PMOS devices; the inverter's PMOS devices are important because any mismatch on their source voltage maps directly to their gate overdrive. The NMOS devices have a correspondingly smaller effect on mismatch.

As the common-mode input falls and the current increases, the offset voltage worsens. We can see this by considering just the input pair mismatches in  $V_t$  and  $\beta$ . Suppose  $\beta_1 > \beta_2$ but we still want  $i_2 > i_1$  when input voltage  $V_2 > V_1$ . In this case, we can define

$$\bar{\beta} = \frac{\beta_1 + \beta_2}{2} \tag{5.16}$$

$$\Delta \beta = \frac{\beta_1 - \beta_2}{2} \tag{5.17}$$

(5.18)

If we define  $v_1 = V_{gs} - V_{t1} - V_{tail}$  and  $v_2$  similarly, then when the input voltage  $V_2$  exceeds input voltage  $V_1$ , correct behavior of the pair occurs if

$$i_2 > i_1$$
 (5.19)

$$(\bar{\beta} - \Delta\beta)v_2^{\alpha} > (\bar{\beta} + \Delta\beta)v_1^{\alpha}$$
(5.20)

$$\bar{\beta}(v_2^{\alpha} - v_1^{\alpha}) > \Delta\beta(v_2^{\alpha} + v_1^{\alpha})$$
(5.21)

$$\frac{v_1^{\alpha} - v_2^{\alpha}}{v_1^{\alpha} + v_2^{\alpha}} > \frac{\Delta\beta}{\bar{\beta}}$$
(5.22)

The right side of the last equation is fixed. So as the denominator term on the left side (the total current in the amp) rises with decreasing common-mode voltage, the numerator term (the differential input voltage) must also rise to compensate. Hence the offset increases with decreased common-mode voltage.

The receivers were sized to balance input offset voltage and clocking power. For input devices sized W/L= $2.88/0.24\mu$ m in an  $0.18-\mu$ m technology, SPICE simulations predicted  $3\sigma$  offsets of 86mV. However, we added margin to this offset estimate because it was based on projections and extrapolations from other technologies. Therefore the system assumed offsets of 100mV.

In hindsight, we know that power in the drivers and equalization devices dominate that in the receiver's amplifier, so using larger amplifiers with correspondingly lower offsets would have been a better choice. Table 5.1 shows these estimated tradeoffs as the receiver's device widths vary from normal to 2.4x normal. In this table, amplifier offsets fall by the square-root of the width increase. Adding 15% margin to these offsets for a "minimum safe swing" enables calculations of energy savings, because SPICE simulations show that the wire energy of the baseline system is 18% of the total system energy. Also, the increase in energy cost from the larger amplifier comes from simulations showing that the amplifier accounts for 8% of the total system energy. We see that a receiver about 1.6x the size actually used would have been preferable, although the benefit is fairly small. A further

Receiver	Receiver	Safe	Benefit in	Cost in
Size	Offsets	Swing	Power (from wires)	Power (from rcvr)
100%	86mV	100mV	0%	0%
120%	78mV	90mV	3.42%	-1.6%
140%	73mV	84mV	5.22%	-3.2%
160%	68mV	78mV	7.02%	-4.8%
180%	64mV	74mV	7.92%	-6.4%
200%	61mV	71mV	9.00%	-8.0%
220%	58mV	67mV	9.90%	-9.6%
240%	56mV	64mV	10.8%	-11.2%

Table 5.1: Estimated improvements to receiver sizing

improvement to the receiver circuits would use well-known offset compensation schemes to reduce the offsets further.

#### Clocking

Receivers circuits need to be clocked for power efficiency reasons. Not doing so would otherwise leave the amplifiers in a continually high-gain, high-power state and with little room for power optimization [95]. Clocking the receivers can be done in one of two ways.

First, the sender can package a clock signal to indicate that the receivers should strobe the data. Sending a clock with the data complicates the design: the clock must be fullswing, or else the receiver does not know when to sample and amplify the clock. But as a full-swing signal, the clock signal will have a different delay from the low-swing data lines, and so matching these delays requires careful circuit design: timing the clock too late wastes power, and timing the clock too early may return incorrect data if some of the data wires have not yet overcome their receivers' offsets.

A related question involves "kiting": the full-swing clock consumes enough energy that it should be bundled with a large number of receivers, to amortize its energy cost. But buffering up the clock to drive a wide bank of receivers takes time and hence requires the clock to be launched earlier than the data, or "kited"<sup>1</sup>. Getting the kite delay exactly right might add more design complexity. Suppose the low-swing data is an order of magnitude more efficient than the full-swing clock. To keep the clock power overhead minimal, to under 10% of the total power cost, we would bundle on order of 100 data lines per clock line. In an 0.18- $\mu$ m technology, optimally repeating the clock wire would result in a final gate driver of around 20 $\mu$ m of gate. Assuming the receivers each present a clock load of around 6 $\mu$ m, this results in a required step-up of 30, or 2-3 FO4s of "kiting."

Second, a much-simplified clocking scheme uses the global chip clock as our clocking source. Leveraging this existing, low-skew, globally-distributed clock is a cheat, because it makes the job of the clock distribution network marginally more complicated. However, it also dramatically simplifies the design, and for this reason the final design and testchip used such a global clocking scheme.

By clocking the receivers with a global clock, we by necessity peg the transfer latency to a clock phase. This changes the figure of merit for the interconnect system, from ps/mm units into "distance travelled in a clock phase." Also, clock-gating now requires a more

<sup>&</sup>lt;sup>1</sup>Following Sutherland's asynchronous parlance, we need to *kite* the clock, or "float" it ahead of the data [96].

sophisticated scheme, in which a suppress signal sent in a cycle means that the next cycle's data should not be read.

## 5.3 Putting it all together

To explore this design space, we built simulations of the low-swing signalling system and compared its behavior and performance to that of a full-swing CMOS system. Using a commodity TSMC/MOSIS 0.18- $\mu$ m technology we designed the drivers and receivers to communicate over 10mm of length at a 10*FO4* repeat rate. This closely matched the projected wire characteristics of the Smart Memories target vehicle.

After optimizing both sizing and the driver power supply, the bus needed 0.8pJ/bit, or just about an order of magnitude less energy than the CMOS full-swing bus. This energy cost included all of the clocking overhead. In this system, the drivers used drive voltages of 0 and 0.5V, with pre-equalization to 0.25V, and a developed difference of 90mV by the time the sense-amp fired. Simulations showed that 55% of the energy was consumed in the driver and balance devices, 25% in the sense-amp, and the remaining 20% in the wires themselves. Compared to the full-swing system, this system hit the same timing deadlines—one link per clock phase—but for up to an order of magnitude less power. With lower activity factors, the energy improvement decreases, unless clock-gating implemented via supression signals is used.

These results motivated further study into low-swing interconnect. Simulations can be revealing, but they rarely reveal all of the problems in a design, and so we assembled some testchips containing low-swing wire systems. In the next chapter we will discuss the results of those experiments.

# **Chapter 6**

# **Experimental Results**

The last chapter reviewed design considerations for a low-swing interconnect system, and described how a simulation of such a system demonstrated communication over a 10mm on-chip bus with a 10 *FO4* latency and with up to ten times less energy than standard CMOS repeaters.

While this looked great, simulations often look great. They cannot answer several key questions: Does the system actually work? What are the real sense-amplifier offsets? Did we miss any energy consumers in our calculation of power savings? Therefore we built a testchip to attempt to answer these and more questions. We found that the system communicated without errors over the target bus at the target rate, and that it in fact saved an order of magnitude in power over a full-activity full-swing bus; with lower activity factors on the full-swing bus, the energy improvement decreases. We also measured offsets on a 0.25- $\mu$ m technology and found our design to have >  $3.5\sigma$  offsets under 80mV, with residual uncertainties after offset cancellation to be around 15mV at  $3\sigma$ .

# 6.1 Testchip overview

We used a MOSIS 0.18- $\mu$ m technology, fabricated by TSMC, with six layers of aluminum interconnect. The testchip emerged on a fabrication run significantly faster than previous runs through MOSIS, according to the parametric SPICE models on the MOSIS website. The chip had measured *FO4* delays of 95ps at a voltage of 1.8V and room temperature, a

20% speedup over the previous fabrication runs. The testchip had an area of 9.2mm<sup>2</sup> and was edge wirebonded inside an 84-pin Kyocera ceramic package.

The testchip contained four major experiments. First, we designed a number of lowswing buses of varying lengths and bus construction. We used 5mm, 7mm, and 10mm buses, and in each length we built a single-bit bus, a two-bit bus, and a three-bit bus, in order to see coupling effects. These buses used twisted wires for differential noise rejection, but we also duplicated the 10mm bus once without any wire twisting to see the effects of the twists. Because the die itself had edge lengths of only 3mm, the buses needed to be serpentined back along themselves to reach their full length. This also allowed the drivers to be physically close to the receivers.

These low-swing buses had isolated power supplies to enable accurate energy measurements. We also attached comparators between the drivers and the nearby receivers, so that we could compare what had been sent with what had been received. These error-checkers would power up a board LED if they detected a mismatch, although they could not identify which bit or clock generated the error.

Second, we constructed full-swing buses, sized and repeated for a minimum energydelay product, to send signals at a 10 *FO4* latency. These buses, using standard CMOS inverters as repeating elements, spanned 5mm and 10mm lengths and used the same errorcatching circuits as the low-swing buses. They enabled a fair apples-to-apples power comparison between low-swing and full-swing interconnect schemes.

Third, to explore the offsets present in our sense-amplifiers, we constructed a large array of them, all commonly-clocked and all able to shift out their outputs in a chain. By varying the input voltages we could use this to gather statistics on amplifier mismatch. Unfortunately, we mis-designed the scan chain, breaking it logically, and so we had to respin just this particular block onto another vendor's  $0.25-\mu$ m technology in order to see any results. To first order, the  $V_t$  mismatch between two devices does not change if we scale their widths and lengths: the effect of decreased area is countered by the effect of scaled gate oxides. Therefore results from the  $0.25-\mu$ m chip are still interesting.

Fourth, we included a self-contained block of oscillators in order to test the native speed of the chip. This block had its own pads and could be wire-bonded separately on a few parts per lot in order to save pads for the rest of the chip.

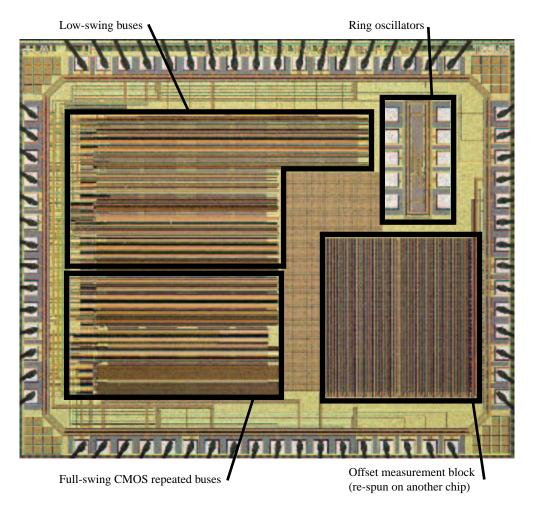


Figure 6.1: Chip photomicrograph

Figure 6.1 shows a die photomicrograph of the chip, with the four major blocks outlined in black.

# 6.2 Bus experiments

The low-swing buses followed the guidelines listed in the previous chapter. A four-way ping-pong device connected to each bus driver's input, in which four scan-loadable flip/flops alternately fed the driver on four successive clocks. This way, any four-bit pattern, including the "quiet-low" pattern of 0000 and the "quiet-high" pattern of 1111, could be transmitted down the wire. The drivers used NMOS linear-mode drivers, with low common-mode voltages, overdrive pre-emphasis and pre-equalization between data tokens. For the overdrive, the buses aimed for a common-mode equalization voltage of 0.25V with drive power supplies of 0 and 0.5V.

The wires ran as a differential pair, twisted along the route to reject outside noise. An extra version of the 10mm bus, without any twisting but identical otherwise, provided a measure of the twisting effectiveness. Balance devices sat every 25% down the wires to achieve fast pre-equalization, along with trickle devices connecting the wires to a  $V_{balance}$  supply to correct mismatched pre-equalization gates.

A receiver, clocked off of the chip clock, amplified a new token each phase, presumably after the development of 90mV of signal. The receiver fed the outputs into an errorchecking comparator that had stored the driver's sent data from the previous clock phase. On any mismatch a pin was driven high to light an LED on the board.

To make the chip more realistic, a number of large inverters placed below the buses could be clocked to inject noise into the system. These inverters were turned off during power measurements, because they used the same supply as the buses, and they do not show up in the figures below, although they were active during exploration of the voltage and speed margins.

For test and measurement, many on-chip voltage sampling oscilloscopes captured waveforms for validation of performance, and they provided the figures in this chapter. Their design will be discussed briefly below.

The reliability of the link was not well instrumented, making unavailable traditional

reliability metrics like pseudo-random data bit error-rate. Claims that the low-swing bus "worked" meant running the chip for several hours and not seeing any errors; in one case the chip ran over a weekend with no errors. However, the bit patterns available for test were only simple repeating four-bit patterns and not more sophisticated PRBS generators.

As the design of the low-swing buses targeted transision of a new token across the 10mm link every 10 *FO4*s, the full-swing CMOS repeated buses next to them were designed to hit the same 10 *FO4* latency for the same length hop. This would provide the fairest energy comparison. In addition, because the low-swing buses used two non-minimal wires per bit, the full-swing buses were given equivalent wire resources and thus had nice, fat, widely-spaced wires. This was again to ensure a fair energy comparison.

### 6.2.1 Overhead

Because full-swing CMOS repeaters do not need a clock, they have no need to cross the entire hop in the 10 *FO4*s: unlike the low-swing repeaters, the penalty for placing a repeater between the hop end-points costs only area and via congestion and does not require generating any special phase-dividing timing signals.

As a result, the full-swing CMOS buses used repeaters placed for optimal energy given a 10 *FO4* latency target, and therefore broke the 10mm wire into three equal sections. In simulation we found that if repeated for optimal latency, they would have reached across the 10mm hop in just over 7.5 *FO4*s, so that they had a native 25% speed improvement over our low-swing buses. However, on the chip they were throttled-back somewhat to hit 10 *FO4*s.

The active devices in the low-swing experiment consumed nearly 75% more area than the active devices in the full-swing CMOS buses. The additional area went to the complex receivers, pre-equalization clamps, and drivers. However, the total area of the bus is completely dominated by the wires and not by the active devices, so the overall area penalty is less than  $\frac{1}{2}$ %.

#### 6.2.2 Measurement circuits

Because this chip explored low-voltage signaling, seeing how the actual waveforms behaved would offer insight into why the chip either worked or failed. To get this on-chip visibility without invasive probing techniques, we made extensive use of on-die sampling circuits (described originally in 1998 [98]) to probe the analog waveforms of each of the buses. Not only did the samplers work well enough to provide invaluable visibility into the die, but we also gained some more experience with using these samplers.

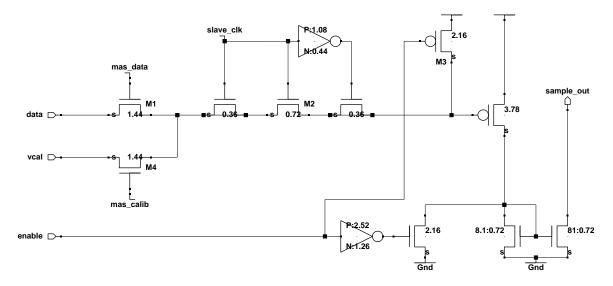


Figure 6.2: Master-slave variant of on-die sampler circuit, sized for a  $0.18 - \mu$  technology

Figure 6.2 shows the basic schematic of the sampler. The sampling head, transistors M1 and M2, perform an "analog flip-flop" function, moving a fascimile of the test voltage onto the gate of the PMOS amplifier M3. Note that transistor M2 is surrounded on either side by half-width devices. These help cancel noise coupled from the slave clock to the slave node, as well as from the master clock to the master node.

By boosting the voltage on the gates of M1 and M2 as well as on the source of M3 to  $V_{dd} + V_t$ , the circuit can sample voltages up to the power supply rail. M3 converts this test voltage into a current, which gets multiplied up twice and sent off-chip into an oscilloscope. To help translate the resulting output voltage into an internal voltage, a separate calibration transistor M4 can be enabled to sample an externally controlled DC voltage.

The sampler can be run in one of two modes. In one mode, the sampler clock and the chip clock have the same frequency, and the chip is placed in a repetitive pattern. Over many cycles of operation, the sampler continually samples the same voltage, ensuring that the gate of M3 gets charged to the full voltage on the node under test. We can record the resulting output voltage and translate it to an on-chip voltage. Then, by stepping the phase relationship between the two (identical frequency) clocks, we can map out another voltage. By repeating this process, we can eventually draw a very accurate waveform.

In the other mode, we set the sampler clock and the chip clock to have very slightly different frequencies. Now, each cycle the sampler examines a slightly different point on the test waveform, and by dialing out the time base of the oscilloscope to the beat frequency of the two clocks, we will get a (slightly inaccurate) picture of the real on-chip waveform on the oscilloscope. In this mode, charge-sharing across gate M2 limits the bandwidth of the sampler, although changing the beat frequency can control this bandwidth limit. Jitter in the sampler clock, on the other hand, blurs the sampled signal over the jitter interval, making it a more serious bandwidth limiter.

In this testchip, not every wire was sampled, but those that were had two samplers each: one on the true wire and on one the complement. The true wire used a sampler very similar to those as used on previous chips (see Figure 6.2), while the sampler on the complement wire inserted a simple buffer-amplifier in-between the NMOS input passgates in an attempt to eliminate the pole that arises from charge-sharing across the slave NMOS passgate. In every case, however, the waveforms from the two samplers did not noticeably differ when we overlaid them (by using one sampler with one data set, and then the other sampler with the complement data set). This indicated that the charge-sharing pole was unimportant. The difference in loading on the true and complement wires from the two different types of samplers, or simply between an inactive and an active sampler, was negligible.

A more robust sampling circuit would have used a differential read path on both of the true and complement wires. However, the testchip used separately-sampled differential wires with two different types of single-ended samplers. This was primarily to reduce the likelihood of design errors in the samplers preventing *any* measurements.

By using an extra sampler to sample the chip clock, we gained a timing reference. Sampling the chip clock returns a clock at the beat frequency between the chip and sampling clock, and we can use this beat clock as the trigger for the oscilloscope. This aligns all sampled data in time, allowing us to line up graphs properly and calculate delays.

Per-sampler calibration was extremely important. The calibration curves, though constant for each sampler, varied significantly between sampler sites. Especially for measuring voltages that hover near 0V, eliminating uncertainty from uncalibrated devices became paramount.

Finally, we made a nearly unrecoverable error in the sampler design. The final output node, carrying the magnified current from the two series current mirrors, routed to a chip pin over a long and winding route. This meant that significant current out of the sampler would cause significant iR drop along that wire, thereby pushing the current mirror out of saturation. This clamped the output current, effectively making the sampler much less responsive to any change in sampled voltage, if the voltage were low enough. As the nodes in question were small-swing at low voltages, this was very nearly fatal.

The fix was to not ground the current mirror output at ground with the oscilloscope input, but rather to drive it down to -2V through a discrete  $50\Omega$  resistor. This way, the *iR* drop along the output wire would still give the current mirror sufficient headroom to stay saturated and maintain sensitivity to the sampler.

### 6.2.3 Performance and results

The low-swing buses performed at 10 *FO4* per cycle, or just about 1GHz, but required a higher driving voltage than predicted. Instead of the planned 0.5V drive voltage, the system required just over 0.6V of drive voltage, and a pre-equalization voltage of 0.3V. In other words, the error LED on the board did not remain off until the drive voltage reached 0.6V.

Upon exploration of the layout, we found that our clock spine, distributing the clock to the various drivers and receivers, was thin, scenically-routed, and replete with clock skew. Although we did not perform FIB experiments to adjust the clock routing, we speculate that this contributed significantly to the bus failures below 0.6V of drive voltage. Unfortunately, we were unable to observe any clock skew because we had no samplers sampling chip clocks near the actual receivers, nor any samplers looking at signals derived from the receiver clock.

#### **Measurement pictures**

Figures 6.3 and 6.4 show sampled oscilloscope measurements taken from the 10mm lowswing bus with three bits on the bus. The middle bit is plotted and the two outer bits are quiet.

Figure 6.3 shows the system running at 500MHz. The outer curves show the waveform at the driver end of the wire, and the inner curves show the waveform at the receiver end of the wire. At 500MHz, or about 20 *FO4*s per token, the system has plenty of time to drive the value down the wire, and the far end of the wire achieves nearly the full voltage split. Due to very generous timing margins, the system worked fine down to a drive voltage of 0.4V and a pre-equalized balance voltage at 0.2V. At 0.5ns and 1.5ns the wires take sudden jumps down and up in common-mode voltage; this is due to the clock signal coupling into the data lines. Although the coupling may seem dramatic, the movement is only around 100mV, which is small compared to the 1.8V swing on the clock wires themselves.

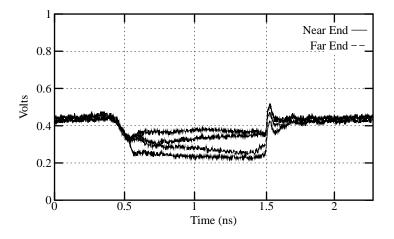


Figure 6.3: Low-swing bus at 500MHz

Figure 6.4 shows the same bus running now at 1GHz. Again, the outer curves show the driver end of the wire and the inner curves the receiver end of the wire. Because the system speed is much faster, the receiver end of the wire barely splits past 90mV before the clock phase ends and the wire begins to pre-equalize. By contrast, the driver-end of the wire opens up to a wider voltage. As mentioned above, this system required a drive voltage

greater than 0.6V before operating without errors; this picture was taken at a drive voltage of 0.7V. The clock coupling is equally evident in this picture, as well.

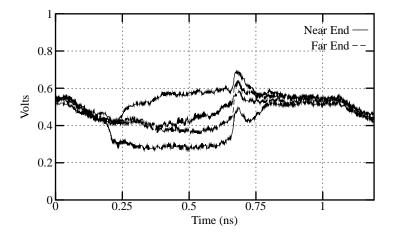


Figure 6.4: Low-swing bus at 1GHz

#### **Power measurements**

By measuring the current required by the chip when only a single bus is operating, we can estimate the energy per bit required to communicate over that bus.

We first operated only the CMOS full-swing bus, and disabled all of the low-swing sections of the chip. Programming the four-bit ping-pong drivers with a "0000" pattern allowed a measurement of the energy required by the error-checking comparators and the four-bit ping-pong drivers themselves. This energy could then be subtracted from all further measurements. Next, we programmed the CMOS full-swing bus to send real data, and measured the energy required to be 9.84pJ/bit over the 10mm link at  $V_{dd} = 1.8$ V.

Next we ran the low-swing link and measured the energy consumed. We ran it at clock speeds of 500MHz and 1GHz. Energy consumed for the full-swing CMOS bus does not depend on clock frequency; however, for the low-swing buses it does.

Figure 6.5 shows the measured energy consumption for the three different length buses at 500MHz, and Figure 6.6 shows the measured energy consumption for the buses at 1GHz. At 500MHz, the system is running slow enough that the receiver end of the wire can split all the way to the drive voltage extremes before being pre-equalized for the next bit. By

contrast, at 1GHz, the system is running fast enough that the receiver end of the wire barely reaches the required 90mV split. As a result, at the slower clock speed the system consumes more energy for the same drive voltage.

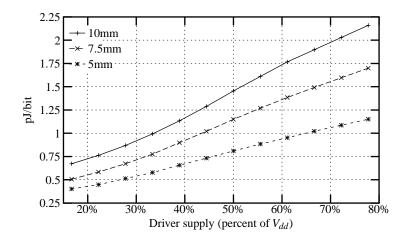


Figure 6.5: Measured energy consumption of low-swing bus at 500MHz. Full-swing bus had measured energy consumption of 9.84pJ/bit.

These figures sweep the driver voltage on the X-axis. As these voltages decrease towards the left side of the figures, the energy falls, and eventually the graphs stop. This indicates that at 500MHz, the system worked with drive voltages down to 0 and 0.3V, with the pre-equalization voltage at 0.15V. At 1GHz, the 10mm wire worked only down to 0.6V, while the 7mm and 5mm buses, being shorter and having less wire *RC*, worked at lower voltages.

At the lower end of the drive voltage spectrum, the low-swing buses required around 0.8pJ/bit, or more than an order of magnitude improvement in energy consumption versus a constantly toggling full-swing bus (carrying bit patterns of "010101..."). If the full-swing bus had a lower activity factor, this energy benefit would decrease. Regaining an energy improvement in this case would require clock gating using supression signals, which we did not build.

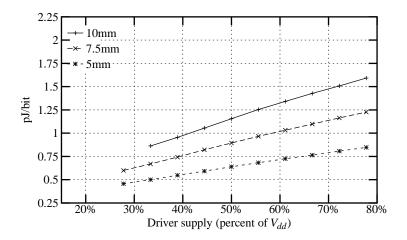


Figure 6.6: Measured energy consumption of low-swing bus at 1GHz. Full-swing bus had measured energy consumption of 9.84pJ/bit.

#### Efficacy of wire twisting

The 10mm bus was twisted along its length four times to reject outside noise. If we label the wires in the three-bit bus *A*,*B*, and *C*, we have the six wires *A*, $\bar{A}$ ,*B*, $\bar{B}$ ,*C*, and  $\bar{C}$  in a row.

Without any twisting, the above wire sequence remains the same all the way down the wire. In this case, if we send a "1" down all three wires, then  $\overline{A}$  will fall as *B* rises, and *C* will rise as  $\overline{B}$  will fall. This pattern generates the worst-case "hurt" coupling, because *A* and *C* will restrict the opening of *B*'s data eye. Conversely, if we send a "0" down *B* as we send "1"s down *A* and *C*, we get the best-case "help" coupling, as *A* and *C* aid in opening *B*'s data eye.

With all wires twisted, all these effects should be nullified, at least to first order. Because we also built a 10mm bus without any twists, we were able to view the actual effects of such a twist. Figure 6.7 and 6.8 show the probed waveforms at the receiver side only. In this experiment the buses ran at 1GHz. In Figure 6.7, the two waveforms are difficult to distinguish; there is no difference between the "help" or "hurt" data patterns.

By contrast, the untwisted wires show that the hurt pattern has a worse data eye than the help pattern. Interestingly, the untwisted wires, even with the hurt pattern, has a data eye very close to that of the twisted wires. After some investigation we realized this resulted from sloppy twists. Because the process uses aluminum interconnect, vias consist of

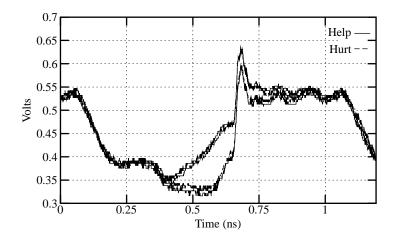


Figure 6.7: Fully-twisted wire: no difference between "help" or "hurt"

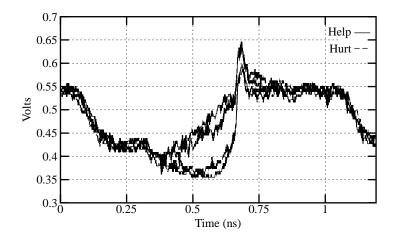


Figure 6.8: Untwisted wire: a difference between "help" and "hurt"

highly-resistive tungsten. The twists in the bus used only single vias and not wide arrays of them, so that the twisted wires, although good at rejecting coupled noise, also had an extra  $40\Omega$  of resistance, assuming via resistance of  $10\Omega$  each. This led to a non-trivial increase in total resistivity and implied that the twisting seemed effective, if poorly implemented.

# 6.3 Offset experiments

In the offset experiment, originally designed for the MOSIS/TSMC 0.18- $\mu$ m technology but ultimately fabbed on a National 0.25- $\mu$ m process, we concentrated 1024 sense-amplifiers together in an array. For these amplifiers, all of the positive inputs were tied together and all the negative inputs were tied together. All of the clocks were also shorted. The outputs were fed into flip-flops that serially connected together in a long scan chain.

With over 1024 amplifiers per die, and with ten die, we had numbers significant to greater than  $3\sigma$ , although over only one lot.

If the inputs were 1V on the positive input and 0V on the negative input, and the group of 1024 sense-amplifiers were clocked, they would all certainly agree, so that the scanchain would hold 1024 "1"s on the output. Conversely, if the inputs were 0V on the positive input and 1V on the negative input, and the entire group of sense-amplifiers were clocked, they would again agree and the scanchain would hold 1024 "0"s.

If the inputs, however, were 0.23V and 0.21V, any sense-amp with an intrinsic offset greater than 20mV would not agree with the rest of the sense-amps, leading to a non-zero number of both "0"s and "1"s in the scanchain. For this particular example, we would define the differential voltage  $\Delta v = 20$ mV and the common-mode voltage  $v_{cm} = 0.22$ V, and record the number of disagreements for that combination of  $\Delta v$  and  $v_{cm}$ . By repeating for a wide range of  $\Delta v$  and  $v_{cm}$  we can build up statistically significant parametrized offset measurements.

## 6.3.1 Offset measurements

Figure 6.9 shows the results of this experiment by plotting the disagreements. The x-axis shows ranges of  $\Delta v$ , and the various curves represent different common-mode voltages

#### 6.3. OFFSET EXPERIMENTS

 $v_{cm}$  from 0.1V to 1.1V; higher common-mode voltages are the curves towards the center. When the input differential voltage is large and negative, all the sense-amplifiers agree, giving no disagreements. When the input differential voltage is large and positive, all the sense-amplifiers again agree, leading to no disagreements. As we move towards the middle from both extreme left at some  $\Delta v_{left}$  the disagreement curve rises above zero. The same happens as we move from the extreme right, at some  $\Delta v_{right}$ . We define the offset of the amplifier to be the difference between these voltages.

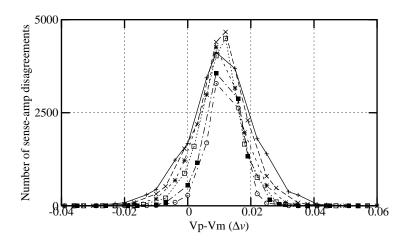


Figure 6.9: Measured scan-chain disagreements. The x-axis is  $\Delta v$  and the different curves are for  $v_{cm}$  ranging from 0.1V to 1.1V (higher  $v_{cm}$  towards the center).

Figure 6.10 shows the same data but with common-mode voltage on the x-axis and offset on the y-axis. This graph more clearly shows the increase in offset voltage with lower common-mode—the input gates are PMOS, so with low common-mode, they have more current. At the lower common-mode voltages used in our design, the offset voltage ended up under 80mV, comparing favorably with our design targets of 90mV<sup>1</sup>. Note that the simulations predict a much larger variation of offset due to common-mode voltage. That the measured data shows a smaller variation means that the true  $\beta$  mismatch is smaller than the 3% we assumed from extrapolation from other technologies; using a  $\beta$  mismatch closer to 1.5% gives curves that show the same common-mode sensitivity.

<sup>&</sup>lt;sup>1</sup>The two numbers should not match, because our initial designs aimed for a 0.18- $\mu$ m technology and these simulations and measurements were done for a 0.25- $\mu$ m technology. They should, however, be close.

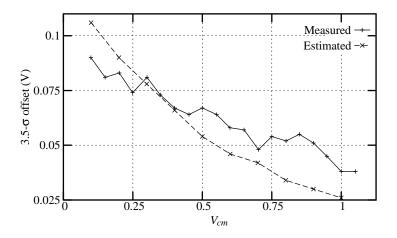


Figure 6.10: Measured offsets versus common-mode voltage

Another set of 1024 sense-amplifiers were built in the same array, but this time with minimum-length transistors, to see how much worse the matching became for smaller devices. Figure 6.11 shows worse results for these sense-amplifiers. Their input gates have the same W/L ratio as the original sense-amplifiers, but with minimum channel length, such that the square root of their area ratios is around 70%. The offsets, as predicted by theory, are approximately 30% higher [93][94].

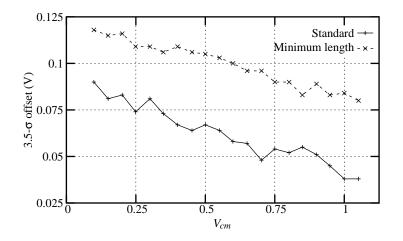


Figure 6.11: Measured offset of minimum channel length amplifiers

Finally, "tilt" across the offset block was negligible. The block itself was relatively

small, leading to fairly uniform offsets across it. Figure 6.12 shows a three-dimensional plot of the offsets by position; the random variation from sense-amp to sense-amp outweighs any offset gradient across the bounding rectangle.

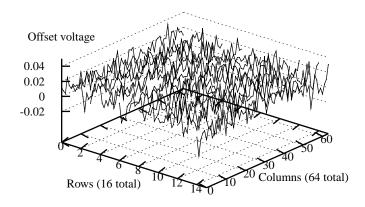


Figure 6.12: Measured offsets in their physical locations (in volts)

## 6.3.2 Systematic errors

In Figure 6.9, the center of the "hump," where most of the sense-amplifier disagreements occur, sits not at  $\Delta v = 0$ V, but rather closer to 10mV. This indicates that the sense-amplifiers have some systematic offset, most likely due to imperfectly balanced layout. The amplifiers were laid out carefully, but not carefully enough to avoid biases or offsets due to manufacturing variations.

Figure 6.13 shows a plot of the sense-amplifier design from the polygon layout editor Magic. Three layout features contribute to a systematic offset:

 The tail device, at the center top of the diagram, ties the power supply to the tail node, but then this current will flow both left and right to each of the input gates. Metal misalignment may effectively move the tail device *closer* to one side than the other. A dual-legged structure, with the tail device split between the two sides, would avoid this problem, and would also allow the two input gates to be closer together, eliminating any minor "tilt" in oxide thickness or implant across the amplifier.

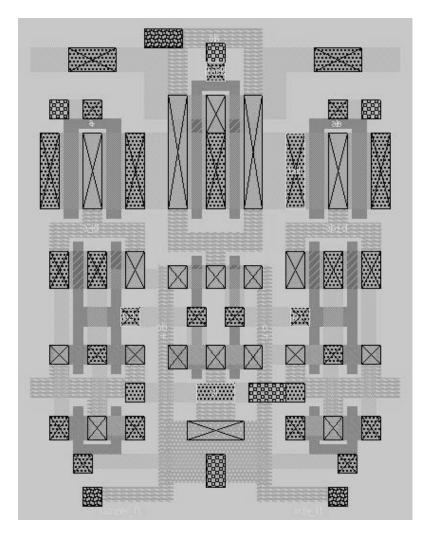


Figure 6.13: Sense amplifier

- 2. The two output inverters, towards the middle of the layout, lack common-centroid layout, so that their current flows are not balanced in both directions. They are the only critical devices not using common-centroid layout. Hence a poly misalignment towards, for example, the left, will move the left vertical poly towards its output, potentially increasing its gate-drain capacitance. At the same time, the right vertical poly will move away from its output, decreasing its gate-drain capacitance. This capacitive imbalance may lead to a small bias.
- 3. The inverters below the input gates "cross-over" using Metal2 strips. One Metal2 strip runs from the upper left of the left-most inverter to the center of the right inverter; the other Metal2 strip runs from the lower right of the right-most inverter to the center of the left. This arrangement means that the left inverter can pullup the right inverter's gate slightly faster than it can pulldown, and the right inverter can pulldown the left inverter's gate slightly faster than it can pullup. Because these inverters are regenerative, this leads to a mild bias.

Creating perfectly balanced layout, especially when the design uses cross-coupled devices, is difficult, as our attempt shows. Although these layout bugs seem obvious now in hind-sight, we noticed them only after knew we had a systematic offset in the design.

### 6.3.3 Residual offsets after compensation

We did not introduce offset compensation schemes into this design, although they would dramatically reduce the input-referred offsets of the sense-amplifiers. We can estimate their effect by a simple further experiment.

If we isolate one sense-amplifier in the test block and feed it input voltages at its flip point, we effectively simulate zeroed inputs into a perfectly offset-compensated amplifier. Now if we clock it repeatedly without changing the inputs, we can count the number of times it returns a "0" and "1." This resulting distribution shows the residual error in the amplifier only, including statistical sources like  $\frac{1}{f}$  noise and thermal noise.

Examining sixteen amplifiers picked from one particular column in the array and clocking them each 250 times over a range of input voltages, we can see their error distributions

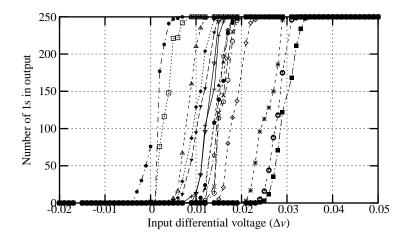


Figure 6.14: Residual noise error distributions

in Figure 6.14. These curves should all be error functions, and we can fit them to the expression below. Here  $\mu$  is the mean, or actual offset voltage, and  $\sigma^2$  the variance of the residual error distribution.

$$Num_1 = 125(1 + erf(\frac{\Delta v - \mu}{\sigma^2})) \tag{6.1}$$

To see how well the graphs in Figure 6.14 conform to error functions, we can replot these graphs mapped to a straight line  $y = \sigma^2 \cdot x + \mu$ . If the graph is a perfect error function, it should be collinear to this line. Figure 6.15 shows this for three example sense-amplifiers. As the figure shows, the results match an error function.

The individual  $\sigma$  terms over the entire array of 1024 amplifiers are plotted below in Figure 6.16. Taking the very worst standard deviation results in a 6- $\sigma$  term of 15mV, as the maximum  $\sigma$  over the array is 2.5mV. Using this 6x multiple of the very worst  $\sigma$  is only slightly conservative.

Because perfect offset compensation is a myth, we assume around 10mV of residual offset compensation noise, leaving an offset-compensated residual offset of about 25mV to 6- $\sigma$ . In other words, even with offset compensation in the sense-amplifiers, there are sufficient statistical noise sources to mandate at least 25mV of signal.

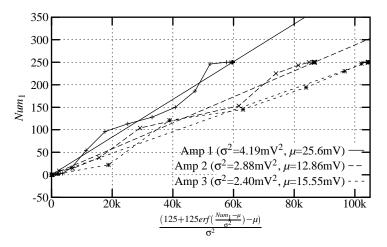


Figure 6.15: Matching the residual noise to an error function

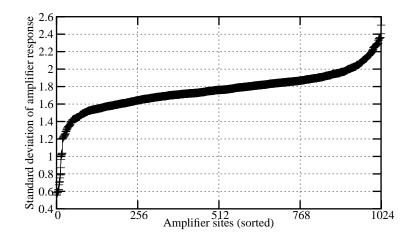


Figure 6.16: Standard deviations across the entire array of amplifiers

## 6.3.4 Offset experiments we wish we had done

Ideally we could have varied the turn-on ramp rate of the amplifier's tail device. Ramping it extremely slowly negates capacitive imbalance in the device's loads, and ramping it extremely quickly emphasizes any such capacitive imbalance. Adjusting the ramp rate would have offered a way to measure this capacitive imbalance. Unfortunately, the clock driver strength and power supply was fixed for the design.

Also, while measuring the residual noise in the sense-amplifier offers insight into the system's reliability, being able to measure the residual error response for a *real channel* would be more interesting. In other words, doing the same experiment with real wires, real crosstalk, and real noise-injecting inverters underneath the bus would have been more realistic and thus more relevant. Unfortunately, we had no easy way of measuring residual noise for a receiver on an unused bus somewhere else in the design.

# 6.4 Summary

This chapter described the implementation and measurements of a low-swing bus testchip. One major experiment on the chip compared the energy requirements of several low-swing buses of varying width and varying lengths to full-swing CMOS buses designed to provide fair energy comparisons. When running at their target latency of 10 *FO4* for a 10mm on-chip hop, the full-swing buses were slightly throttled back from their top velocity, but they also consumed an order of magnitude more energy than the low-swing buses. The low-swing buses required negligible increased area.

A second major experiment, re-spun on a  $0.25 \ \mu m$  technology to fix a logic error, derived statistics for the intrinsic offset and mismatch for the sense-amplifiers used in the design. Our design number of 90mV held up, as the mismatch to >  $3.5\sigma$  was under 80mV. By repeatedly sampling sense-amps held at their offset-zeroed input voltages, we could also determine the statistic residual noise to be 15mV at  $6-\sigma$ .

The engineering challenges in building such a system are nontrivial. For example, our sense-amplifiers, despite careful layout, still suffered from a systematic offset of about 10mV. Nonetheless, these experiments validated that not only could the low-swing system

## 6.4. SUMMARY

work in simulation, but also on a real chip, and with substantial energy savings.

# Chapter 7

# Conclusions

This work described the future of wires, and what designers can do about it. Wire delays scale up as gate delays scale down. Understanding these trends quantitatively requires us to forecast wire and gate characteristics and to combine them into performance metrics. In Chapter 2 we used wire geometry information from the SIA roadmap to predict a range of future wire resistances and capacitances; resistances scaled up linearly with time while capacitances changed only slightly. Using a simple model for gate delays, we showed how the ratio of wire delays to gate delays either scales slowly for scaled-length wires or grows rapidly for fixed-length wires.

That future on-chip wires will display this duality, of "fast local wires" contrasted with "slow global wires," affects how we approach VLSI designs. Chapter 3 considered two such design implications. First, CAD place-and-route tools must improve. With growing die complexity and more local blocks gathered on a chip, maintaining constant design time and designer productivity requires fewer wire exceptions per block and thus improved CAD tools.

Second, modular architectures can effectively exploit the dual nature of wires. Fast local wires in small compute cores, connected by slow but high-bandwidth wide global buses, allow both high performance and high integration. Such modular architectures connect their global wires into regular on-chip networks, using wide buses of high bandwidth to offset long wire latencies. Using such wide and long global buses can burn a great deal of power: the global network in the Stanford Smart Memories modular architecture can consume up to 100W at peak bandwidth, if built with traditional delay-optimal CMOS repeaters. Chapter 4 considers the energy efficiency of such standard repeater architectures. By accepting a 10% penalty in performance, traditional repeaters can be sized and spaced to save nearly 30% in energy.

This 30% of energy savings is not a lot. Chapter 5 discusses techniques for running global wires at a reduced voltage, thereby saving considerably more energy. Drivers using NMOS pullups and pulldowns swing differential twisted wires between ground and a drive voltage significantly higher than the target voltage swing required by the receiver, for decreased delay. Between transmitted bits, pre-equalization devices balance the differential wires together; because we cannot overdrive this pre-equalization, we distribute balance devices over the wire length. Receivers, made of simple latching amplifiers, use the system clock, exploiting the synchronous nature of the global network. Their offsets fundamentally limit how low of a swing we can use on the wires.

We roughed out these ideas in simulation, finding an order of magnitude in energy savings for no effective slowdown of the Smart Memories global buses. Chapter 6 describes the chip we built to more fully test these low-swing buses. This chip validated a 10x savings in energy while running 10mm per clock phase. Another experiment validated our < 90mV offset estimates, and using that same experiment we could determine that if we had perfect input offset compensation, the residual input uncertainty in the amplifier was still around 15mV. The test chip demonstrated that relatively simple low-swing architectures do work and provide substantial energy savings for global on-chip buses.

This work can continue in many different directions. First, the low-swing architecture can be modified to eliminate the receiver's dependence on a global clock. By sending a strobe signal along with the data we can self-time the wire system, although due to rampup delay at the receiver this strobe signal will have to be sent slightly ahead of the data for best performance. This allows some timing leeway in the global networks and if combined with some form of handshaking between modules, can free the global network from a strict relationship with the chip clock. Whether or not this flexibility is useful depends on the system. Second, a careful reckoning of receiver offsets could both modify the existing design as well as add offset compensation schemes. While the swing level is ultimately limited by residual uncertainties (15mV in our design), this number is design-dependent as well. Improving the receiver's uncompensated offset as well as compensated residual uncertainty would let us push the voltage swing and energy savings to extremes without worrying about the non-scaling behavior of transistor offsets.

Third, formulating the low-swing ideas and techniques into a methodology or CAD framework would dramatically improve its usability. Chip builders can be astonishingly conservative, because the extra cost of more robust system components such as power supplies or heat removers is much less than the cost of a yield drop or of field escapes due to sensitive designs. In high-volume manufacturing, technical design considerations often play second-fiddle to yield, manufacturing, and margins, so formalizing this exploratory research will enhance its applicability.

## Appendix A

## Fanout in a buffered repeater

If we have a two inverters as a repeating stage, what should be the fanout between these stages? Compare a fanout of four versus a fanout of one. The FO4 case has a longer intrinsic delay and hence farther-spaced repeaters: fewer but slower stages. The FO1 case has a shorter repeater delay and hence closer-spaced repeaters: more but faster stages. Which is faster?

In any repeated system, if we think of a stage delay as broken into a gate-drivinggate delay ("GateGate"), a gate-driving-wire delay ("GateWire"), a wire-driving-wire delay ("WireWire"), and a wire-driving-gate delay (WireGate), we know that at the delay-optimal length and driver size, we have

$$GateGate = WireWire$$
 (A.1)

$$GateWire = WireGate$$
(A.2)

Let the fanout within the repeating stage be f. Then the GateGate delay is FOf + FO(1/f), or, when normalized to a FO4, (f + 1/f + 1)/4.5 (assuming that diffusion capacitance is half of gate capacitance). Thus at the delay-optimal length, and assuming center-skewed devices and a full power supply,

GateGate = 
$$\frac{f+1/f+1}{4.5}FO4$$
 (A.3)

WireWire = 
$$0.35R_w C_w l_{opt}^2$$
 (A.4)

$$l_{opt} = \sqrt{\frac{f+1/f+1}{0.35 \cdot 4.5}} \sqrt{\frac{FO4}{R_w C_w}}$$
(A.5)

The other half of the story comes from equating the GateWire and WireGate delays. Because of the repeater's internal fanout, the gate drive is f times more powerful than its input capacitance would indicate:

GateWire = 
$$\frac{R_v}{fw_{opt}}C_w l_{opt}$$
 (A.6)

WireGate = 
$$R_w l_{opt} 3C_g w_{opt}$$
 (A.7)

$$w_{opt} = \sqrt{\frac{R_v C_w}{3f R_w C_g}} \tag{A.8}$$

Rewriting and then dividing the stage delay by  $l_{opt}$  for delay per unit-length gives us the following. Note that all of this can be restated by using the general repeater model with a = 3, b = 1.5f, r = 1/f, and d = (f + 0.5)/4.5.

GateWire = 
$$0.360\sqrt{\frac{f+1/f+1}{f}}FO4$$
 (A.9)

Delay/unit-length 
$$\propto \sqrt{\frac{f+1/f+1}{2}} + \sqrt{\frac{1}{f}}$$
 (A.10)

We can also specify the energy of a buffer by using e = 3(f + 0.5), giving a total energy of

Energy 
$$\propto 1 + 1.5 \frac{f+1}{\sqrt{2(f^2 + f + 1)}}$$
 (A.11)

Whether we consider just delay, or the energy delay product, the optimal fanouts sit around 2.25–2.5, with fairly low sensitivity. This doesn't change much with other  $\hat{l}$  or  $\hat{w}$  design points, either. At an energy delay optimal of  $\hat{l} = 1.5$  and  $\hat{w} = 0.4$ , the optimal fanout is still between 2.5 and 3.

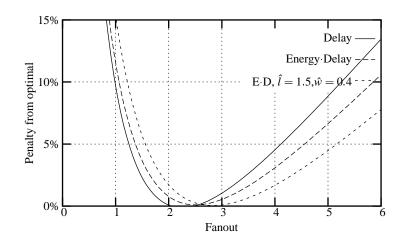


Figure A.1: Delay and Energy-Delay vs. buffer internal fanout

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